

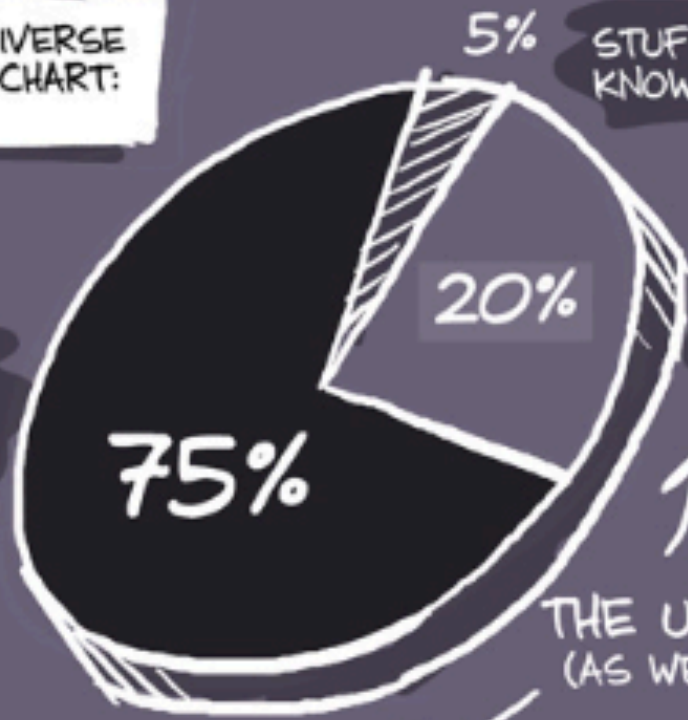
# Exploring the Unknown Universe

Daniel Whiteson, UC Irvine



IF YOU LOOK AT WHAT THE UNIVERSE IS MADE OUT OF, LIKE A PIE CHART:

WE HAVE  
**NO IDEA**



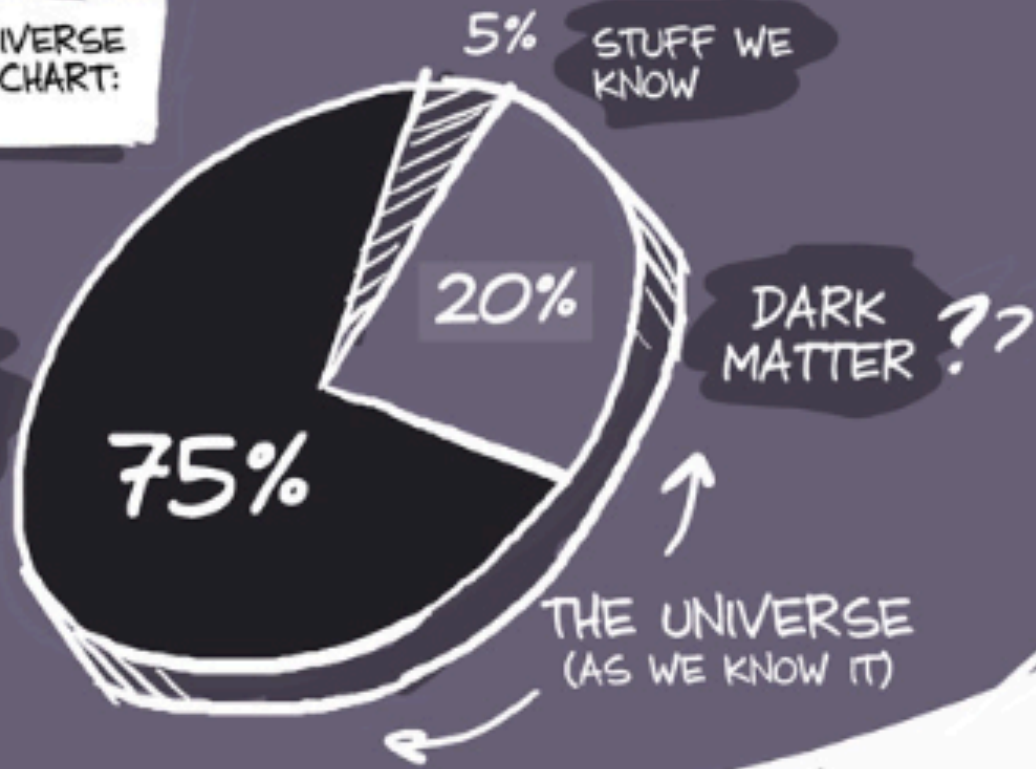
5% STUFF WE KNOW

DARK MATTER ??

THE UNIVERSE  
(AS WE KNOW IT)

IF YOU LOOK AT WHAT THE UNIVERSE IS MADE OUT OF, LIKE A PIE CHART:

WE HAVE  
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THE UNIVERSE  
(AS WE KNOW IT)

DARK MATTER IS 5 TIMES AS HEAVY AS ALL THE MATTER WE KNOW ABOUT,

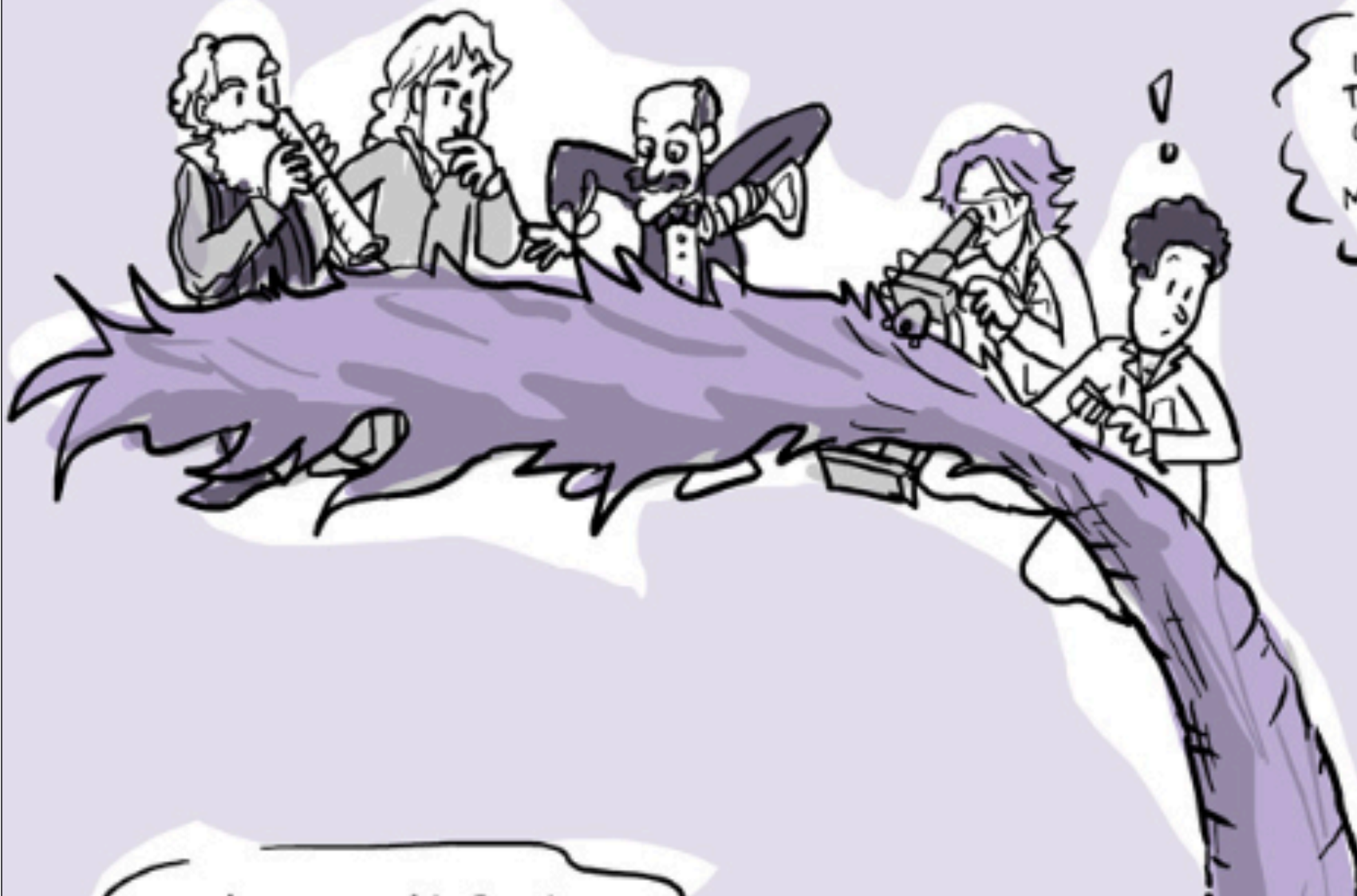


WE'VE BEEN  
STUDYING MATTER  
FOR A COUPLE  
HUNDRED YEARS...

...AND WE HAVE A  
FINE UNDERSTANDING  
OF CHEMICALS, ETC...

AND ALL OF A SUDDEN WE  
DISCOVER THAT ALL THAT  
WORK WE'VE BEEN DOING,

IS ONLY ON A  
TINY FRACTION  
OF WHAT THE  
UNIVERSE IS  
MADE OUT OF!



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IT'S LIKE YOU'VE BEEN  
STUDYING AN ELEPHANT'S  
TAIL FOR TWO HUNDRED  
YEARS AND YOU DISCOVER...

IT'S ONLY THE TAIL!



# Motivation

## The Standard Model

$$\begin{aligned} & -\frac{1}{2}\partial_\nu g_\mu^a \partial_\nu g_\mu^a - g_s f^{abc} \partial_\mu g_\nu^a g_\nu^b g_\nu^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_\mu^b g_\nu^c g_\mu^d g_\nu^e + \\ & \frac{1}{2}ig_s^2 (\bar{q}_i^c \gamma^\mu q_i^c) g_\mu^a + G^a \partial^2 G^a + g_s f^{abc} \partial_\mu G^a G^b G^c - \partial_\mu W_\nu^+ \partial_\nu W_\mu^- - \\ & M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2}M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\nu A_\mu - \frac{1}{2}\partial_\mu H \partial_\mu H - \\ & \frac{1}{2}m_H^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2\epsilon_0^2} M \phi^0 \phi^0 - \beta_h \left[ \frac{2M^2}{g^2} + \right. \\ & \left. \frac{2M}{g} H + \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right] + \frac{2M^4}{g^4} \alpha_h - ig_{c_w} [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\ & W_\mu^- W_\nu^+) - Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\nu^+) + Z_\mu^0 (W_\nu^- \partial_\nu W_\mu^+ - \\ & W_\mu^+ \partial_\nu W_\nu^-)] - ig_{s_w} [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\mu^- W_\nu^+) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - \\ & W_\mu^- \partial_\nu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)] - \frac{1}{2}g^2 W_\mu^+ W_\mu^- W_\nu^+ W_\nu^- + \\ & \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\nu^+ W_\mu^- + g^2 c_w^2 (Z_\mu^0 W_\nu^+ Z_\mu^0 W_\nu^- - Z_\mu^0 W_\nu^+ W_\nu^-) + \\ & g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\nu W_\mu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\mu^0 (W_\mu^+ W_\nu^- - \\ & W_\mu^- W_\nu^+) - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-] - g\alpha [H^3 + H\phi^0 \phi^0 + 2H\phi^+ \phi^-] - \\ & \frac{1}{8}g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2] - \\ & gM W_\mu^+ W_\mu^- H - \frac{1}{2}g \frac{M}{c_w} Z_\mu^0 Z_\mu^0 H - \frac{1}{2}ig [W_\mu^+ (\partial_\nu \phi^0 \phi^- - \phi^- \partial_\nu \phi^0) - \\ & W_\mu^- (\partial_\nu \phi^0 \phi^+ - \phi^+ \partial_\nu \phi^0)] + \frac{1}{2}ig [W_\mu^+ (H\partial_\nu \phi^- - \phi^- \partial_\nu H) - W_\mu^- (H\partial_\nu \phi^+ - \\ & \phi^+ \partial_\nu H)] + \frac{1}{2}g \frac{1}{c_w} (Z_\mu^0 (H\partial_\nu \phi^0 - \phi^0 \partial_\nu H) - ig \frac{s_w^2}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \\ & ig_{s_w} M A_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - ig \frac{1-2c_w^2}{2c_w} Z_\mu^0 (\phi^+ \partial_\nu \phi^- - \phi^- \partial_\nu \phi^+) + \\ & ig_{s_w} A_\mu (\phi^+ \partial_\nu \phi^- - \phi^- \partial_\nu \phi^+) - \frac{1}{4}g^2 W_\mu^+ W_\mu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\ & \frac{1}{4}g^2 \frac{1}{c_w^2} Z_\mu^0 Z_\mu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-] - \frac{1}{2}g^2 \frac{s_w^2}{c_w} Z_\mu^0 \phi^0 (W_\mu^+ \phi^- + \\ & W_\mu^- \phi^+) - \frac{1}{2}ig^2 \frac{s_w^2}{c_w} Z_\mu^0 H (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + \\ & W_\mu^- \phi^+) + \frac{1}{2}ig^2 s_w A_\mu H (W_\mu^+ \phi^- - W_\mu^- \phi^+) - g^2 \frac{s_w}{c_w} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - \\ & g^1 s_w^2 A_\mu A_\nu \phi^+ \phi^- - e^\lambda (\gamma \partial + m_\nu^2) e^\lambda - \bar{\nu}^\lambda \gamma \partial \nu^\lambda - \bar{u}_i^\lambda (\gamma \partial + m_u^2) u_i^\lambda - \\ & \bar{d}_i^\lambda (\gamma \partial + m_d^2) d_i^\lambda + ig_{s_w} A_\mu [-(e^\lambda \gamma^\mu e^\lambda) + \frac{2}{3}(\bar{u}_i^\lambda \gamma^\mu u_i^\lambda) - \frac{1}{3}(\bar{d}_i^\lambda \gamma^\mu d_i^\lambda)] + \\ & \frac{ig}{4c_w} Z_\mu^0 [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (e^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (\bar{u}_i^\lambda \gamma^\mu (\frac{2}{3}s_w^2 - \\ & 1 - \gamma^5) u_i^\lambda) + (\bar{d}_i^\lambda \gamma^\mu (1 - \frac{2}{3}s_w^2 - \gamma^5) d_i^\lambda)] + \frac{ig}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) e^\lambda) + \\ & (\bar{u}_i^\lambda \gamma^\mu (1 + \gamma^5) C_{\lambda\nu} d_i^\nu)] + \frac{ig}{2\sqrt{2}} W_\mu^- [(e^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_i^\lambda C_{\lambda\nu}^\dagger \gamma^\mu (1 + \\ & \gamma^5) u_i^\lambda)] + \frac{ig}{2\sqrt{2}} \frac{m_\nu^2}{M} [-\phi^+ (\bar{\nu}^\lambda (1 - \gamma^5) e^\lambda) + \phi^- (e^\lambda (1 + \gamma^5) \nu^\lambda)] - \\ & \frac{g}{2} \frac{m_\nu^2}{M} [H (e^\lambda e^\lambda) + i\phi^0 (e^\lambda \gamma^5 e^\lambda)] + \frac{ig}{2M\sqrt{2}} \phi^+ [-m_\nu^2 (\bar{u}_i^\lambda C_{\lambda\nu} (1 - \gamma^5) d_i^\nu) + \\ & m_\nu^2 (\bar{u}_i^\lambda C_{\lambda\nu} (1 + \gamma^5) d_i^\nu)] + \frac{ig}{2M\sqrt{2}} \phi^- [m_\nu^2 (\bar{d}_i^\lambda C_{\lambda\nu}^\dagger (1 + \gamma^5) u_i^\nu) - m_\nu^2 (\bar{d}_i^\lambda C_{\lambda\nu}^\dagger (1 - \\ & \gamma^5) u_i^\nu)] - \frac{g}{2} \frac{m_\nu^2}{M} H (\bar{u}_i^\lambda u_i^\lambda) - \frac{g}{2} \frac{m_\nu^2}{M} H (\bar{d}_i^\lambda d_i^\lambda) + \frac{ig}{2} \frac{m_\nu^2}{M} \phi^0 (\bar{u}_i^\lambda \gamma^5 u_i^\lambda) - \\ & \frac{ig}{2} \frac{m_\nu^2}{M} \phi^0 (\bar{d}_i^\lambda \gamma^5 d_i^\lambda) + \bar{X}^+ (\partial^2 - M^2) X^+ + \bar{X}^- (\partial^2 - M^2) X^- + \bar{X}^0 (\partial^2 - \\ & \frac{M^2}{c_w^2}) X^0 + Y \partial^2 Y + ig_{c_w} W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \partial_\mu \bar{X}^+ X^0) + ig_{s_w} W_\mu^+ (\partial_\nu \bar{Y} X^- - \\ & \partial_\mu \bar{X}^+ Y) + ig_{c_w} W_\mu^- (\partial_\mu \bar{X}^- X^0 - \partial_\mu \bar{X}^0 X^+) + ig_{s_w} W_\mu^- (\partial_\mu \bar{X}^- Y - \\ & \partial_\mu \bar{Y} X^+) + ig_{c_w} Z_\mu^0 (\partial_\mu \bar{X}^+ X^+ - \partial_\mu \bar{X}^- X^-) + ig_{s_w} A_\mu (\partial_\mu \bar{X}^+ X^+ - \\ & \partial_\mu \bar{X}^- X^-) - \frac{1}{2}igM [\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w} \bar{X}^0 X^0 H] + \\ & \frac{1-2c_w^2}{2c_w} igM [\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} igM [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \\ & igMs_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \frac{1}{2}igM [\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0] \end{aligned}$$

Can this be right?

