

# Theory of $D^0 - \overline{D}^0$ mixing

Alexey A Petrov

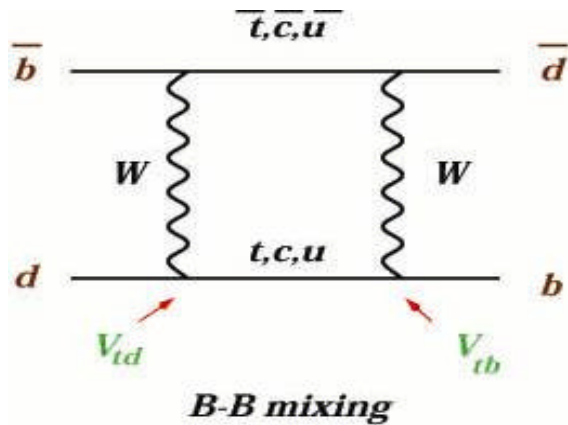
Wayne State University

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Based on works with S. Bergmann,  
A.F. Falk, Y. Grossman, Z. Ligeti, Y. Nir

# Introduction: why do we care?

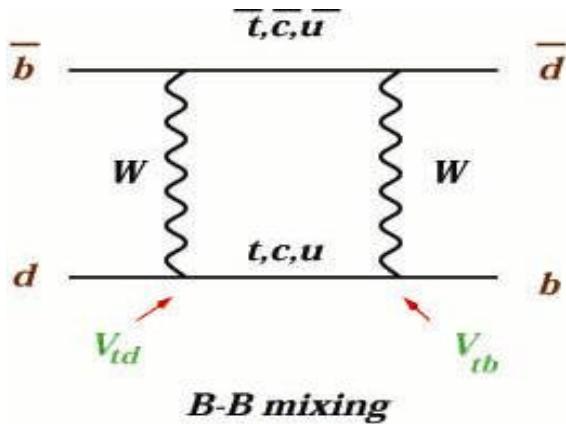


B-mixing:

- ➔  $\Delta Q=2$ : only at one loop in the Standard Model
- ➔ GIM mechanism:  $rate \propto m_1^2 - m_2^2$
- ➔ sensitive to ultra-heavy particles in the loop

Expectation: *rate* is “large” in B system

# Introduction: why do we care?



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$$|B(t)\rangle = \begin{pmatrix} a(t) \\ b(t) \end{pmatrix} = a(t) |B^0\rangle + b(t) |\overline{B^0}\rangle$$

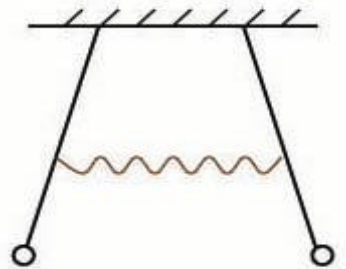
Time-dependence: coupled Schrödinger equations

$$i \frac{\partial}{\partial t} |B(t)\rangle = \left( M - \frac{i}{2} \Gamma \right) |B(t)\rangle = \begin{bmatrix} A & p^2 \\ q^2 & A \end{bmatrix} |B(t)\rangle$$

Diagonalize: mass eigenstates  $\neq$  flavor eigenstates

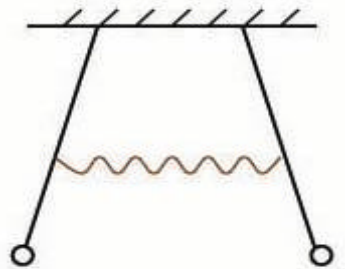
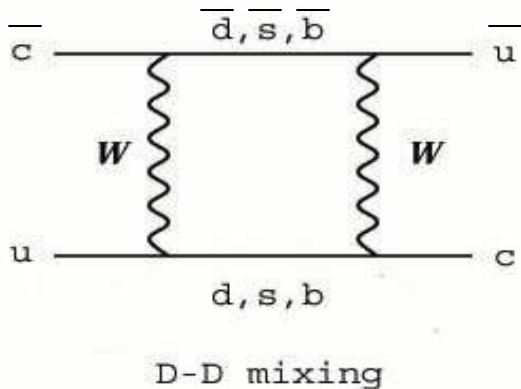
$$|B_{1,2}\rangle = p |B^0\rangle \pm q |\overline{B^0}\rangle$$