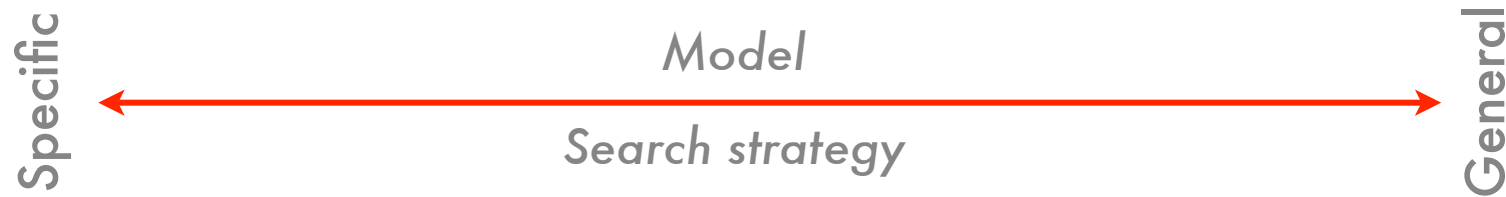
# Outline

I. Motivation

**II. Strategy**

III. Results

# Searching for new physics

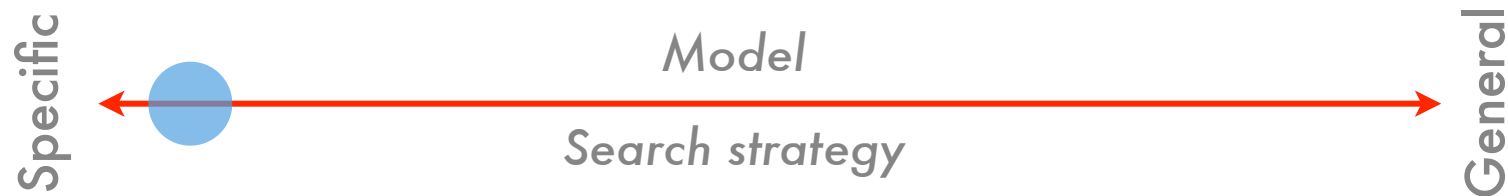


## Our goals:

- Maximize possibility for discovery
- Learn something no matter what we see

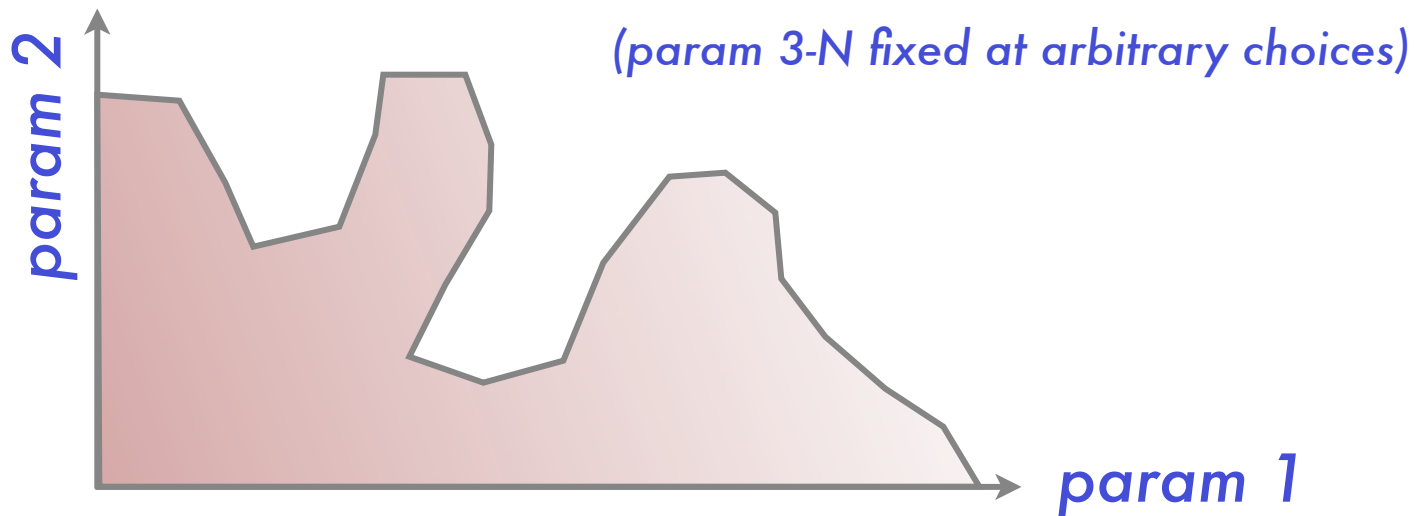


# Traditional approach

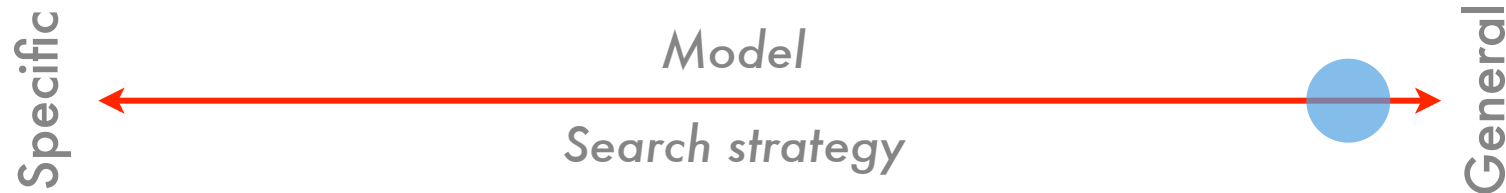


## Bet on a specific full theory

Optimize analysis to squeeze out maximal sensitivity to new physics.

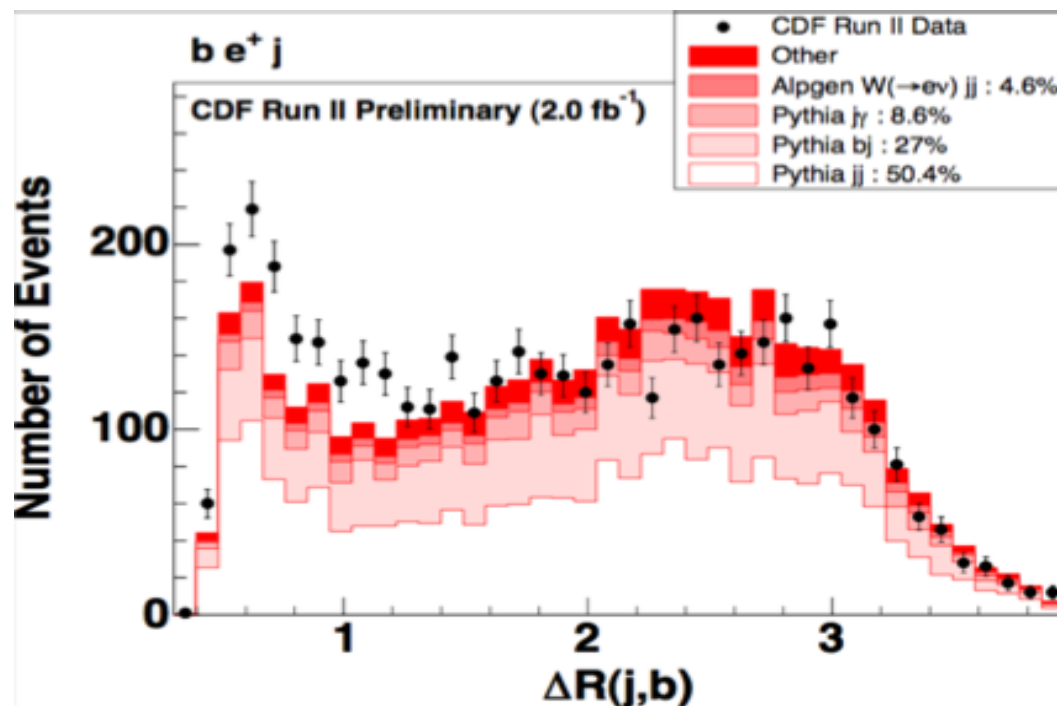


# Model independent search



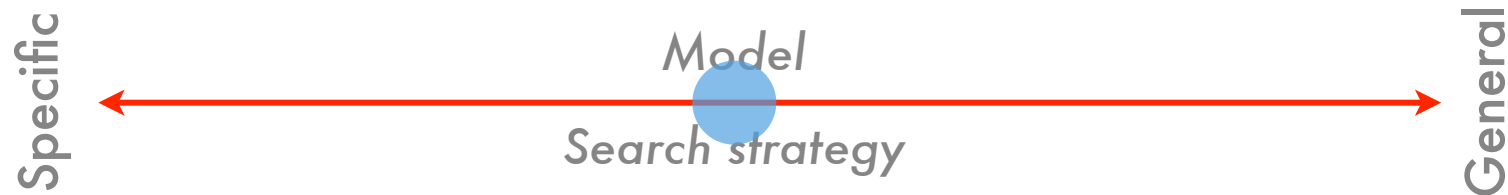
## Discard the model

compare data to standard model



“Never listen to theorists.  
Just go look for it”  
-Aaron Pierce, Theorist

# Compromise



## Admit the need for a model

New signal requires a coherent physical explanation,  
even trivial or effective

## Generalize your model

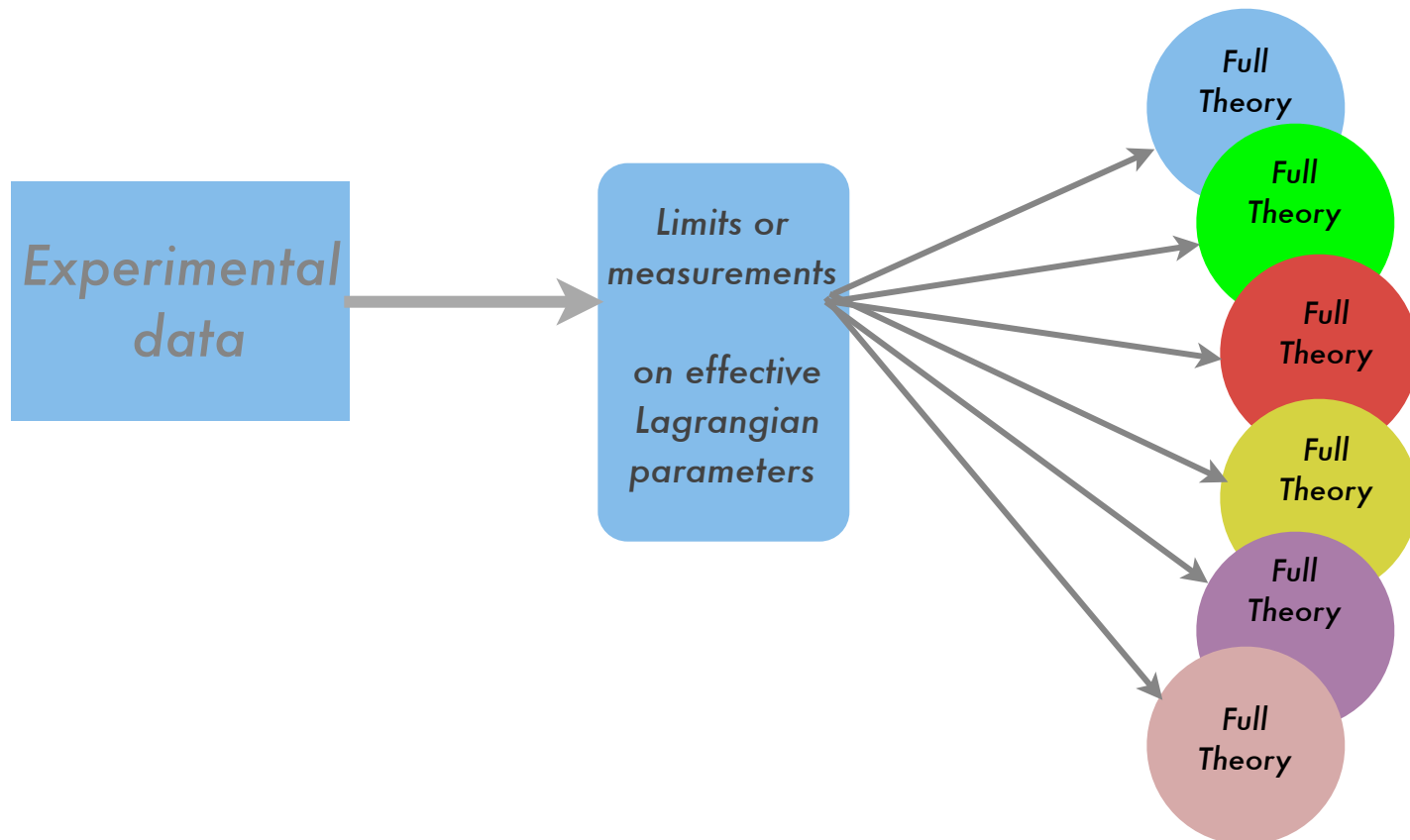
Focus on the general experimental sensitivity  
Construct simple models that describe classes of new physics

## Examples

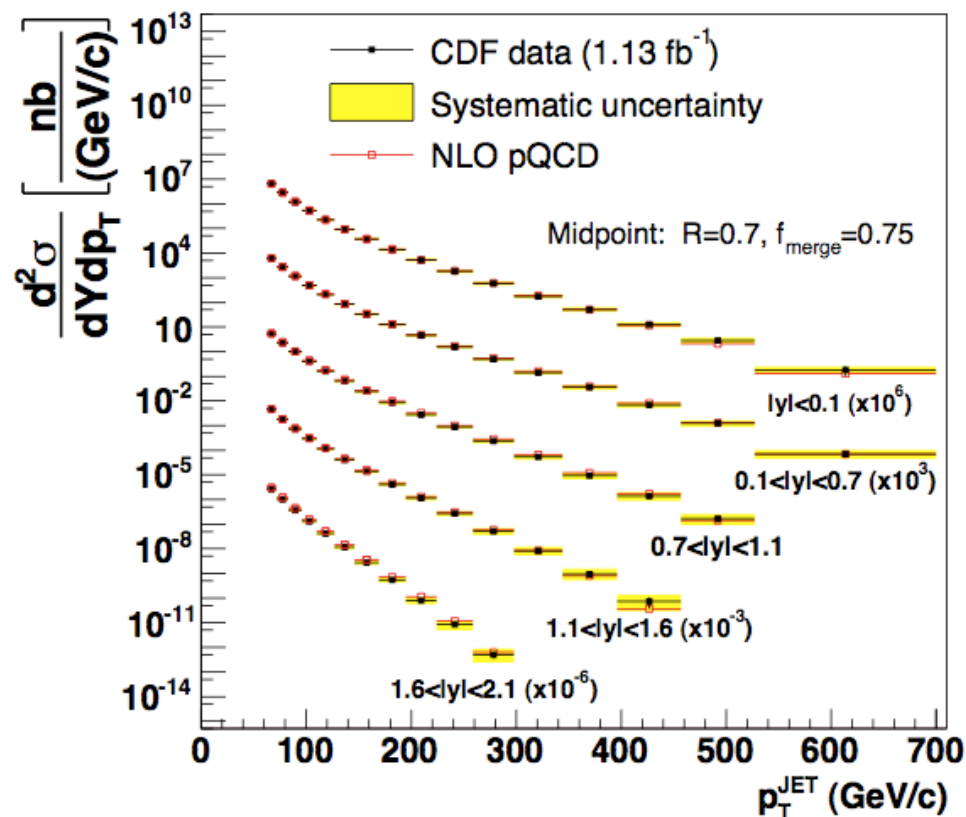
Simple SM extensions: fourth generation,  $Z'$ , resonances ( $X \rightarrow t\bar{t}$ ) etc

# Effective Lagrangian

A natural, compact language for communication between theory and experiment.



# A Theorist's dream?



## Unfolded cross-sections

Deconvolution to remove detector effects

Publish measured differential cross-sections

Theorists don't need to know/have detector description

*This is hard!*

# Outline

I. Motivation

II. Strategy

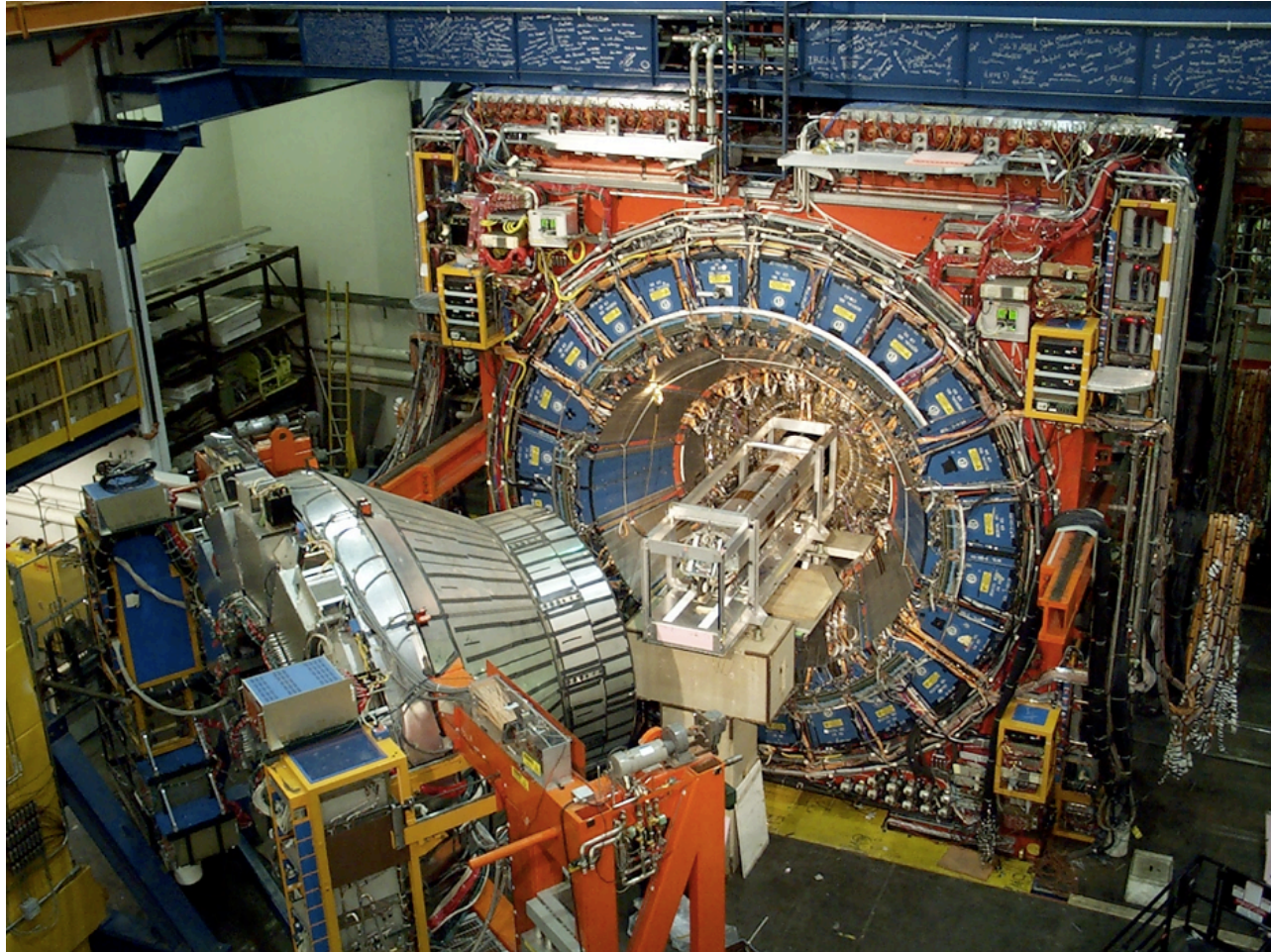
**III. Results**

*a. Heavy resonances ( $Z'$ )*

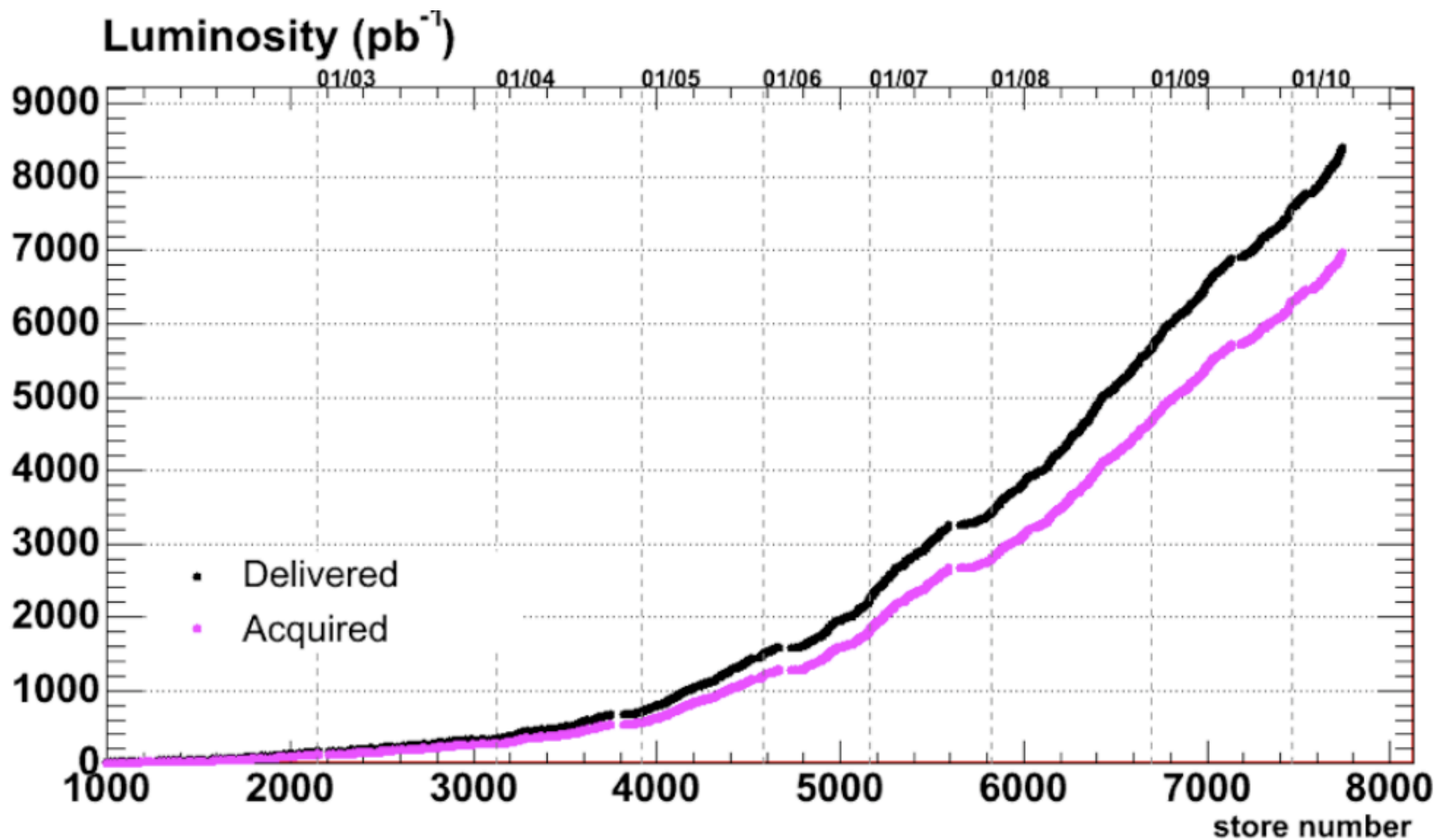
*b. Heavy quarks ( $b'$ ,  $t'$ )*

*c. Simplified SUSY*

# CDF



# Dataset

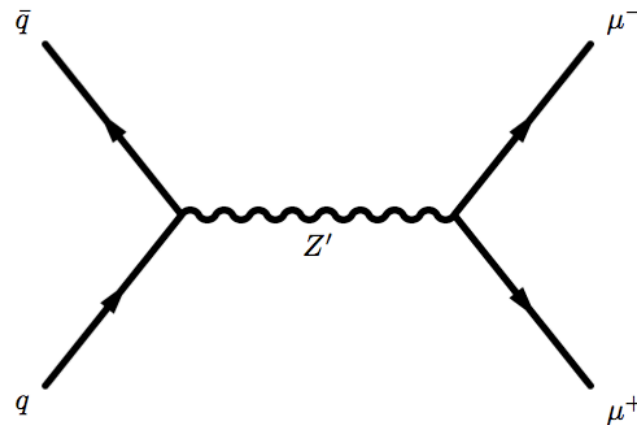




# High mass resonances

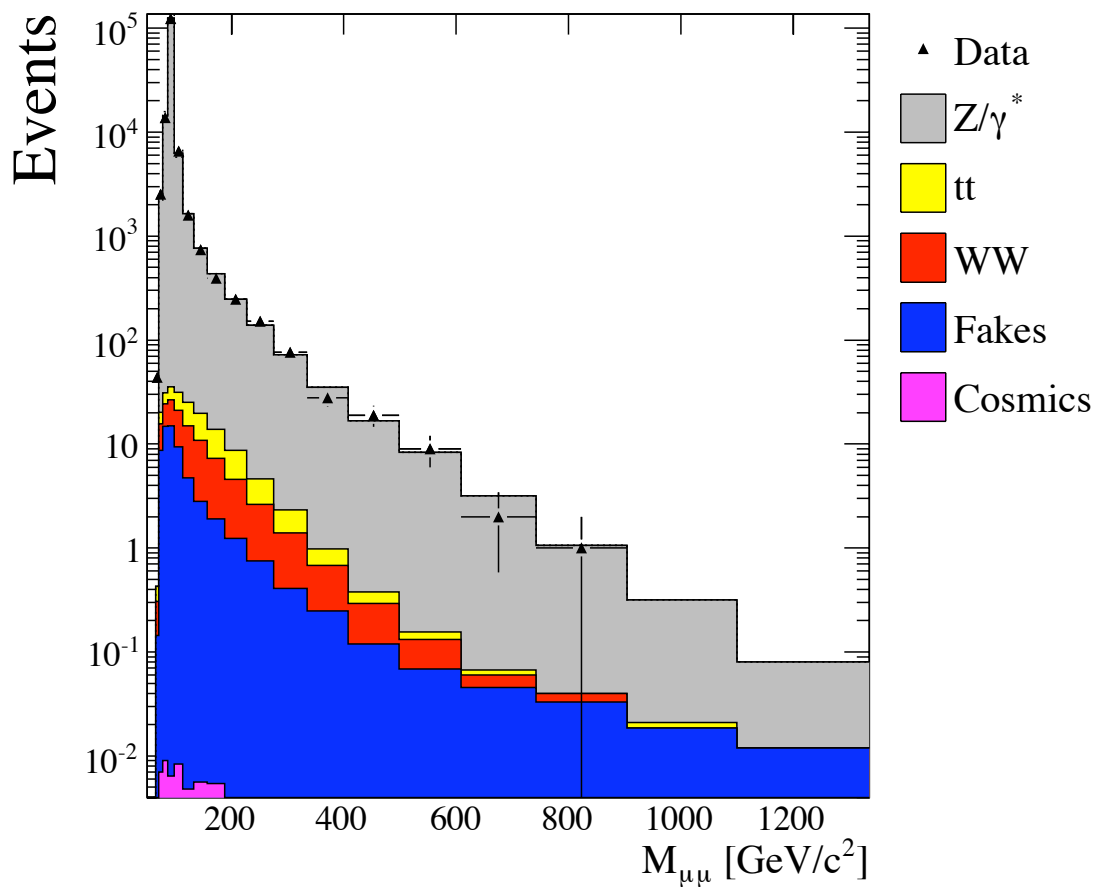
UCI Undergrad  
Eddie Quinlan

$Z'$  to di-muons

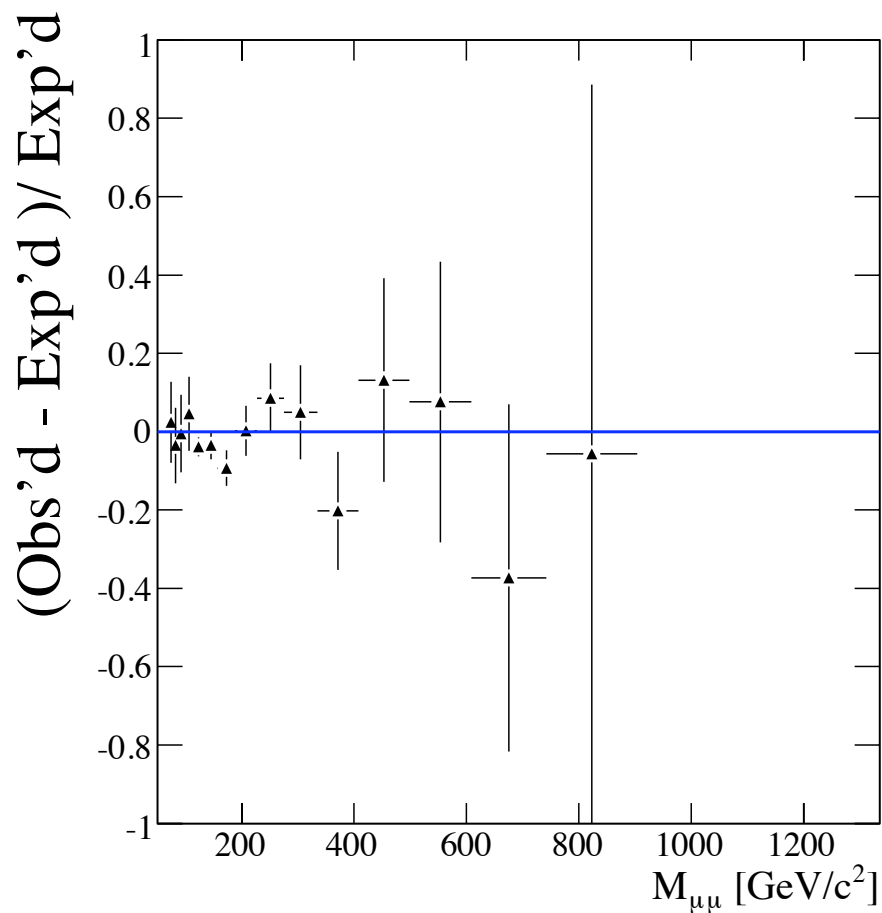


# High mass dimuon res.

CDF Run II Preliminary 4.6 fb<sup>-1</sup>



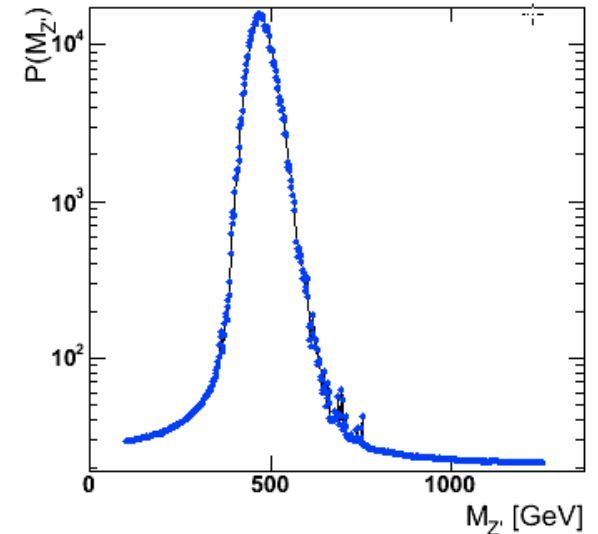
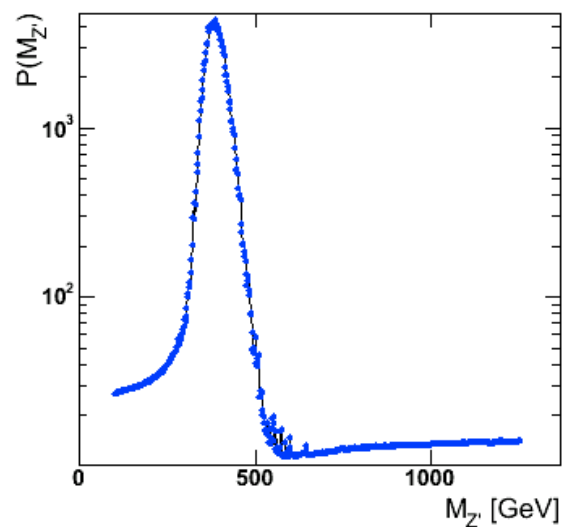
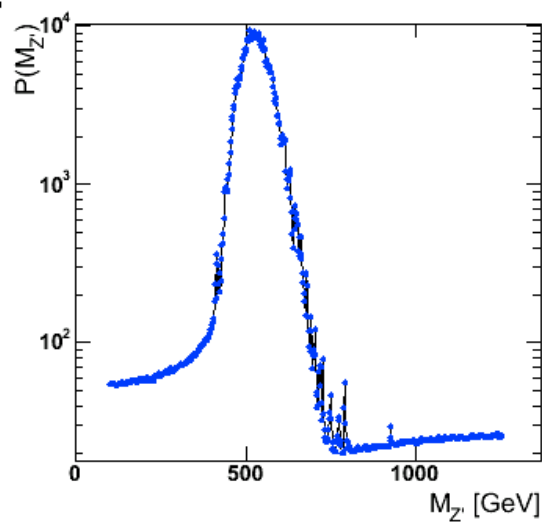
CDF Run II Preliminary 4.6 fb<sup>-1</sup>



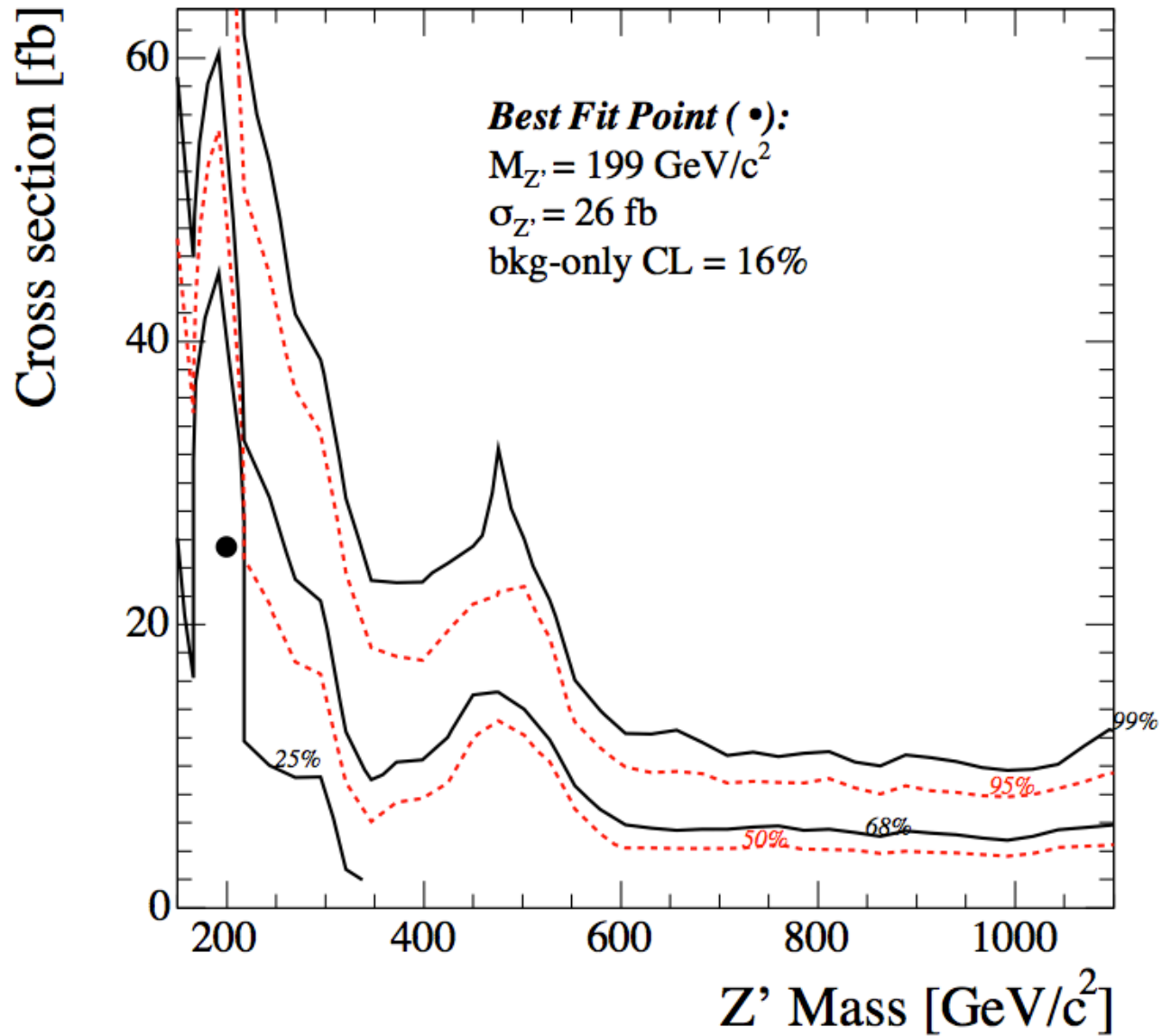
PRL 2011, to appear

# Z' to muons

$$P_{Z'}(x_i|M_{Z'}) = \int dq_1 dq_2 |\mathcal{M}_{Z'}(M_{Z'})|^2 \\ \times f_{PDF}(x_p) f_{PDF}(x_{\bar{p}}) T(p_1, q_1) T(p_2, q_2) P_{PT}(q_1 + q_2, N_{jets})$$