$$x = \frac{M_2 - M_1}{\Gamma}, \quad y = \frac{\Gamma_2 - \Gamma_1}{2\Gamma}$$



*Coupled oscillators*

# Introduction: why do we care?



## D-mixing:

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$$|D(t)\rangle = \begin{pmatrix} a(t) \\ b(t) \end{pmatrix} = a(t) |D^0\rangle + b(t) |\overline{D^0}\rangle$$

Time-dependence: coupled Schrodinger equations

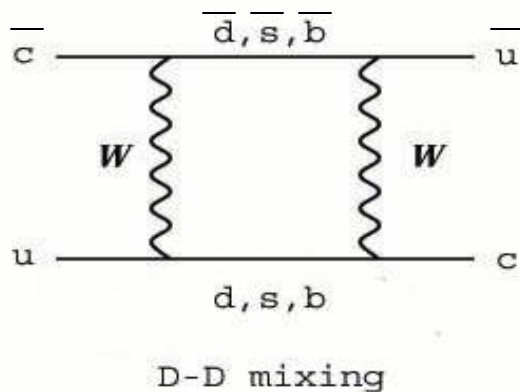
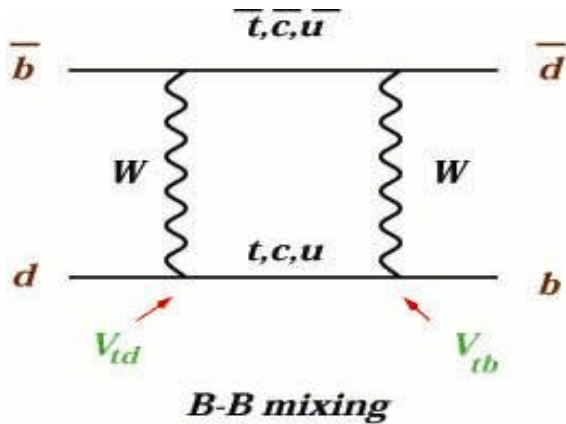
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## D-mixing:

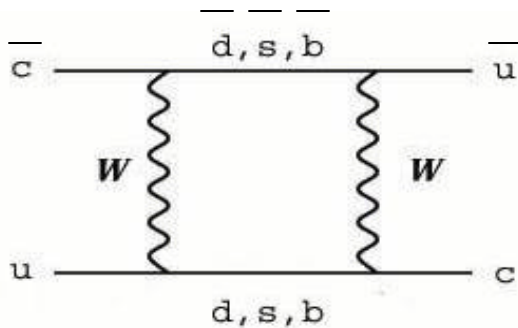
- ➔ the **only** probe of down-type quark dynamics
  - ➔ GIM mechanism: **no** ultra-heavy quarks in the loop
  - ➔ **b-quark** contribution  $\propto V_{ub} m_b^2$  can be neglected
- $$rate \propto f(m_s) - f(m_d) = 0 \quad (\text{SU}(3)_F \text{ limit})$$
- ➔ very sensitive to **long-distance QCD**, as  $m_c \sim 1 \text{ GeV}$

Clean probe of New Physics?

# How would new physics affect mixing?

D-D mass matrix:

Local operator  $\Rightarrow$  possible  
**New Physics!**



D-D mixing

$$\left( M - \frac{i}{2} \Gamma \right)_{ij} = m_D^{(0)} \mathbf{d}_{ij} + \frac{1}{2m_D} \langle D_i^0 | H_W^{\Delta C=2} | D_j^0 \rangle + \frac{1}{2m_D} \sum_I \frac{\langle D_i^0 | H_W^{\Delta C=1} | I \rangle \langle I | H_W^{\Delta C=1} | D_j^0 \rangle}{m_D^2 - m_I^2 + i\epsilon}$$