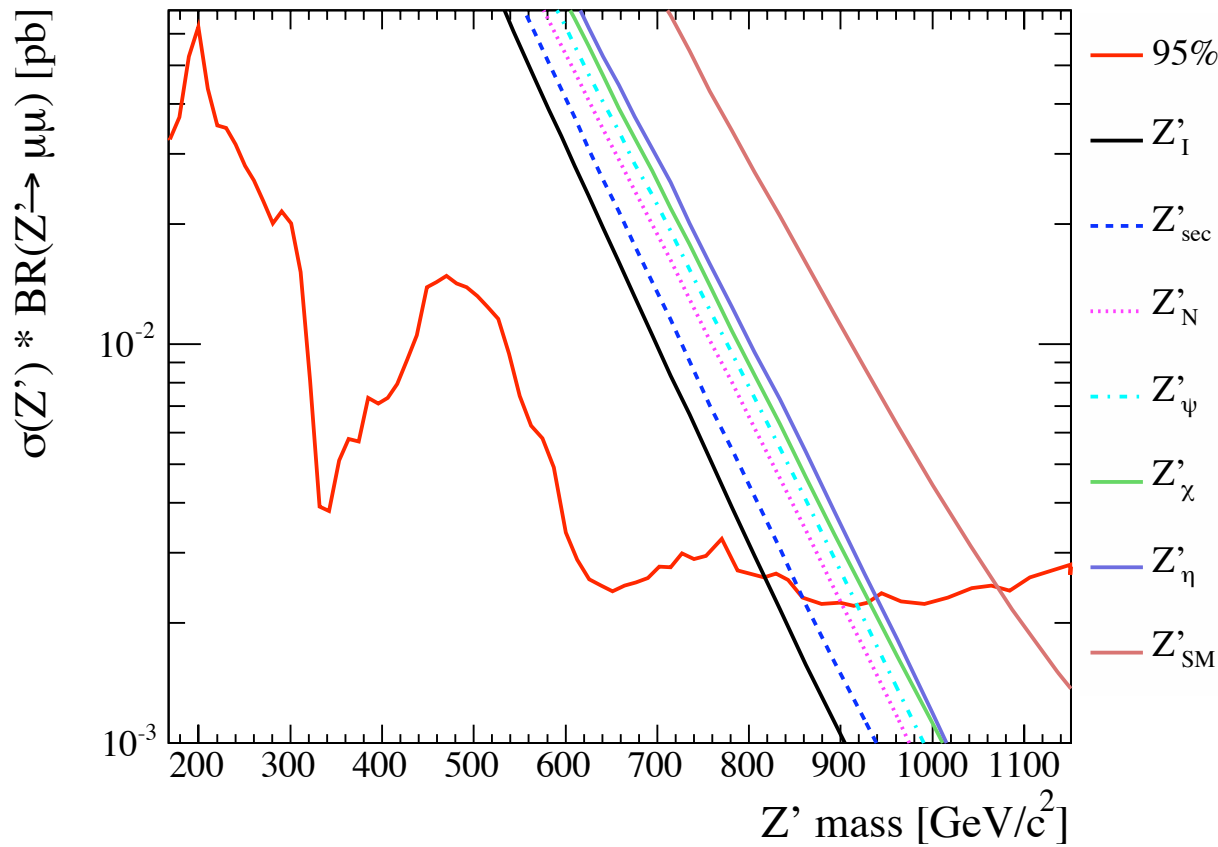


# Z' to muons



# Z' to muons

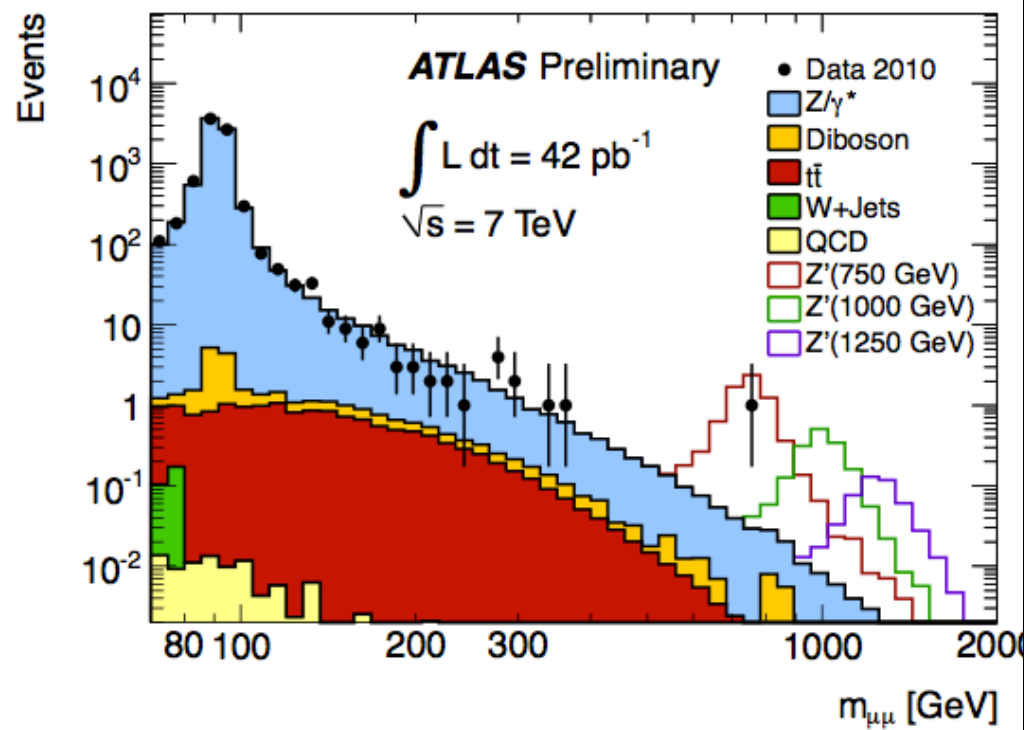
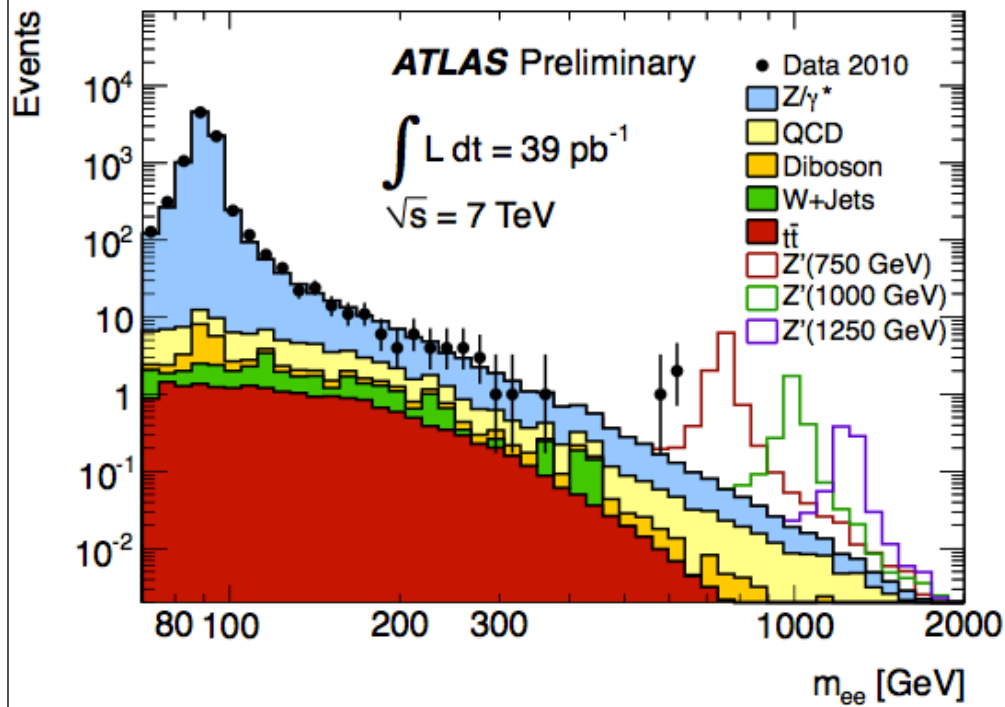
CDF Run II Preliminary 4.6 fb<sup>-1</sup>



Model	Mass Limit (GeV/c <sup>2</sup> )
Z' <sub>1</sub>	817
Z' <sub>sec</sub>	858
Z' <sub>N</sub>	900
Z' <sub>ψ</sub>	917
Z' <sub>χ</sub>	930
Z' <sub>η</sub>	938
Z' <sub>SM</sub>	1071

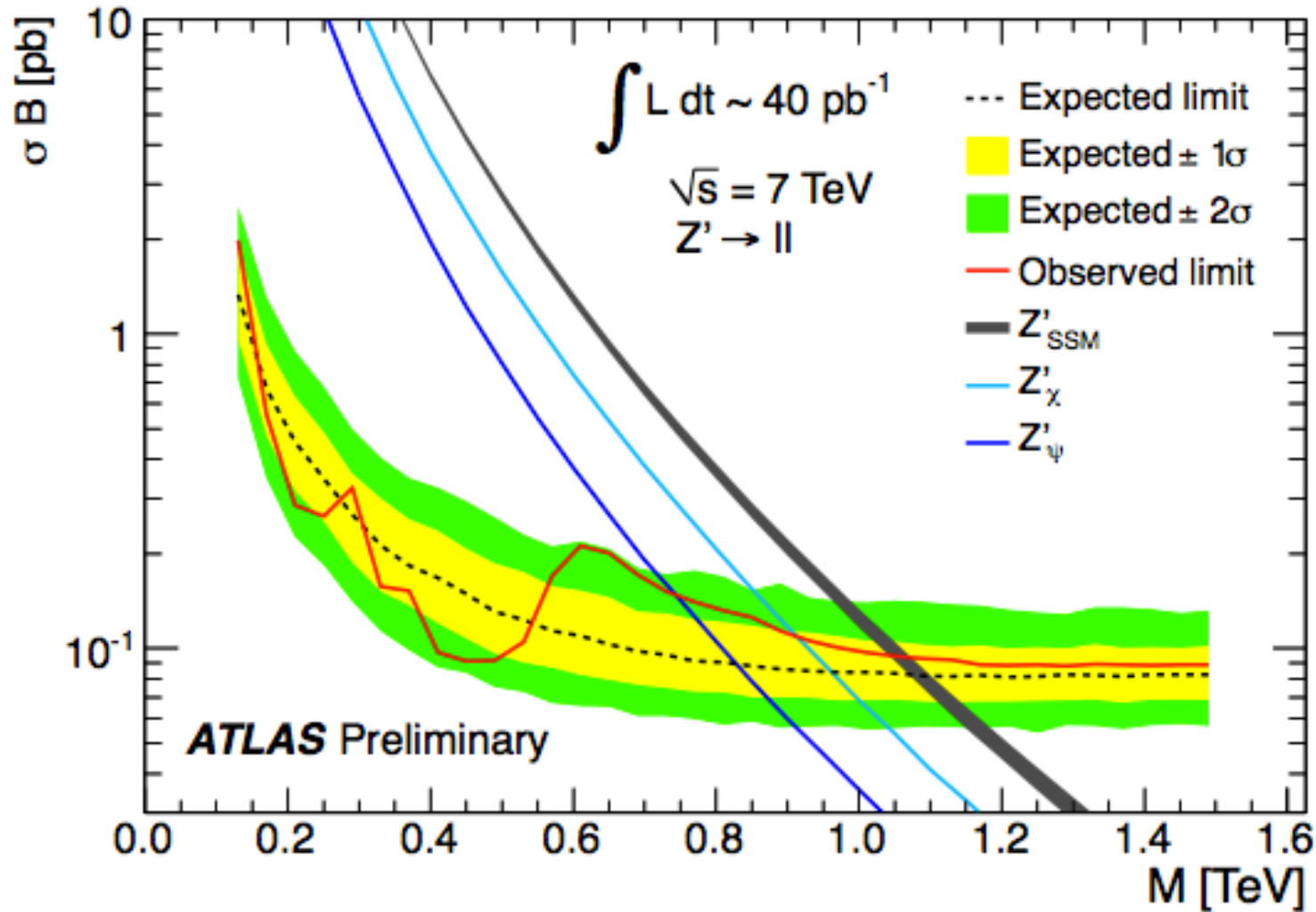
# ATLAS Z'

Penn  
+other groups



# Limits

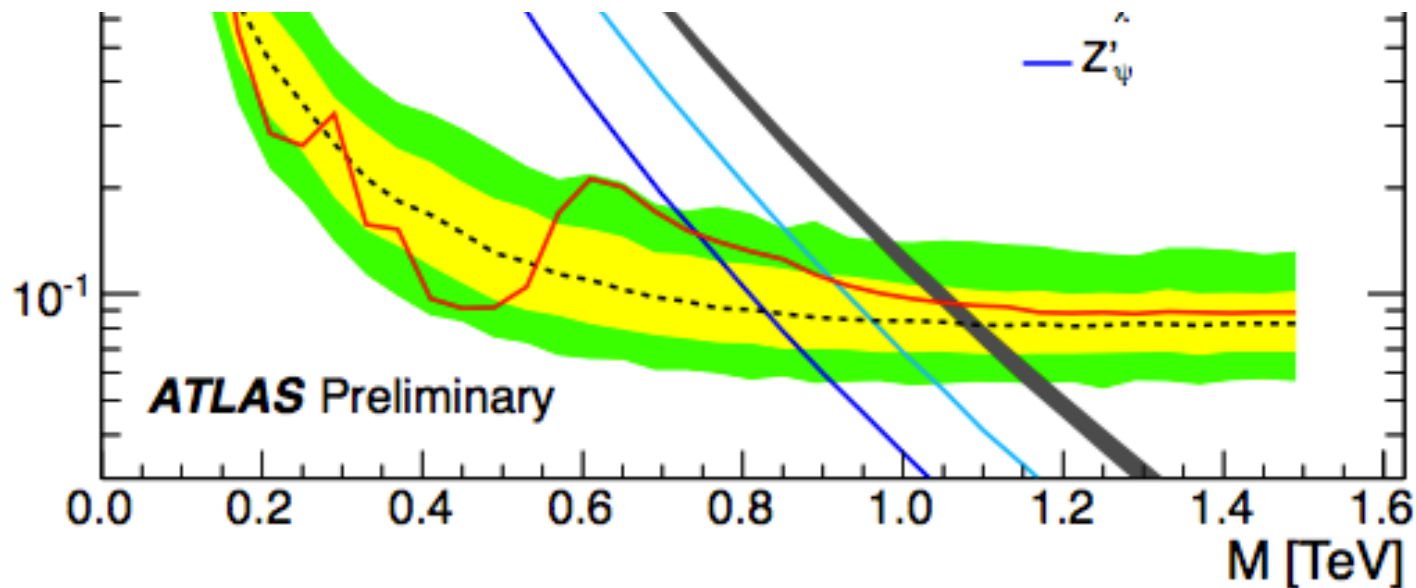
Penn  
+other groups



# Limits

Penn  
+other groups

	Observed limit		Expected limit	
	mass [TeV]	$\sigma B$ [pb]	mass [TeV]	$\sigma B$ [pb]
$Z'_{SSM} \rightarrow e^+e^-$	0.957	0.155	0.964	0.148
$Z'_{SSM} \rightarrow \mu^+\mu^-$	0.834	0.297	0.895	0.206
$Z'_{SSM} \rightarrow \ell^+\ell^-$	1.048	0.094	1.084	0.082





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a. *Heavy resonances ( $Z'$ )*

**b. *Heavy quarks ( $b'$ ,  $t'$ )***

c. *Simplified SUSY*

# 4th generation

PDG says it's  
ruled out to  $6\sigma$ ....

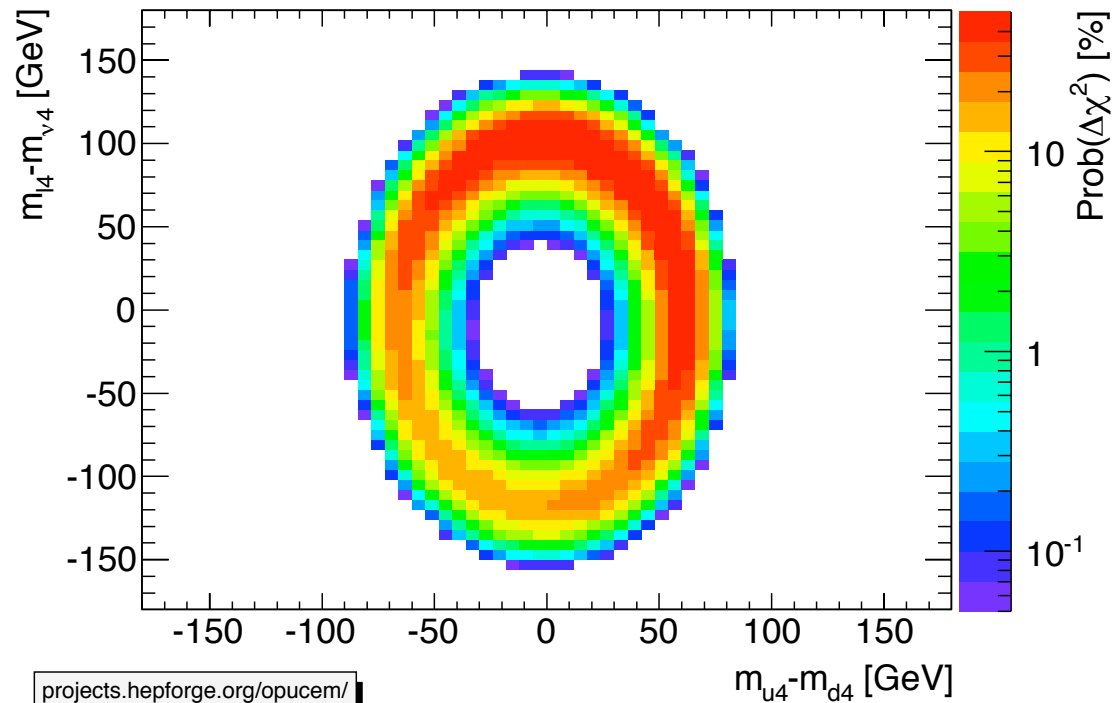
Quarks	u	c	t	t'
	d	s	b	b'
Leptons	$\nu_e$	$\nu_\mu$	$\nu_\tau$	$\nu'$
	e	$\mu$	$\tau$	$\tau'$
	I	II	III	IV

# 4th generation

PDG says it's ruled out to  $6\sigma$ ....

..that's true if the masses are degenerate

Quarks	u	c	t	t'
	d	s	b	b'
Leptons	$\nu_e$	$\nu_\mu$	$\nu_\tau$	$\nu'$
	e	$\mu$	$\tau$	$\tau'$
	I	II	III	IV



$t'$

UC Davis

## Selection

1 lepton

$p_t > 20 \text{ GeV}$

4 jets

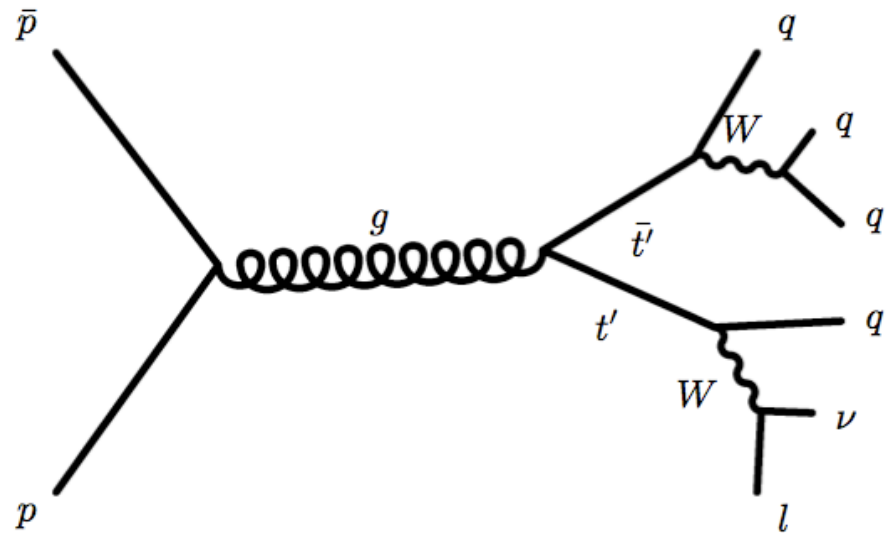
$p_t > 20 \text{ GeV}$

Missing transverse energy

$> 20 \text{ GeV}$

## Sample

4.6/fb

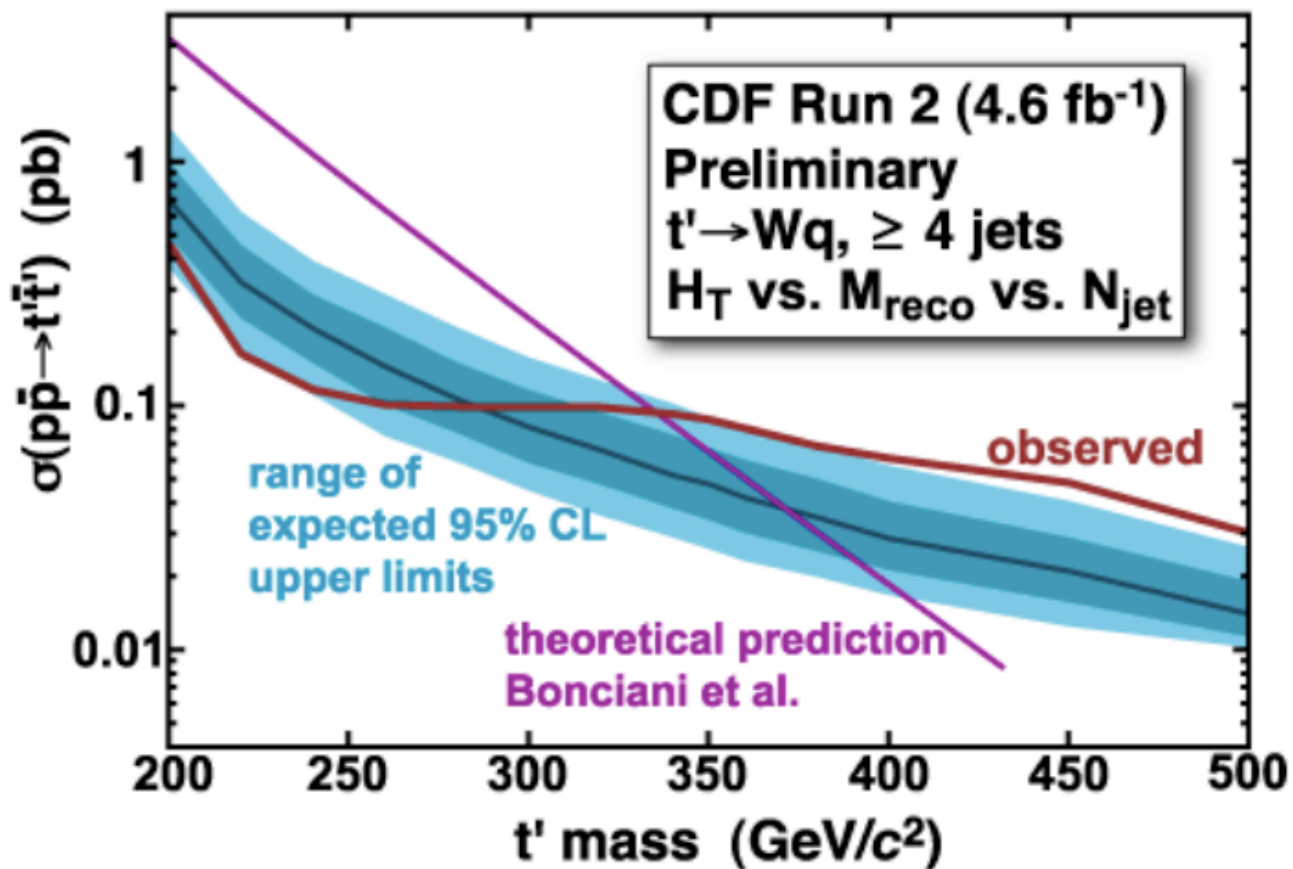


$t'$

UC Davis

Limit

$m_{t'} > 335 \text{ GeV}$



# $b'$

UCI Undergrad  
Matt Hickman



## Selection

2 like-signed leptons

$pt > 20$  GeV

at least **one** isolated

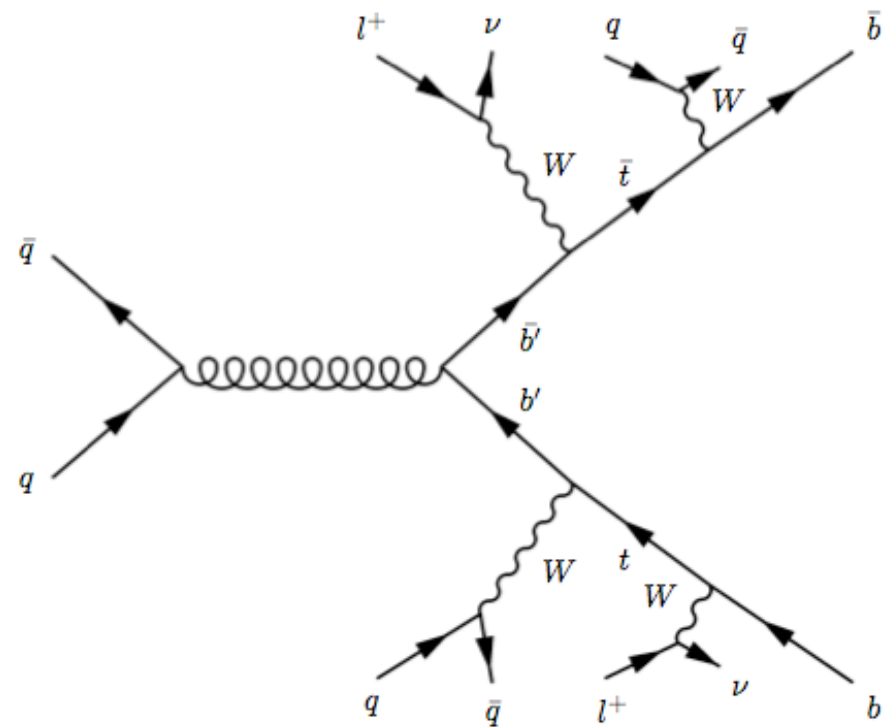
2 jets

$pt > 20$  GeV

$\geq 1$  btags

Missing transverse energy

$> 20$  GeV



## Sample

2.7/fb

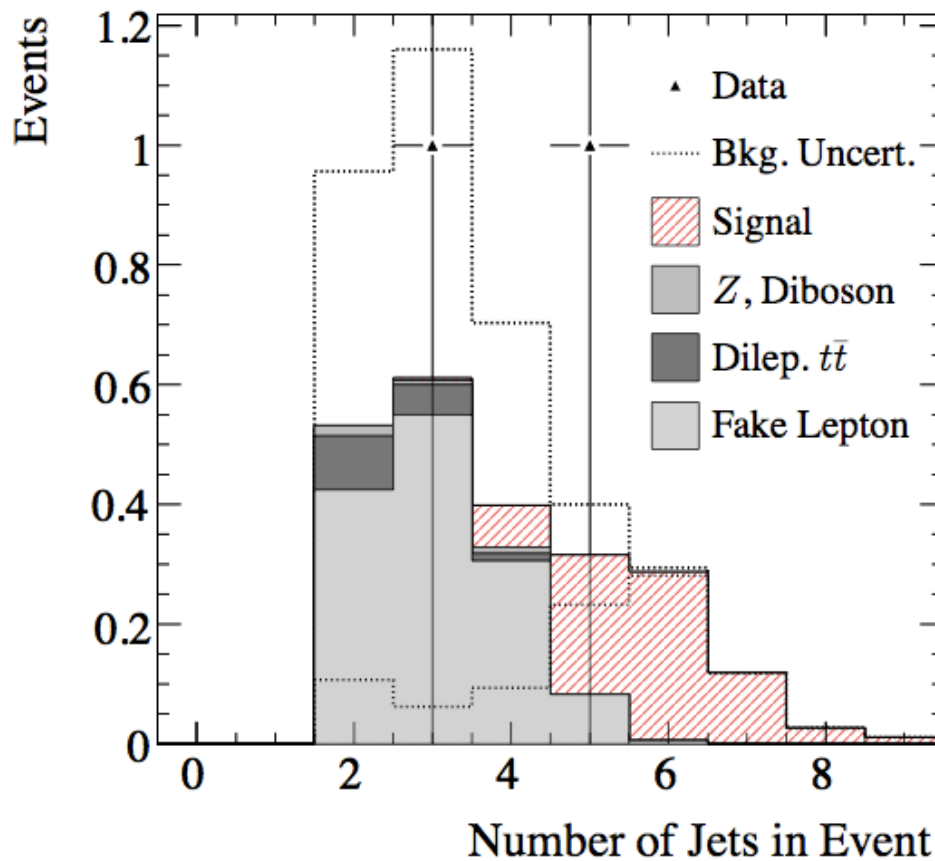
# $b'$

## Final selection

2 like-signed leptons

2 jets  $\geq 1$   $b$ tags

Missing transverse energy

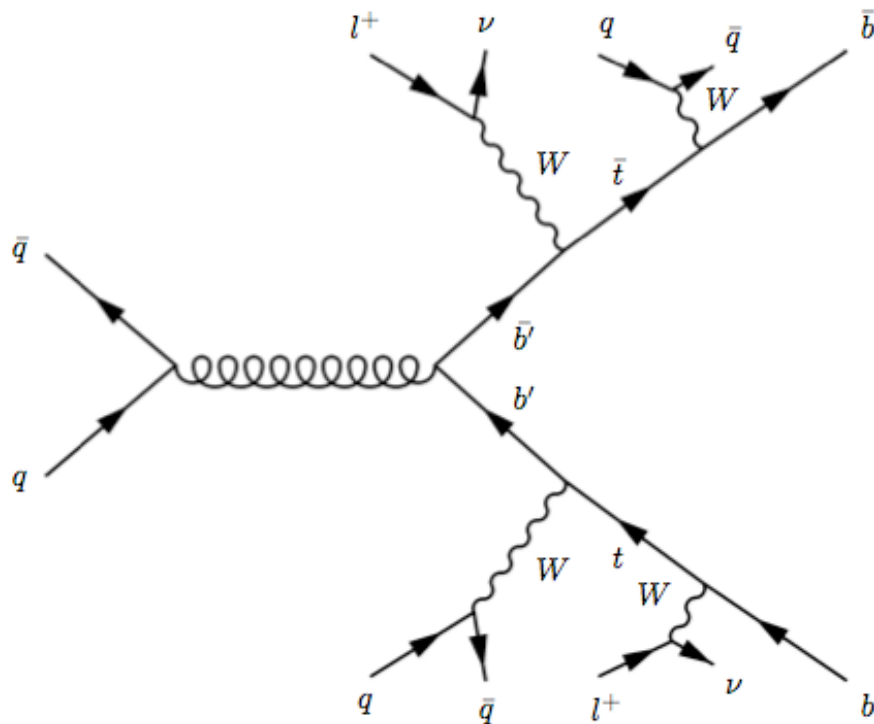


$m_{b'} > 338$  GeV

# b' decays

UCI undergrad  
Reza AmirArjomand

If  $b' \rightarrow Wt$

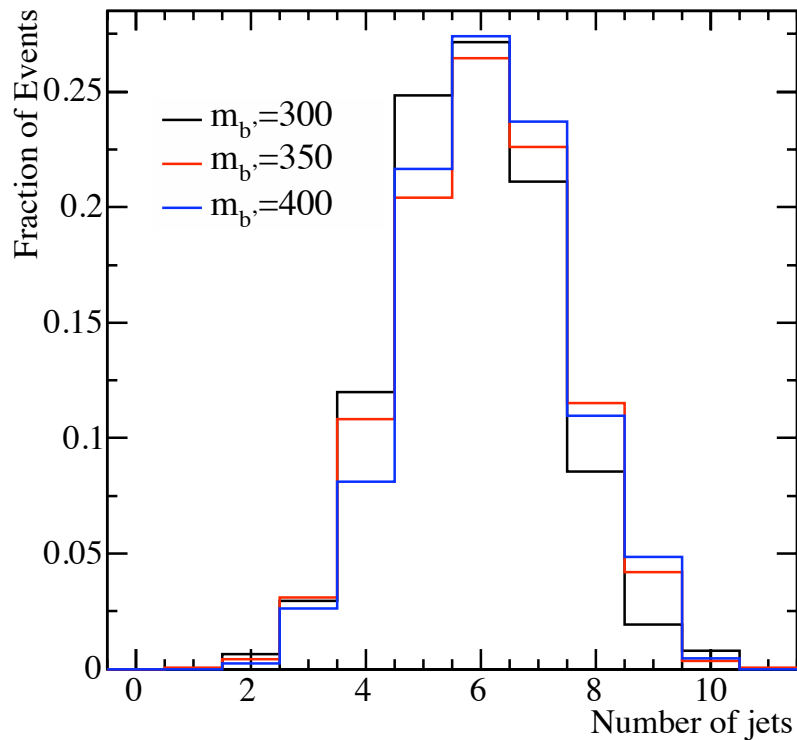


same-sign lepton selection:  $\sim 2\%$   
*consider single-lepton mode*

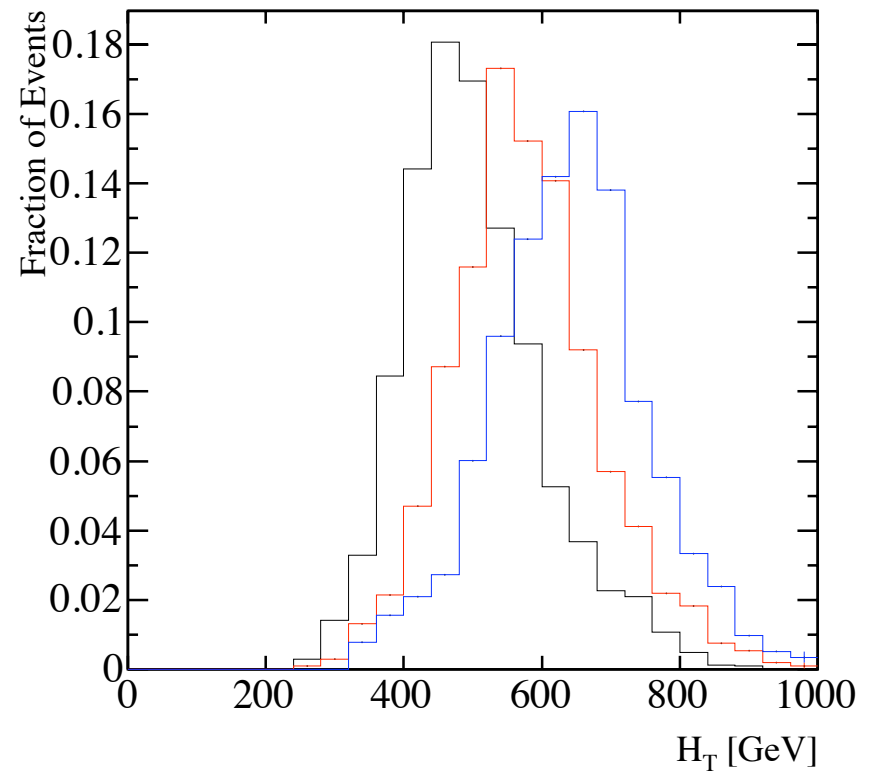


# Signal (madgraph)

CDF Run II Preliminary



CDF Run II Preliminary



Eight hard partons,  $\sim 6$  jets

# Signal (madgraph)

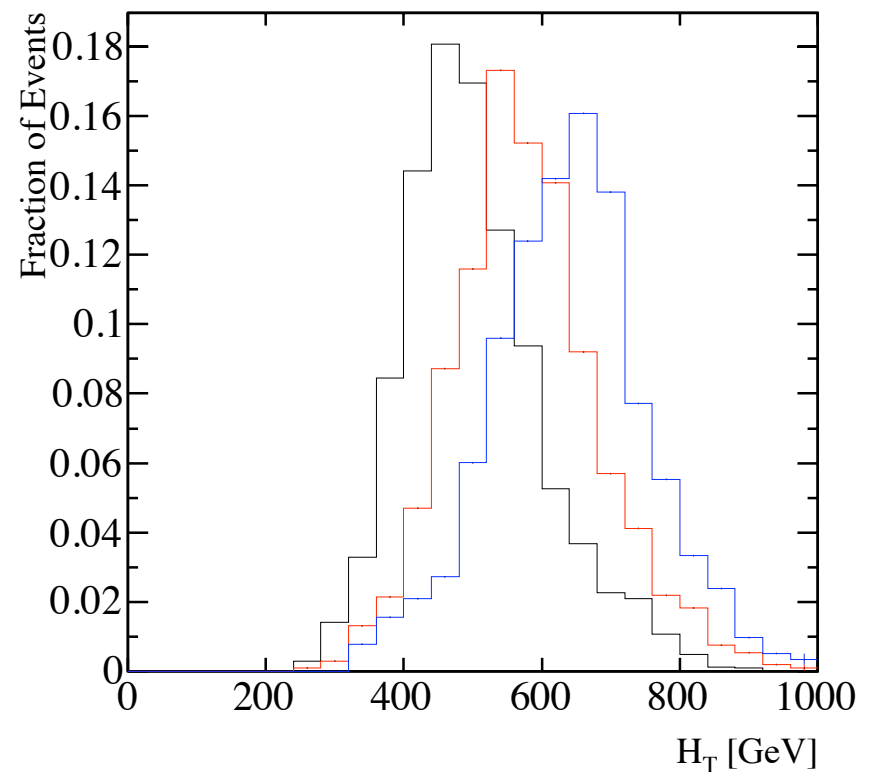
## HT

Scalar sum of transverse energy in the event

Includes jets, lepton and missing transverse energy

Captures soft recoil and unclustered jets

CDF Run II Preliminary



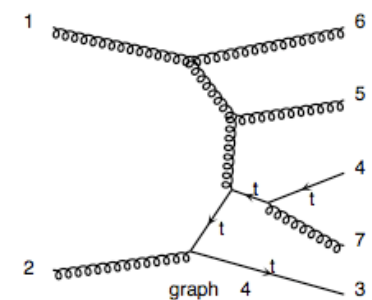
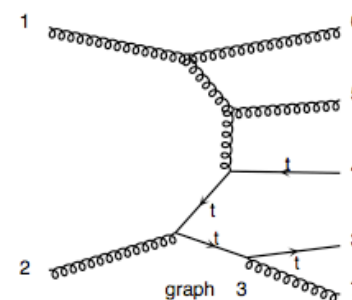
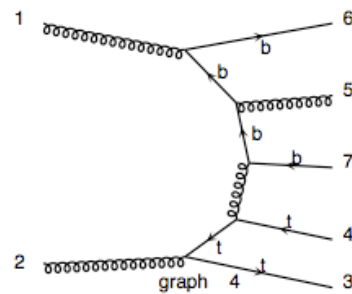
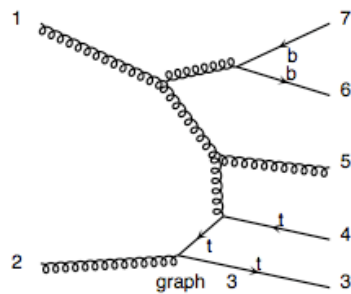
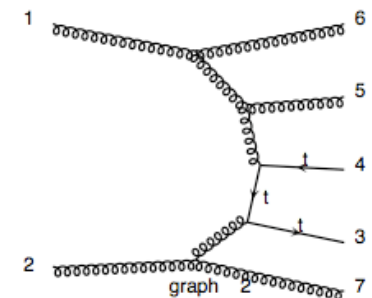
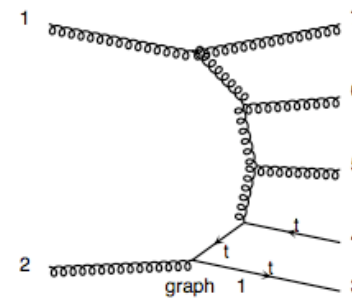
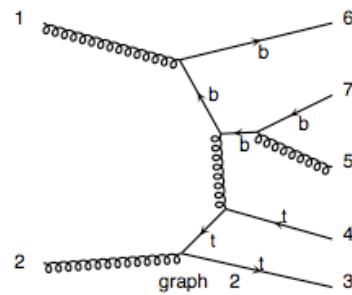
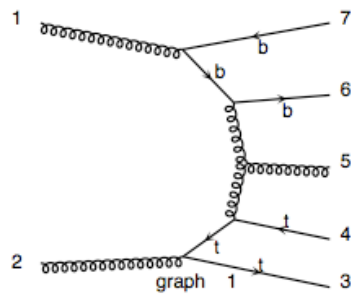
# top quark pair background