Real intermediate states, affect both x and y  $\Rightarrow$  **Standard Model**

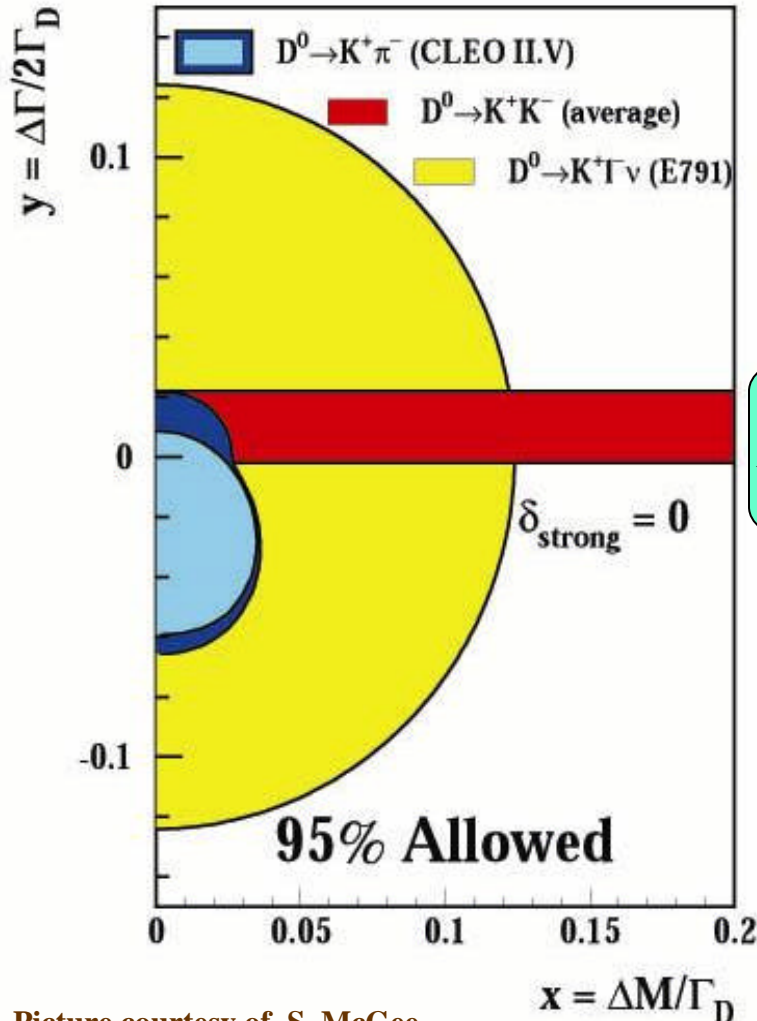
1.  $x \gg y$  : signal for New Physics?  
 $x \approx y$  : Standard Model?

2. CP violation in mixing/decay

With *b*-quark contribution neglected:  
only **2** generations contribute  
 $\Rightarrow$  **real 2x2 Cabibbo matrix**

# Experimental constraints

## $D^0$ - $\bar{D}^0$ Mixing Limits



Picture courtesy of S. McGee

### 1. Time-dependent $D^0(t) \rightarrow K^+ p^-$ analysis

$$\Gamma[D^0(t) \rightarrow K^+ p^-] = e^{-\Gamma t} |A_{K^+ p^-}|^2 \times \left[ R + \sqrt{R} R_m (y' \cos \mathbf{f} - x' \sin \mathbf{f}) \Gamma t + \frac{R_m^2}{4} (y^2 + x^2) (\Gamma t)^2 \right]$$

$$R = \left| \frac{A_{K^+ p^-}}{A_{K^+ K^-}} \right|^2$$

$$R_m e^{i\mathbf{f}} = \frac{q}{p}$$

if  $D^0 \leftrightarrow \bar{D}^0$ :

$$R_m \leftrightarrow R_m^{-1}$$

$$x' \leftrightarrow -x'$$

### 2. Lifetime difference analysis

$$y_{CP} = \frac{t(D \rightarrow p^+ K^-)}{t(D \rightarrow K^+ K^-)} - 1 = y \cos \mathbf{f} - x \sin \mathbf{f} \frac{A_m}{2}$$