tt + 0,1,2,3p

p = udscb

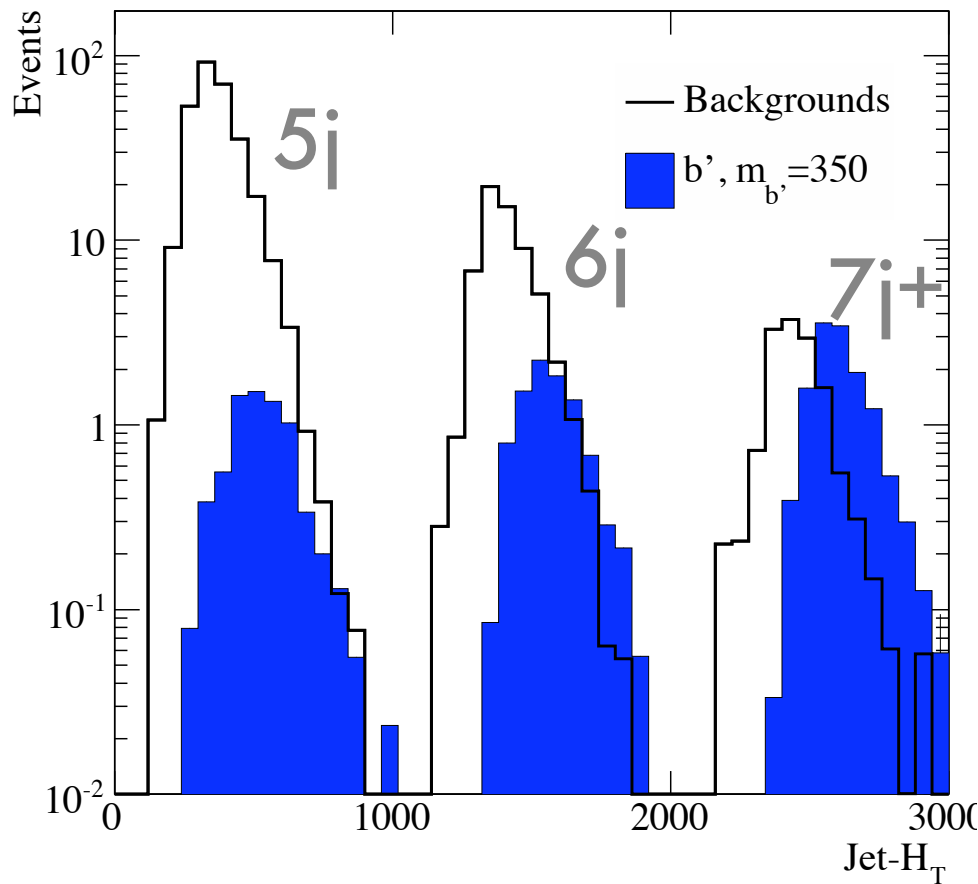
Madgraph+Pythia

MLM matching



# Analysis technique

CDF Run II Preliminary



normalized to 5/fb

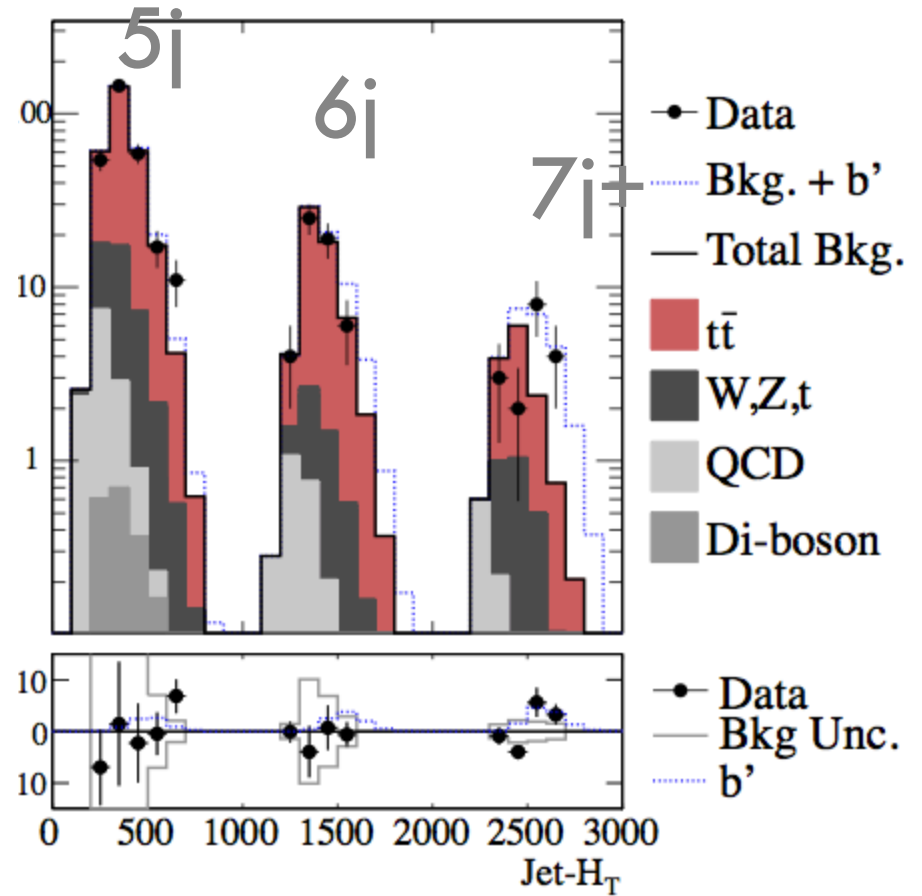
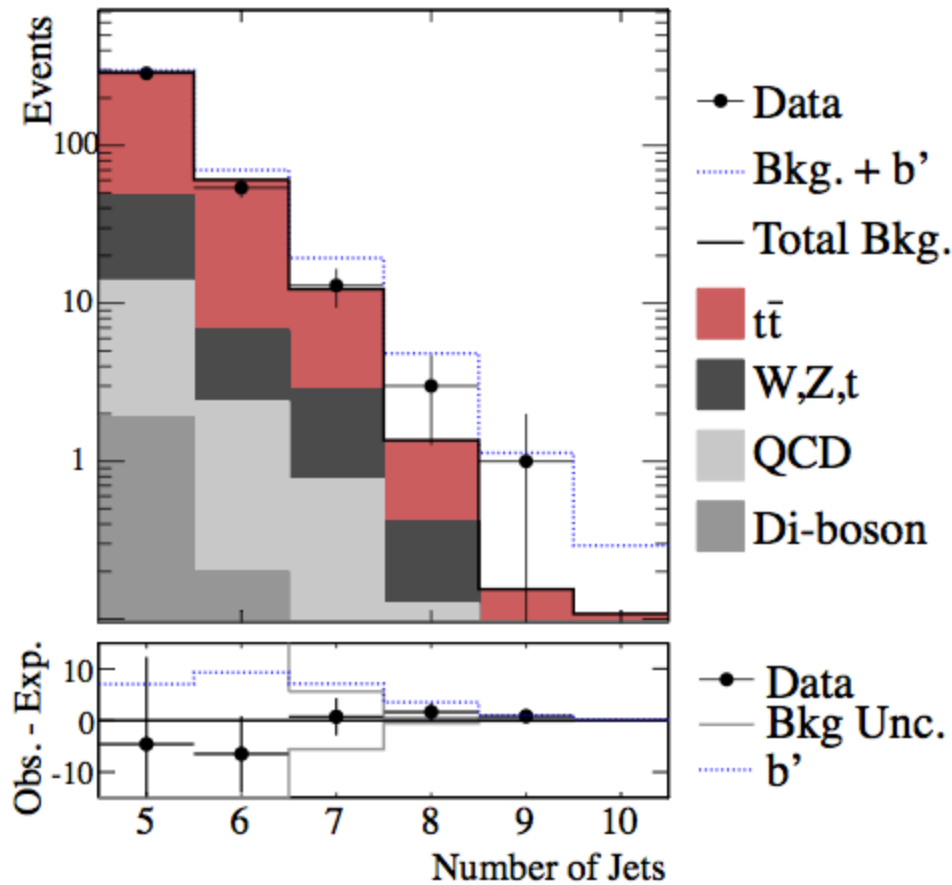
## Events

heavy *and* jetty

## Analysis variable

- if  $N_{jets} == 5$ ,  $Jet-H_T = H_T$ .
- if  $N_{jets} == 6$ ,  $Jet-H_T = H_T + 1000$  GeV.
- if  $N_{jets} \geq 7$ ,  $Jet-H_T = H_T + 2000$  GeV.

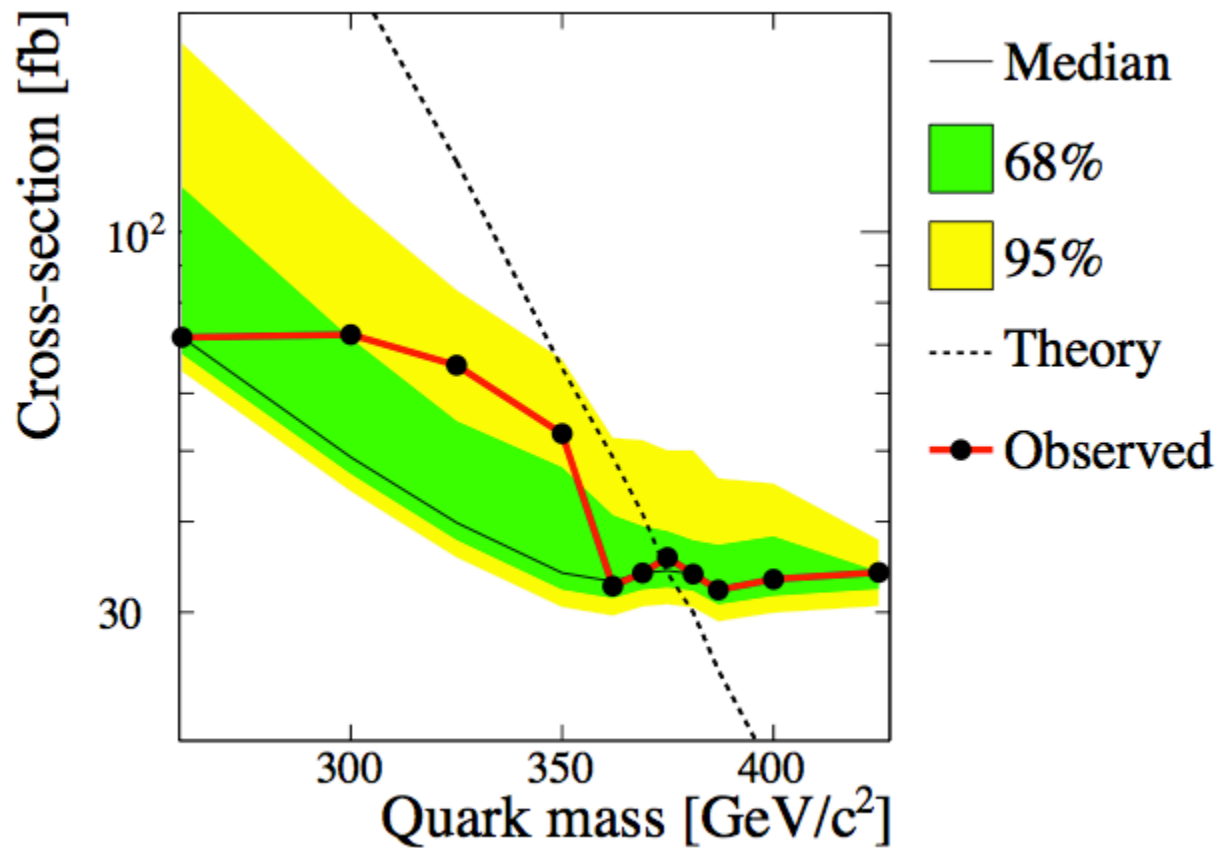
# Data, $\geq 1$ b-tag



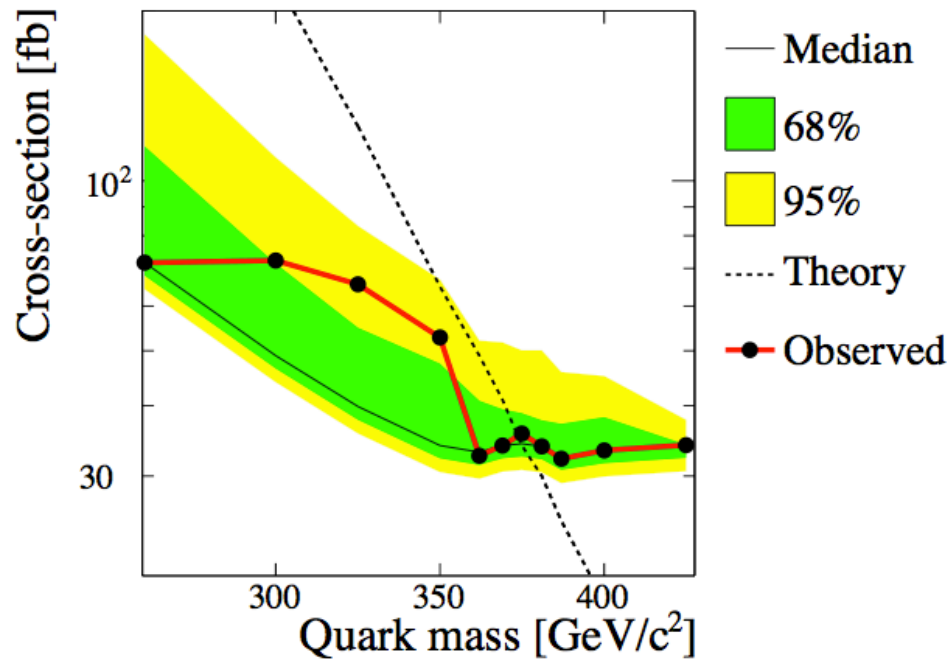
# The numbers

Jets	Control Region		Signal Region		Sum	
	Exp.	Obs.	Exp.	Obs.	Exp.	Obs.
5	207 ± 125	199	84 ± 65	87	291 ± 190	286
6	43 ± 31	40	18 ± 12	14	61 ± 43	54
≥ 7	11 ± 3.9	5	3.4 ± 3.4	12	14 ± 7.1	17