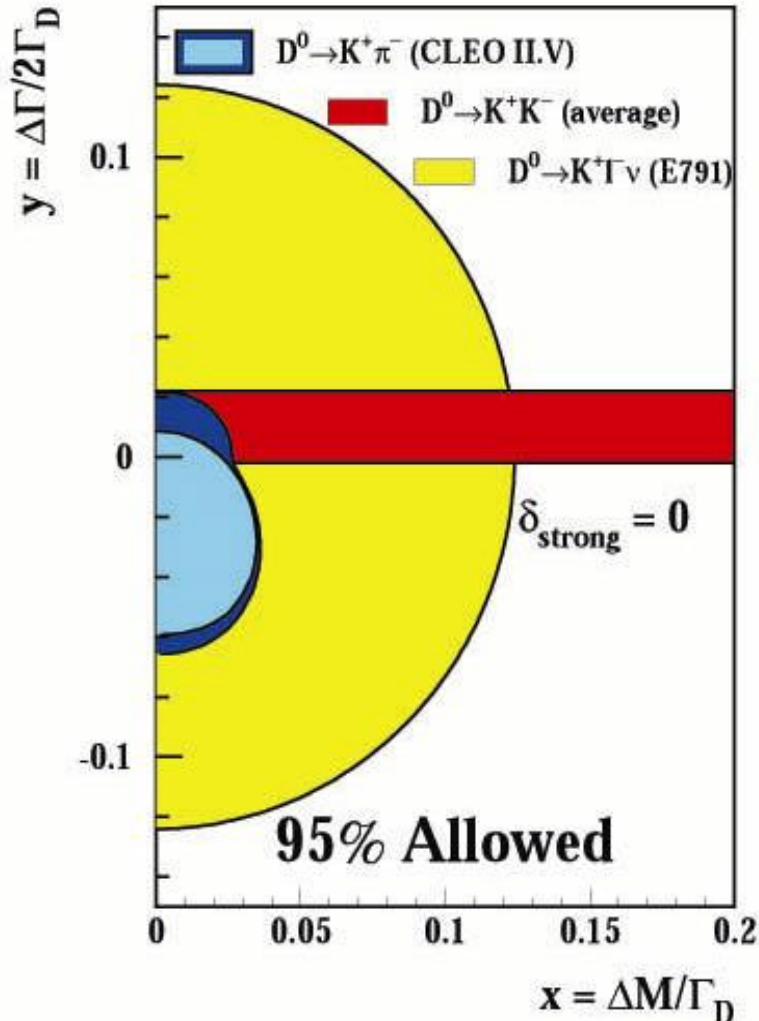
### 3. Semileptonic analysis

$$\text{rate} \propto x^2 + y^2$$

Quadratic in x,y: not so sensitive

# Experimental constraints

## $D^0$ - $\bar{D}^0$ Mixing Limits



Several groups have measured  $y_{\text{CP}}$

Experiment	Value
<b>FOCUS (2000)</b>	<b><math>(3.42 \pm 1.39 \pm 0.74)\%</math></b>
E791(2001)	$(0.8 \pm 2.9 \pm 1.0)\%$
CLEO (2002)	$(-1.2 \pm 2.5 \pm 1.4)\%$
Belle (2002)	$(-0.5 \pm 1.0^{+0.7}_{-0.8})\%$
BaBar (2002)	$(1.4 \pm 1.0^{+0.6}_{-0.7})\%$

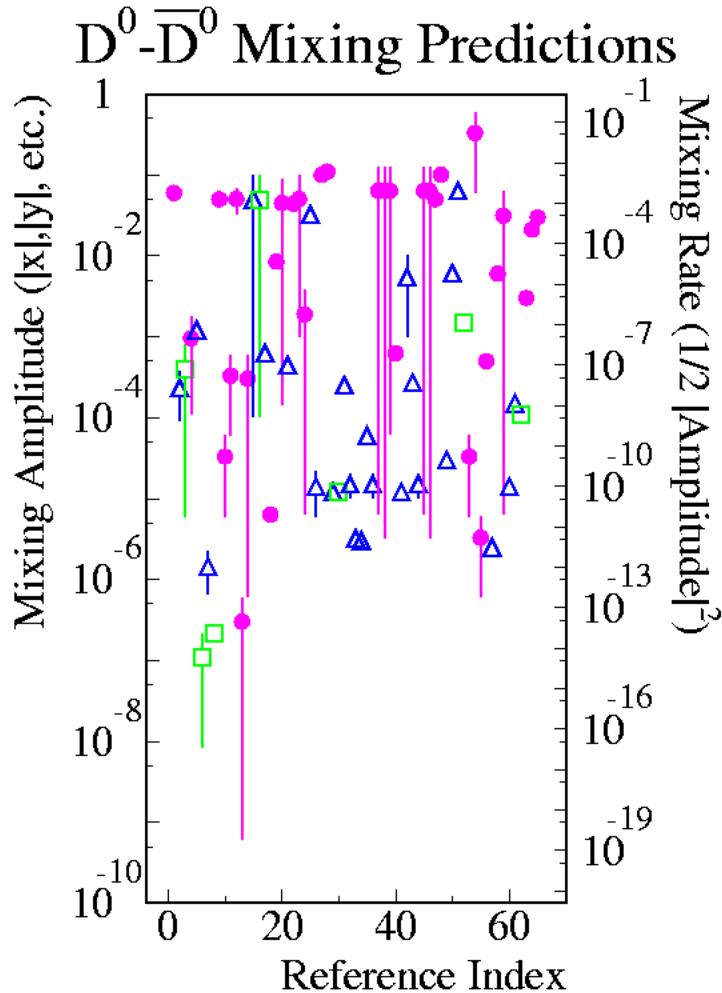
World average:  $(1.0 \pm 0.7)\%$

G. Raz

What are the expectations for x and y?



# Theoretical estimates

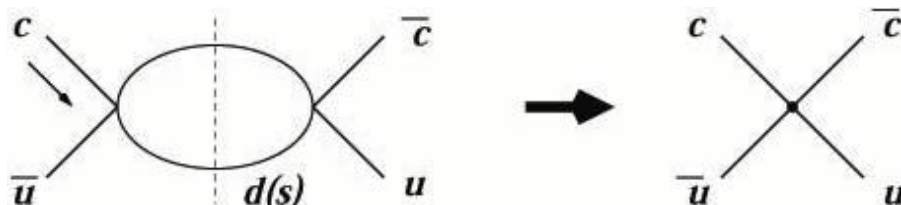


- Theoretical predictions are all over the board
- Can  $y \sim 1\%$  be convincingly accommodated?
- Is it possible to have  $y \gg x$ ?
- Does it still mean that  $y \sim x$ ?

# Theoretical estimates I

A. Short distance gives a tiny contribution, consider  $y$  as an example

$m_c$  is quite large !!!



$$y = \frac{1}{m_D \Gamma} \langle D^0 | ? | \bar{D}^0 \rangle$$

... as can be seen from the straightforward computation...



$$y_{sd} = \frac{N_C + 1}{2 \mathbf{p} N_C \Gamma} X_D \frac{(m_s^2 - m_d^2)^2}{m_c^2} \frac{m_s^2 + m_d^2}{m_c^2} [C_2^2 + 2C_1 C_2 + C_1^2 N_C - \frac{2(2N_C - 1) B_D'}{1 + N_C} \frac{M_D^2 C_2^2}{B_D (m_c + m_u)^2} \left( 1 + \left( N_C \frac{C_1^2}{C_2^2} + 2 \frac{C_1}{C_2} \right) \frac{2 - N_C}{2N_C - 1} \right)]$$

with  $\langle D^0 | \bar{u} \Gamma_{m_c} c \bar{u} \Gamma^{m_c} c | D^0 \rangle = \frac{1 + N_C}{N_C} \frac{4 F_D^2 m_D^2}{2 m_D} B_D, \text{ etc.}$

4 unknown matrix elements

similar for  $x$  (trust me!)