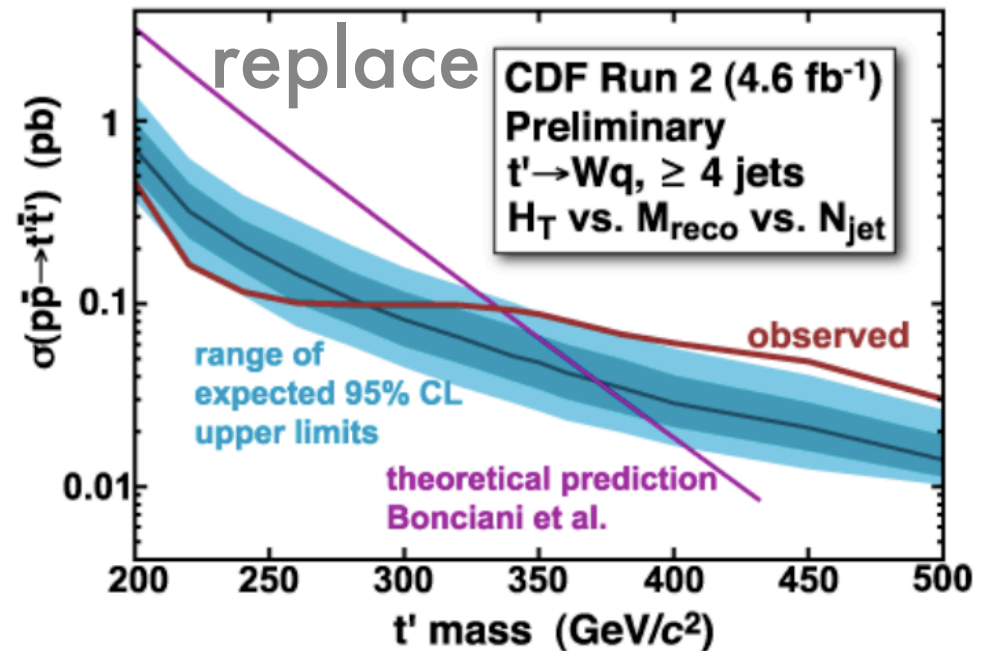
# The limits



# Direct searches



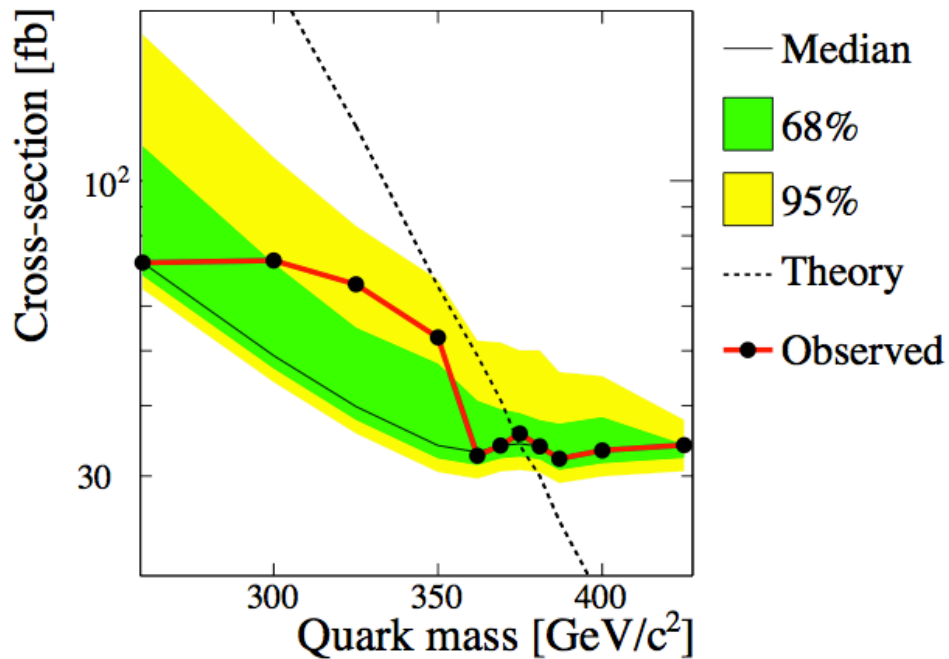
$m_{b'} > 372 \text{ GeV}$



$m_{t'} > 335 \text{ GeV}$

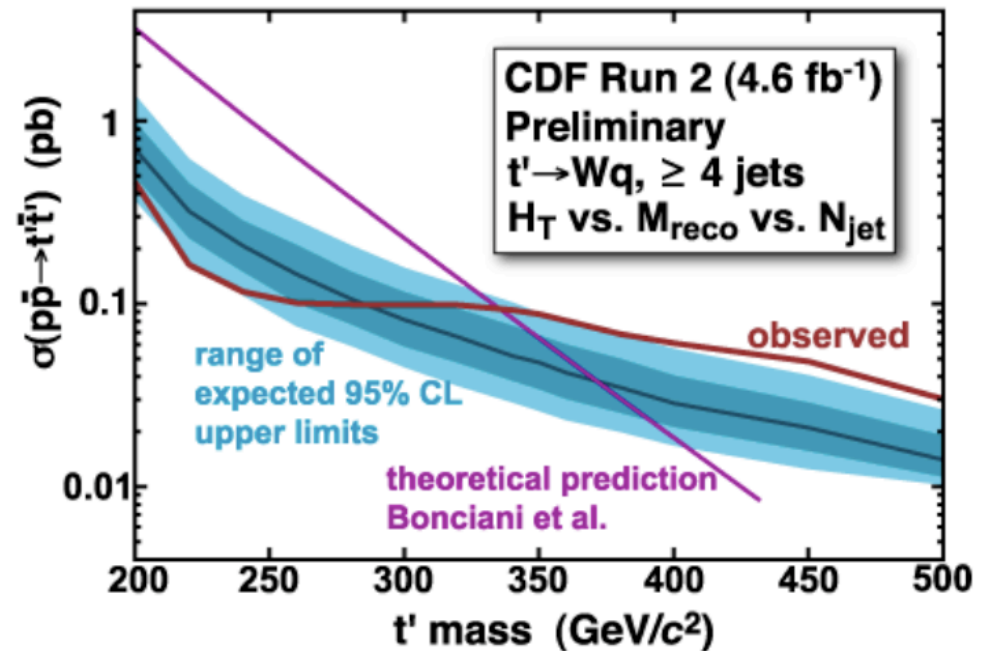


# Direct searches



$$\underline{m_{b'} > 372 \text{ GeV}}$$

*If  $BR(b' \rightarrow Wt) = 100\%$*



$$\underline{m_{t'} > 335 \text{ GeV}}$$

*If  $BR(t' \rightarrow Wq) = 100\%$*

# b' and t'

UCI postdoc

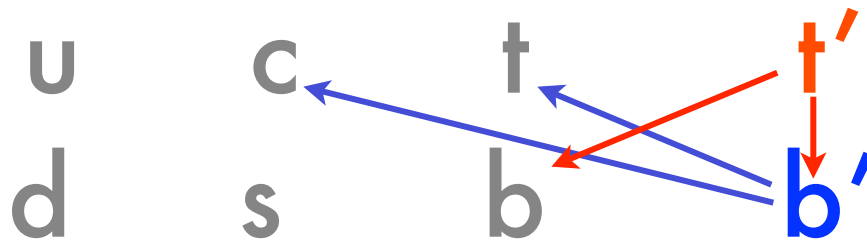
Christian Flacco

UCI undergrad

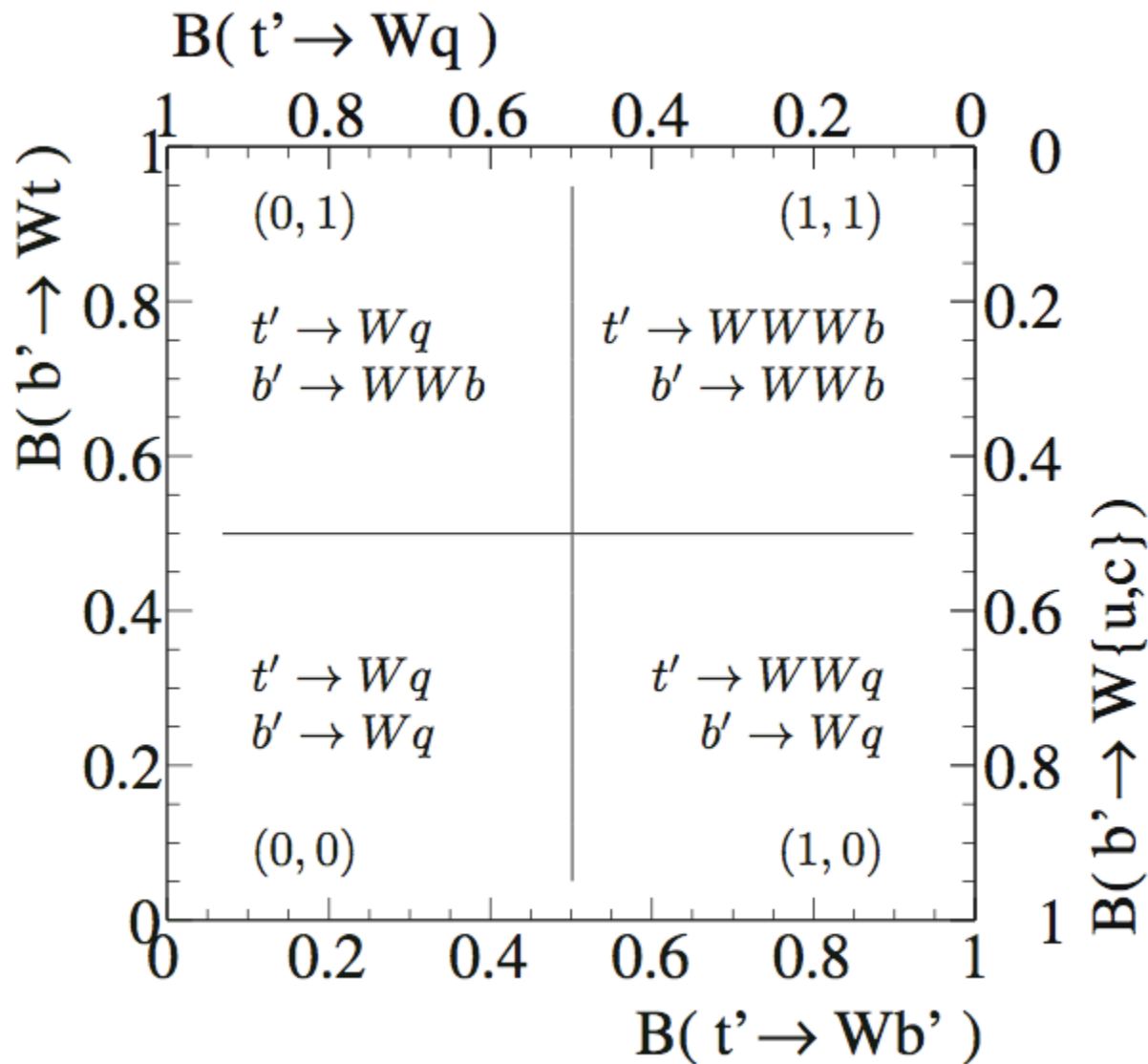
Matt Kelly



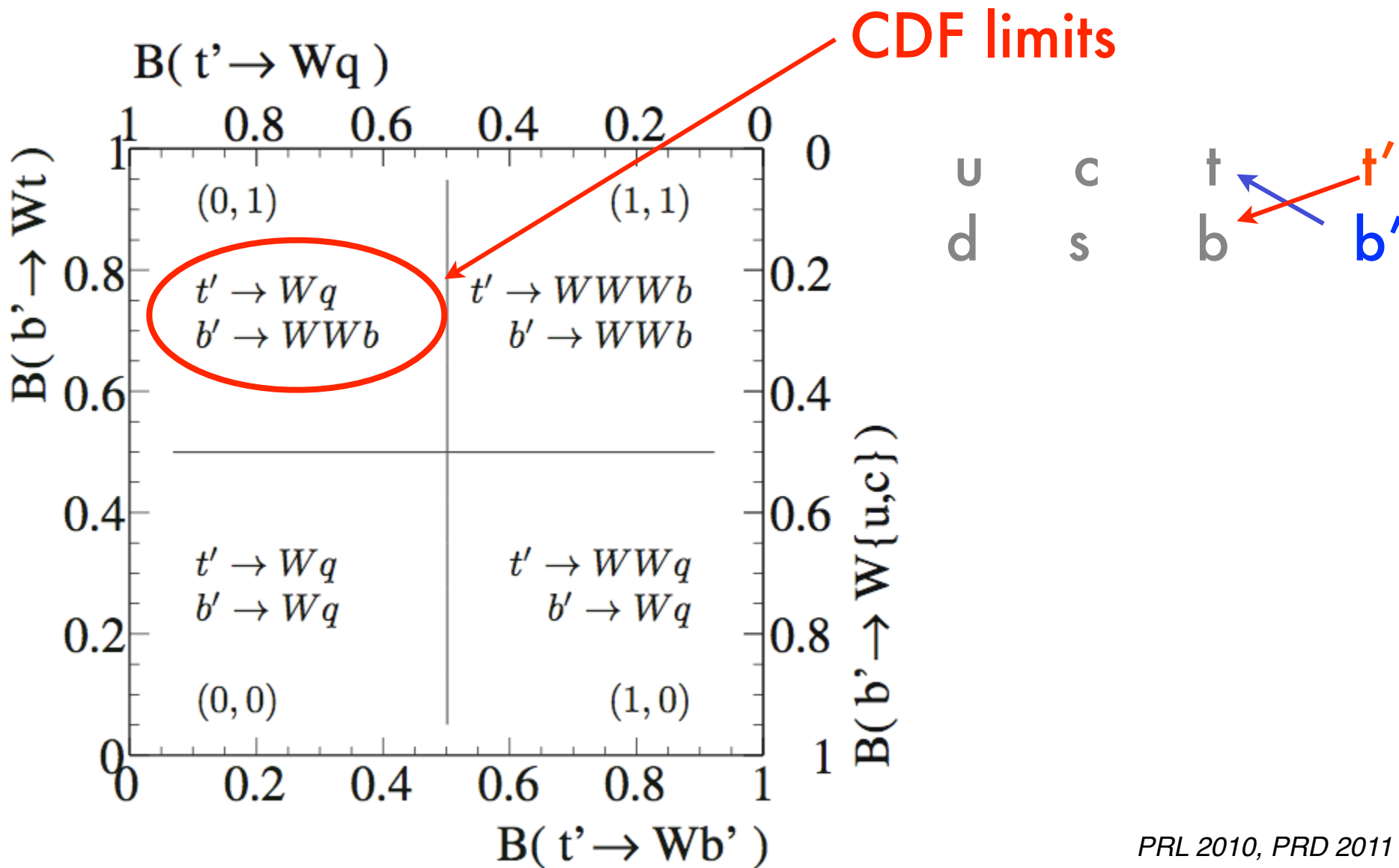
If  $m_{t'} > m_{b'}$



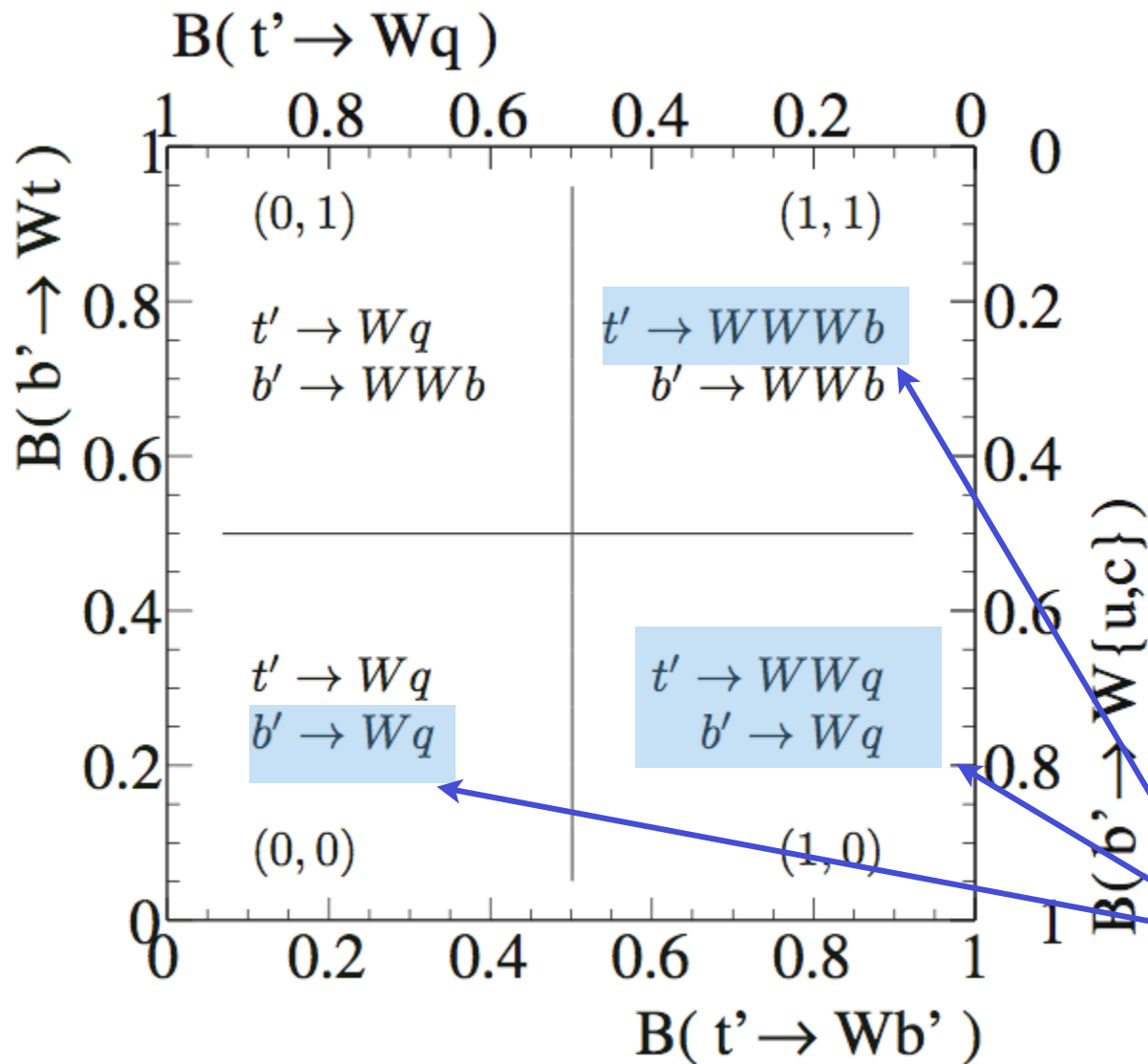
# $b'$ and $t'$



# $b'$ and $t'$

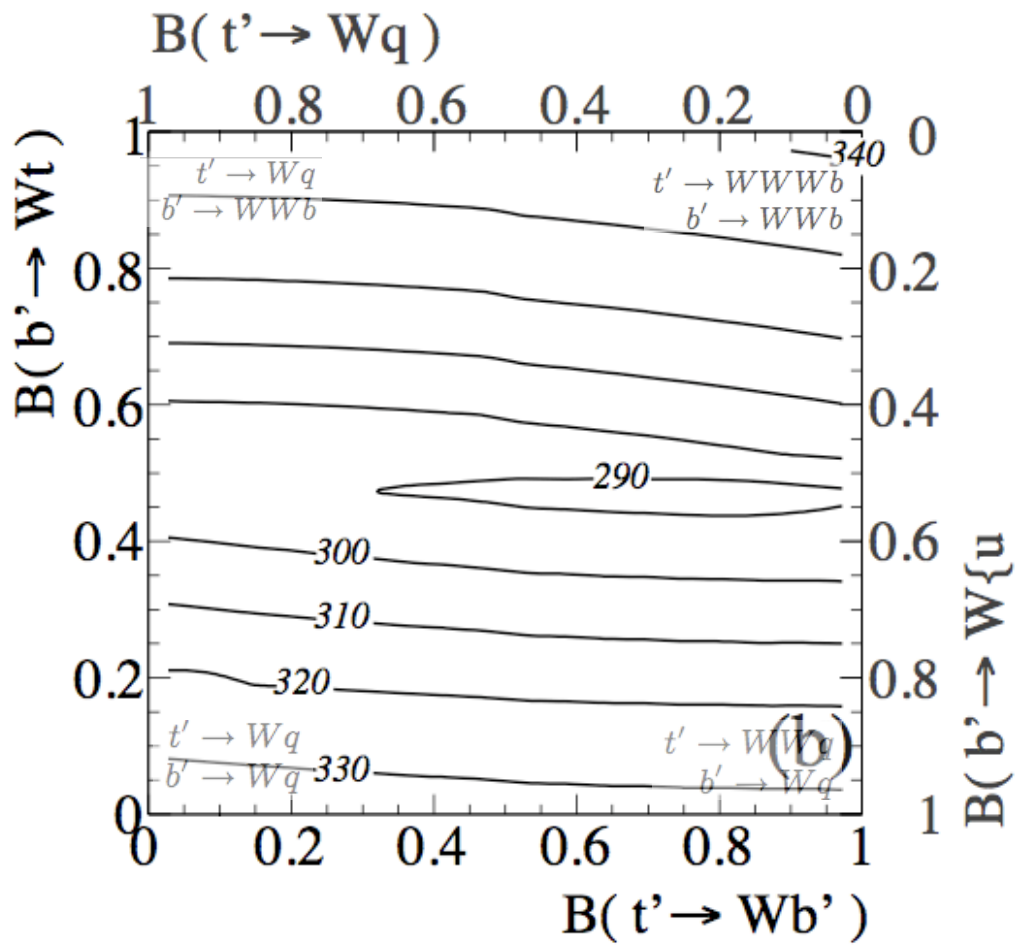


# $b'$ and $t'$

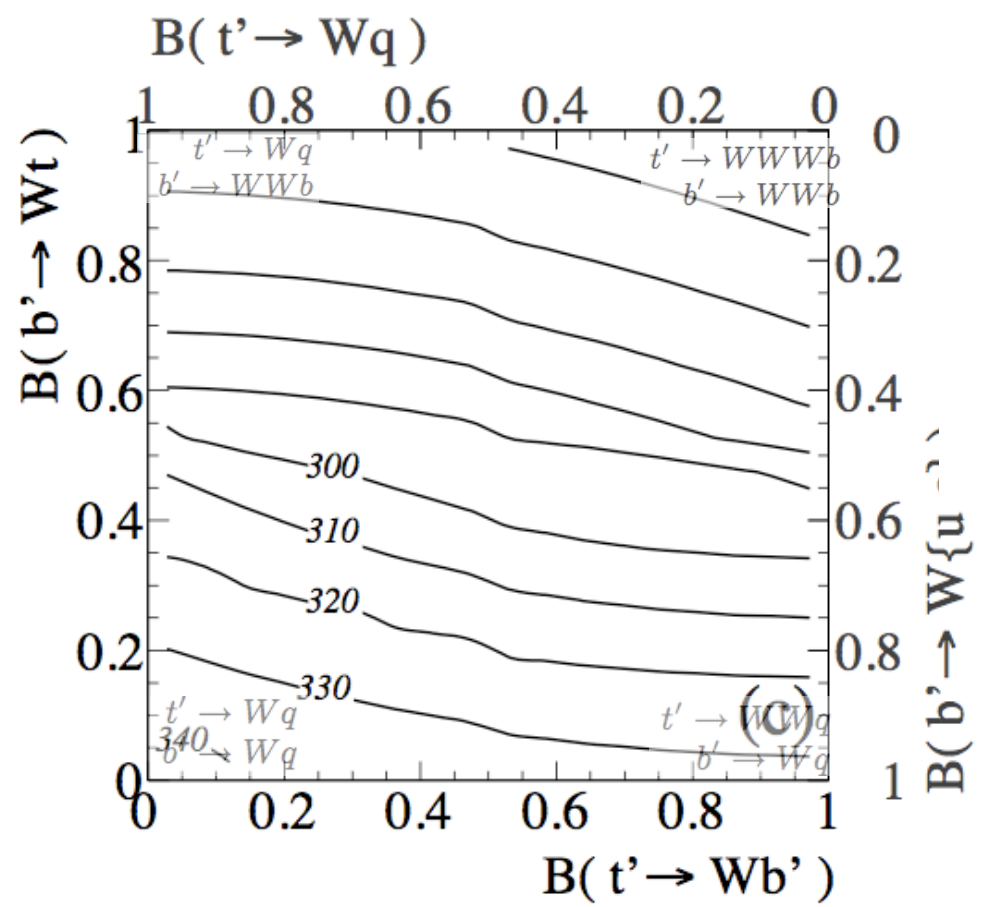


No direct limits!

# $t'$ and $b'$



$$m_{t'} = m_{b'} + 100$$

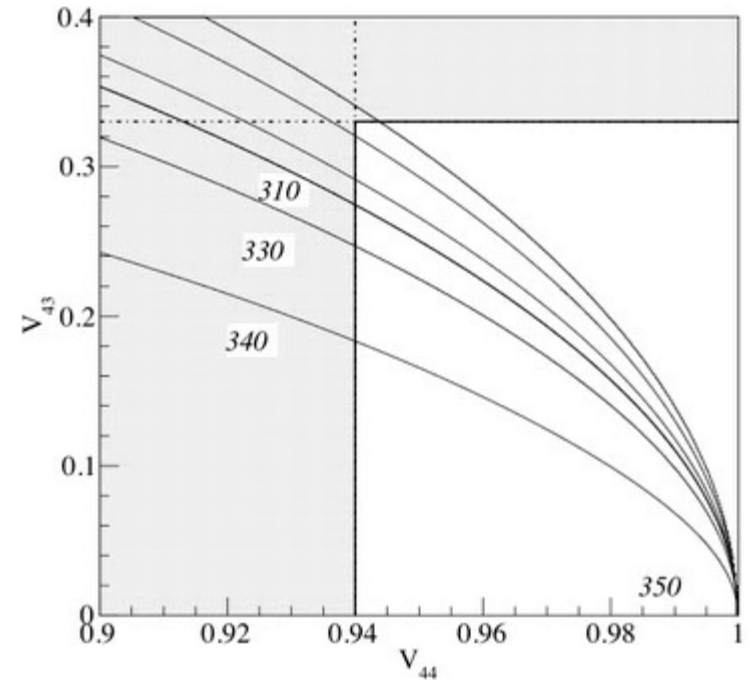
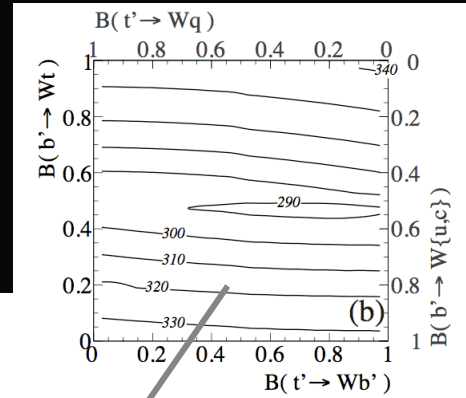


$$m_{t'} = m_{b'} + 50$$

# Limits

$$V_{CKM4} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} & V_{ub'} \\ V_{cd} & V_{cs} & V_{cb} & V_{cb'} \\ V_{td} & V_{ts} & V_{tb} & V_{tb'} \\ V_{t'd} & V_{t's} & V_{t'b} & V_{t'b'} \end{pmatrix}$$

$V_{ud} = 0.97418 \pm 0.00027$	Nuclear Beta decay
$V_{us} = 0.2255 \pm 0.0019$	Semileptonic K-decay
$V_{ub} = 0.00393 \pm 0.00036$	Semileptonic B-decay
$V_{cd} = 0.230 \pm 0.011$	Semileptonic D-decay
$V_{cs} = 1.04 \pm 0.06$	Semi- /Leptonic D-decay
$V_{cb} = 0.0412 \pm 0.0011$	Semileptonic B-decay
$V_{tb} > 0.74$	Single Top-production