# Theoretical estimates I

A. Short distance + “subleading corrections” (in  $1/m_c$  expansion):

$$y_{sd}^{(6)} \propto \frac{(m_s^2 - m_d^2)^2}{m_c^2} \frac{m_s^2 + m_d^2}{m_c^2} \mathbf{m}_{had}^{-2} \propto m_s^6 \Lambda^{-6}$$

$$x_{sd}^{(6)} \propto \frac{(m_s^2 - m_d^2)^2}{m_c^2} \mathbf{m}_{had}^{-2} \propto m_s^4 \Lambda^{-4}$$

4 unknown matrix elements

...subleading effects?

$$y_{sd}^{(9)} \propto m_s^3 \Lambda^{-3}$$

$$x_{sd}^{(9)} \propto m_s^3 \Lambda^{-3}$$

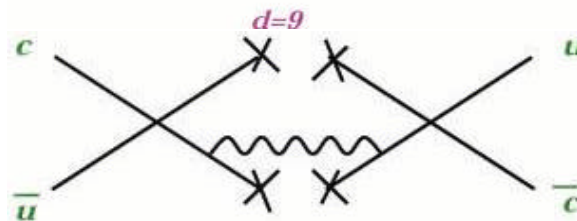


15 unknown matrix elements

Georgi, ...  
Bigi, Uraltsev

$$y_{sd}^{(12)} \propto \mathbf{b}_0 \mathbf{a}_s^2(\mathbf{m}) m_s^2 \Lambda^{-2}$$

$$x_{sd}^{(12)} \propto \mathbf{a}_s(\mathbf{m}) m_s^2 \Lambda^{-2}$$



Twenty-something unknown  
matrix elements

 Leading contribution!!!

Guestimate:  $x \sim y \sim 10^{-3}?$

**Resume:** model-independent computation  
with model-dependent result

# Theoretical estimates II

$m_c$  is NOT large !!!

B. Long distance might give a large result? Let's see...

$$y = \frac{1}{2\Gamma} \sum_n \mathbf{r}_n \left[ \langle D^0 | H_W^{\Delta C=1} | n \rangle \langle n | H_W^{\Delta C=1} | \bar{D}^0 \rangle + \langle \bar{D}^0 | H_W^{\Delta C=1} | n \rangle \langle n | H_W^{\Delta C=1} | D^0 \rangle \right]$$

... with  $n$  being all states to which  $D^0$  and  $\bar{D}^0$  can decay. Consider  $\pi\pi$ ,  $\pi K$ ,  $KK$  intermediate states as an example...

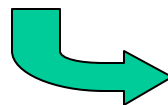
$$y_2 = Br(D^0 \rightarrow K^+ K^-) + Br(D^0 \rightarrow p^+ p^-) - 2 \cos \mathbf{d} \sqrt{Br(D^0 \rightarrow K^+ p^-) Br(D^0 \rightarrow p^+ K^-)}$$

cancellation expected!

If every Br is known up to  $O(1\%)$   $\rightarrow$  the result is expected to be  $O(1\%)!$

The result here is a series of large numbers with alternating signs, SU(3) forces 0

$x = ?$  Extremely hard...



need to restructure the calculation...

**Resume:** model-dependent computation  
with model-dependent result

## Questions:

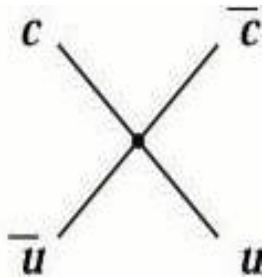
1. Can any model-independent statements be made for  $x$  or  $y$  ?

What is the order of SU(3) breaking?  
i.e. if  $x, y \propto m_s^n$  what is n?

2. Can one claim that  $y \sim 1\%$  is natural?

# Theoretical expectations

At which order in  $SU(3)_F$  breaking does the effect occur? Group theory?



$$\langle D^0 | H_w H_w | \bar{D}^0 \rangle \Rightarrow \langle 0 | D H_w H_w D | 0 \rangle$$

is a singlet with  $D \rightarrow D_i$  that belongs to  $\mathbf{3}$  of  $SU(3)_F$  (one light quark)

The  $\Delta C=1$  part of  $H_w$  is  $(\bar{q}_i c)(\bar{q}_j q_k)$ , i.e.  $3 \times \bar{3} \times \bar{3} = \bar{15} + 6 + \bar{3} + \bar{3} \Rightarrow H_k^{ij}$

$$O_{\bar{15}} = (\bar{s}d)(\bar{u}d) + (\bar{u}c)(\bar{s}d) + s_1(\bar{d}c)(\bar{u}d) + s_1(\bar{u}c)(\bar{d}d) \\ - s_1(\bar{s}c)(\bar{u}s) - s_1(\bar{u}c)(\bar{s}s) - s_1^2(\bar{d}c)(\bar{u}s) - s_1^2(\bar{u}c)(\bar{d}s)$$



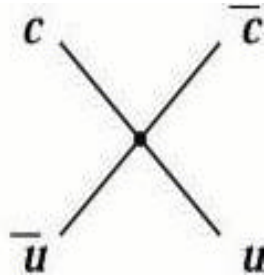
$$O_6 = (\bar{s}d)(\bar{u}d) - (\bar{u}c)(\bar{s}d) + s_1(\bar{d}c)(\bar{u}d) - s_1(\bar{u}c)(\bar{d}d) \\ - s_1(\bar{s}c)(\bar{u}s) + s_1(\bar{u}c)(\bar{s}s) - s_1^2(\bar{d}c)(\bar{u}s) + s_1^2(\bar{u}c)(\bar{d}s)$$

Introduce  $SU(3)$  breaking via the quark mass operator  $M_j^i = \text{diag}(m_u, m_d, m_s)$

All nonzero matrix elements built of  $D_i, H_k^{ij}, M_j^i$  must be  $SU(3)$  singlets



# Theoretical expectations



note that  $D_i D_j$  is symmetric  $\Rightarrow$  belongs to  $\mathbf{6}$  of  $SU(3)_F$

$$\langle D^0 | H_W H_W | \bar{D}^0 \rangle \Rightarrow \langle 0 | D H_W H_W D | 0 \rangle$$

Explicitly,

$$DD \Rightarrow D_6$$

$$H_W H_W \Rightarrow O_{\bar{6}0} + O_{42} + O_{15'}$$

1. No  $\bar{6}$  in the decomposition of  $H_W H_W \Rightarrow$  **no**  $SU(3)$  singlet can be formed

$\Rightarrow$  **D mixing is prohibited by  $SU(3)$  symmetry**

2. Consider a single insertion of  $M_j^i \Rightarrow D_6 M$  transforms as  $6 \times 8 = 24 + \bar{15} + 6 + \bar{3} \Rightarrow$  still **no**  $SU(3)$  singlet can be formed

$\Rightarrow$  **NO D mixing at first order in  $SU(3)$  breaking**

3. Consider double insertion of  $M \Rightarrow DDM : 6 \times (8 \times 8)_S = (60 + \bar{42}) + 24 + \bar{15} + \bar{15} + 6 + (24 + 15 + 6 + \bar{3}) + 6$

$\Rightarrow$  **D mixing occurs only at the second order in  $SU(3)$  breaking**

A.F., Y.G., Z.L., and A.A.P.

# Theoretical expectations

- Does it always work? SU(3) breaking must enter perturbatively...

$$A_i = A_{SU(3)} + \mathbf{d}_i$$

- Known counter-examples:

1. Very narrow light quark resonance with  $m_R \sim m_D$

$$x, y \sim \frac{g_{DR}^2}{m_D^2 - m_R^2} \sim \frac{g_{DR}^2}{m_D^2 - m_0^2 - 2m_0 \mathbf{d}_R}$$

Most probably don't exist...

see E.Golowich and A.A.P.

2. Part of the multiplet is kinematically forbidden

Example: both  $D^0 \rightarrow 4p$  and  $D^0 \rightarrow 4K$  are from the same multiplet, but the latter is kinematically forbidden

see A.F., Y.G., Z.L., and A.A.P.

# Theoretical expectations

- Two major sources of SU(3) breaking

1. phase space

$$m_K \neq m_p \neq m_h \dots$$

2a. matrix elements (absolute value)

$$f_K \neq f_p \dots$$

2b. matrix elements (phases aka FSI)

$$\text{Im} \frac{A(D^0 \rightarrow K^+ p^-)}{A(\bar{D}^0 \rightarrow K^+ p^-)} \neq 0$$

Take into account only the first source (computable)!

# SU(3) and phase space

- “Repackage” the analysis: look at the complete multiplet contribution

$$y = \sum_{F_R} y_{F,R} Br(D^0 \rightarrow F_R) \sim \sum_{F_R} y_{F,R} \frac{1}{\Gamma} \sum_{n \in F_R} \Gamma(D^0 \rightarrow n)$$

y for each SU(3) multiplet

Each is **0** in SU(3)

- Does it help? If only phase space is taken into account: **no (mild)** model dependence

$$y_{F,R} = \frac{\sum_{n \in F_R} \langle \bar{D}^0 | H_W | n \rangle \mathbf{r}_n \langle n | H_W | D^0 \rangle}{\sum_{n \in F_R} \langle D^0 | H_W | n \rangle \mathbf{r}_n \langle n | H_W | D^0 \rangle}$$

$$= \frac{\sum_{n \in F_R} \langle \bar{D}^0 | H_W | n \rangle \mathbf{r}_n \langle n | H_W | D^0 \rangle}{\sum_{n \in F_R} \Gamma(D^0 \rightarrow n)}$$

if CP is conserved

Can consistently compute !

# Results

Final state representation		$y_{F,R}/s_1^2$	$y_{F,R}$ (%)
$PP$	8	-0.0038	-0.018
	27	-0.00071	-0.0034
$PV$	$8_S$	0.031	0.15
	$8_A$	0.032	0.15
	10	0.020	0.10
	$\overline{10}$	0.016	0.08
	27	0.040	0.19
$(VV)_{s\text{-wave}}$	8	-0.081	-0.39
	27	-0.061	-0.30
$(VV)_{p\text{-wave}}$	8	-0.10	-0.48
	27	-0.14	-0.70
$(VV)_{d\text{-wave}}$	8	0.51	2.5
	27	0.57	2.8

Final state representation		$y_{F,R}/s_1^2$	$y_{F,R}$ (%)
$(3P)_{s\text{-wave}}$	8	-0.48	-2.3
	27	-0.11	-0.54
$(3P)_{p\text{-wave}}$	8	-1.13	-5.5
	27	-0.07	-0.36
$(3P)_{\text{form-factor}}$	8	-0.44	-2.1
	27	-0.13	-0.64
$4P$	8	3.3	16
	27	2.2	9.2
	27'	1.9	11

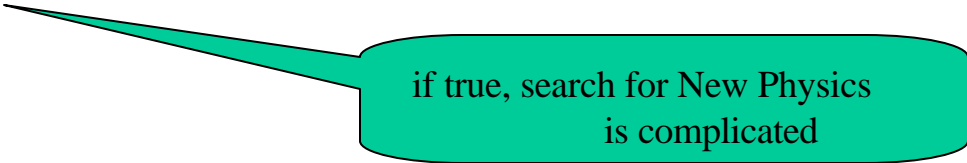
- Product is naturally  $O(1\%)$
- No (symmetry-enforced) cancellations
- Does NOT occur for  $x$

naturally implies that  $y \sim 1\%$  and  $x < y$  !

Final state	fraction
$PP$	5%
$PV$	10%
$(VV)_{s\text{-wave}}$	5%
$(VV)_{d\text{-wave}}$	5%
$3P$	5%
$4P$	10%

# Conclusions

- $x, y=0$  in the SU(3) limit (as  $V_{ub}$  is very small)
- it is a **second** order effect
- it is quite possible that  $y \sim 1\%$  with  $x < y$



if true, search for New Physics  
is complicated

- expect new data from BaBar/Belle/CLEO/CLEOc/CDF(?)
- currently:  $x \approx (\pm 2.8 \pm 2.5)\%$ ,  $y \approx (\pm 0.9 \pm 3.6)\%$  (allowing NP)
- CP-violation in mixing is a “smoking gun” signal for New Physics