# heavy quarks

If the lifetime is short enough  
so the decay is in the central detector:

$$m_{Q'} > 290 \text{ GeV}$$



# ATLAS $t'$

UCI grad student  
Michael Werth

## Selection

2 OS leptons

$p_t > 20$  GeV

2 jets

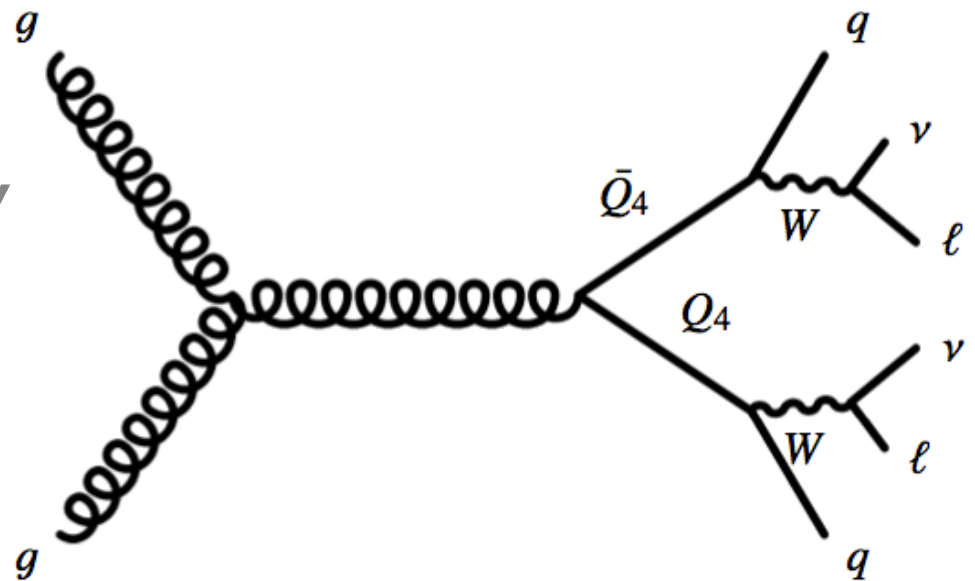
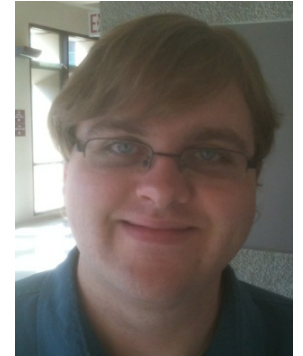
$p_t > 20$  GeV

Missing transverse energy

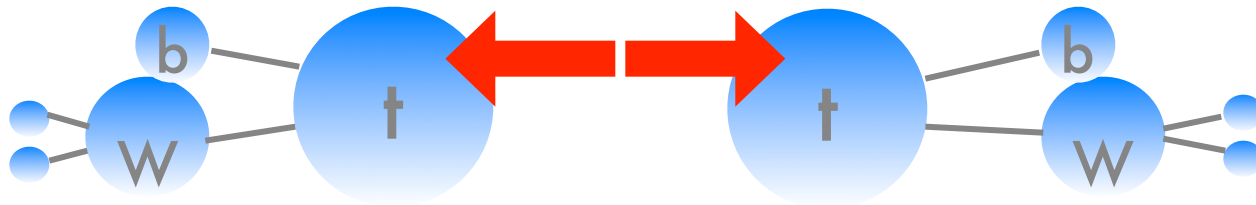
$> 20$  GeV

## Sample

35/pb

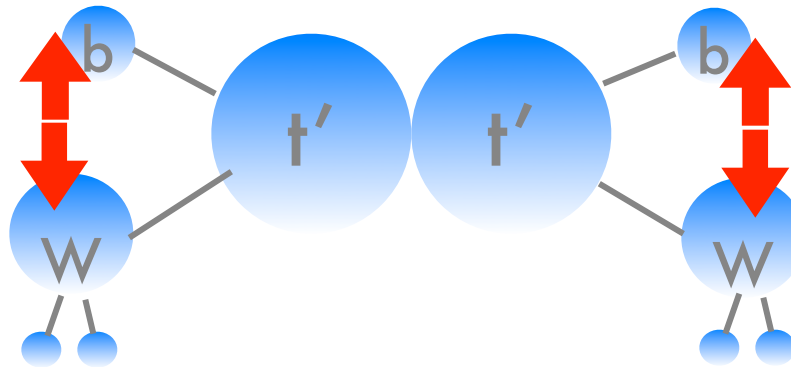
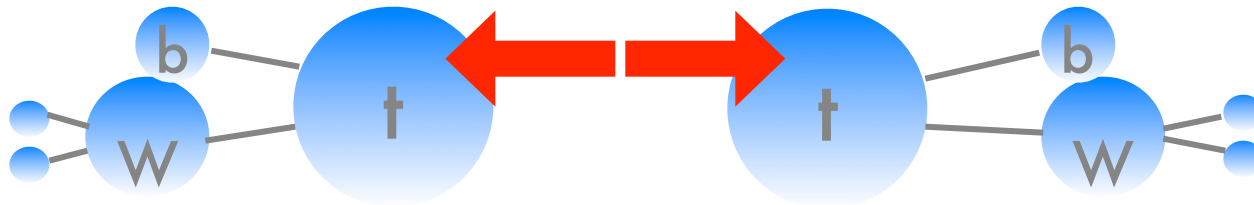


# topology



Boosted tops

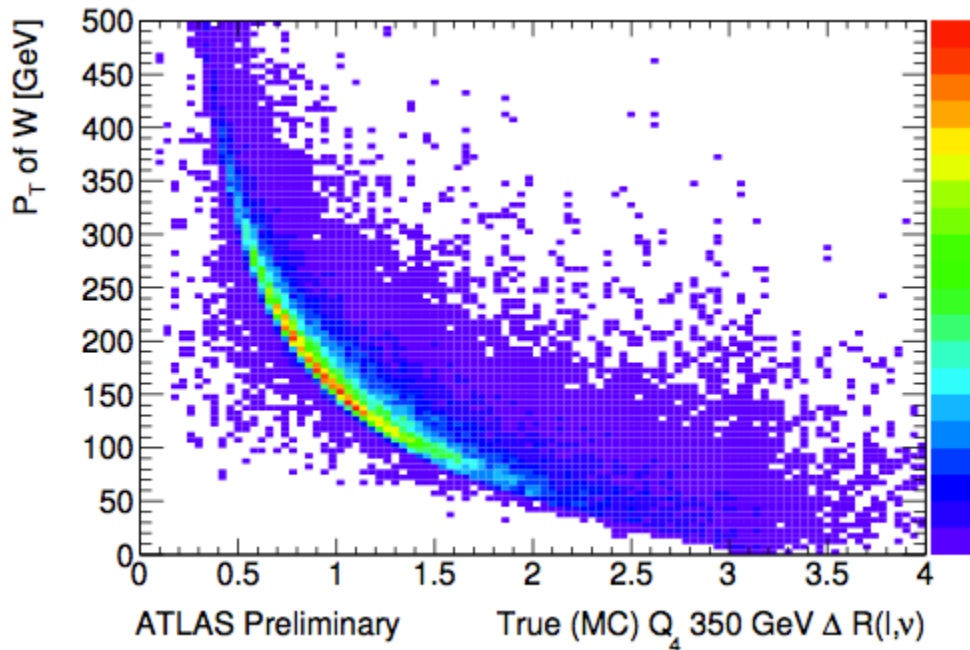
# topology



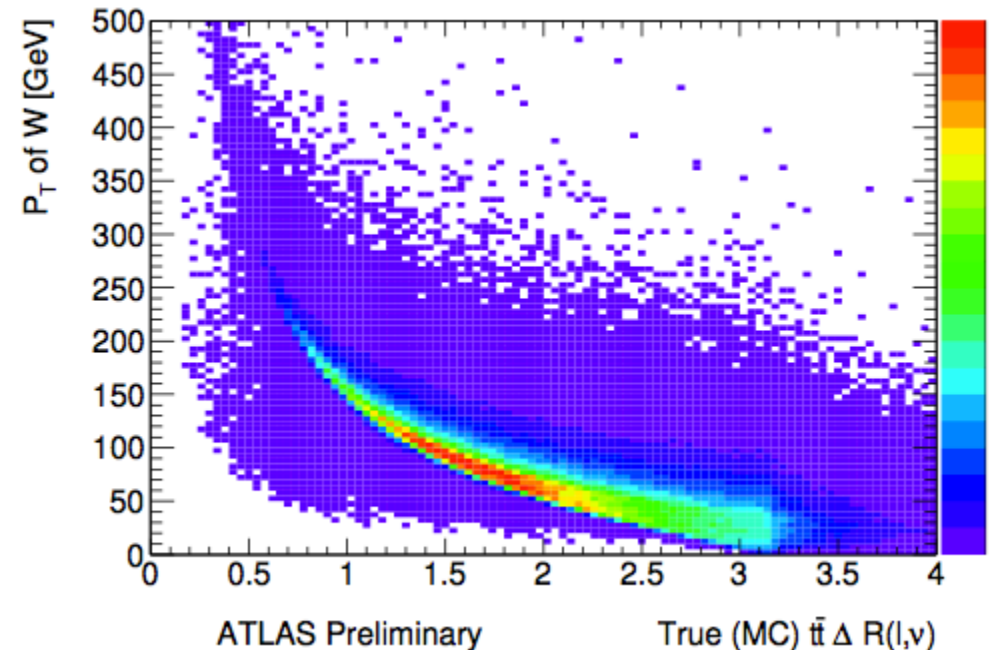
Boosted Ws!

# Lepton-neutrino angles

Heavy  $t'$



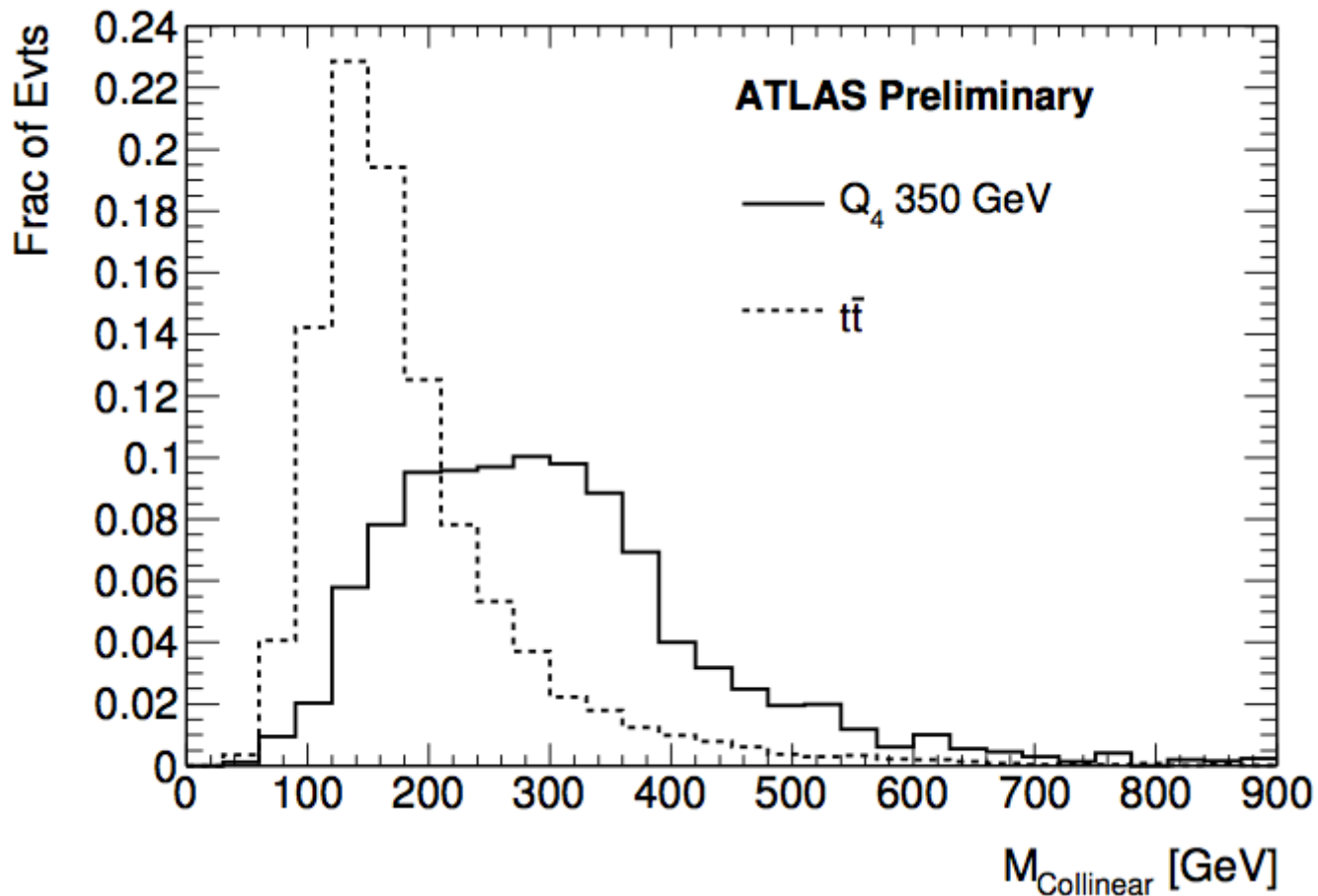
SM top



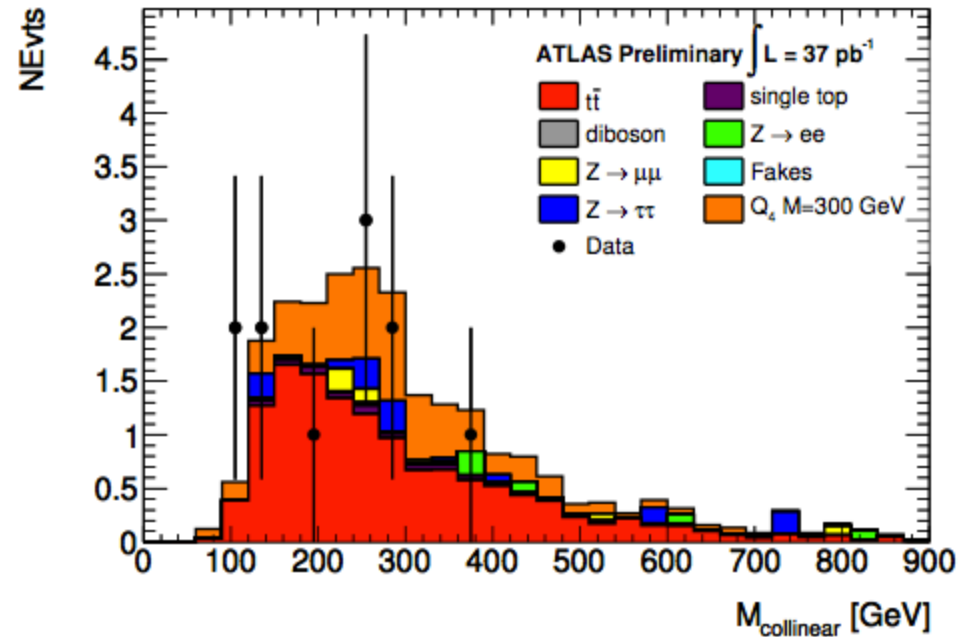
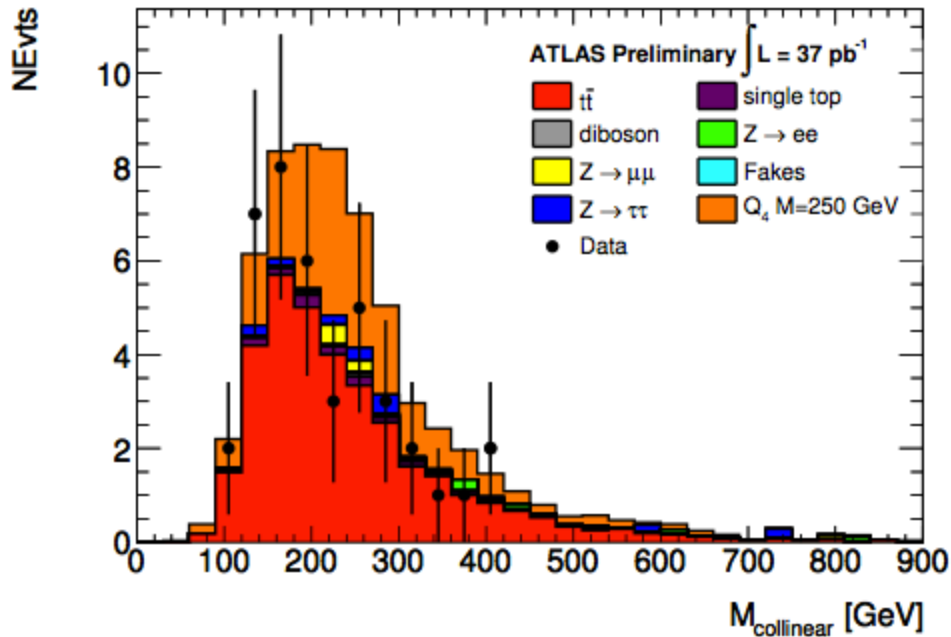
More W  $p_T$  means smaller opening angle

# Mass reconstruction

Assume lepton and neutrino are  $\sim$  collinear

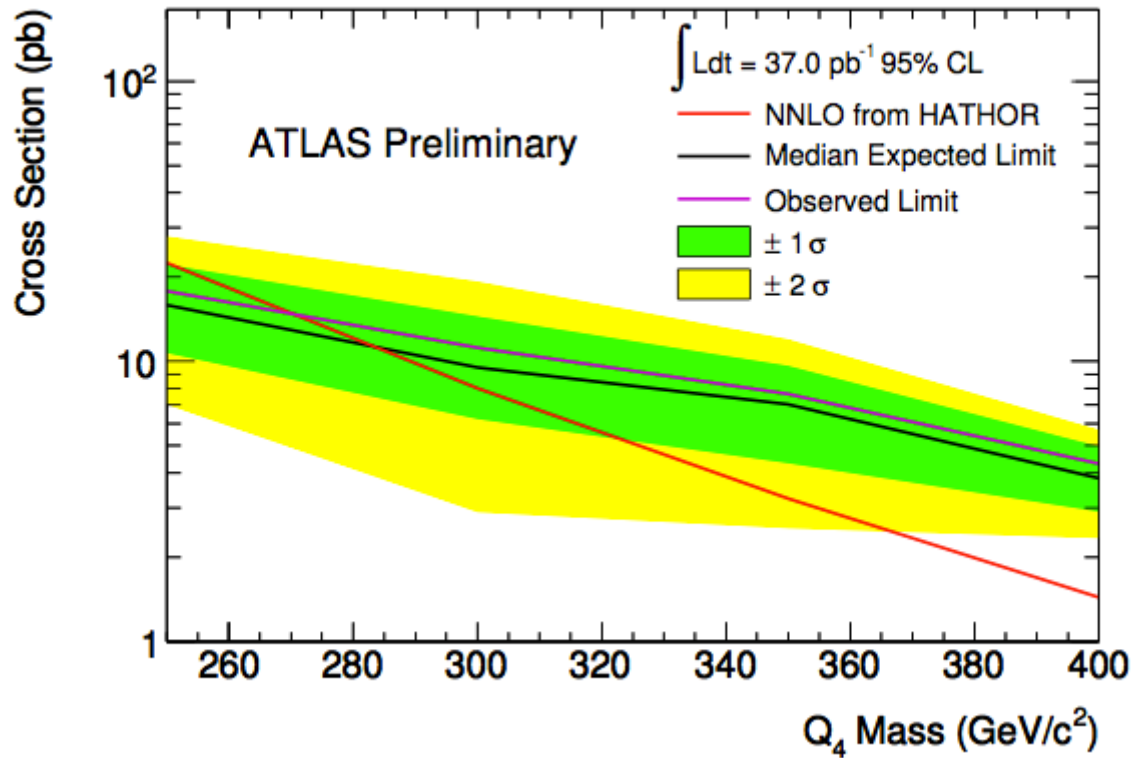


# Data



No sign of heavy quarks...

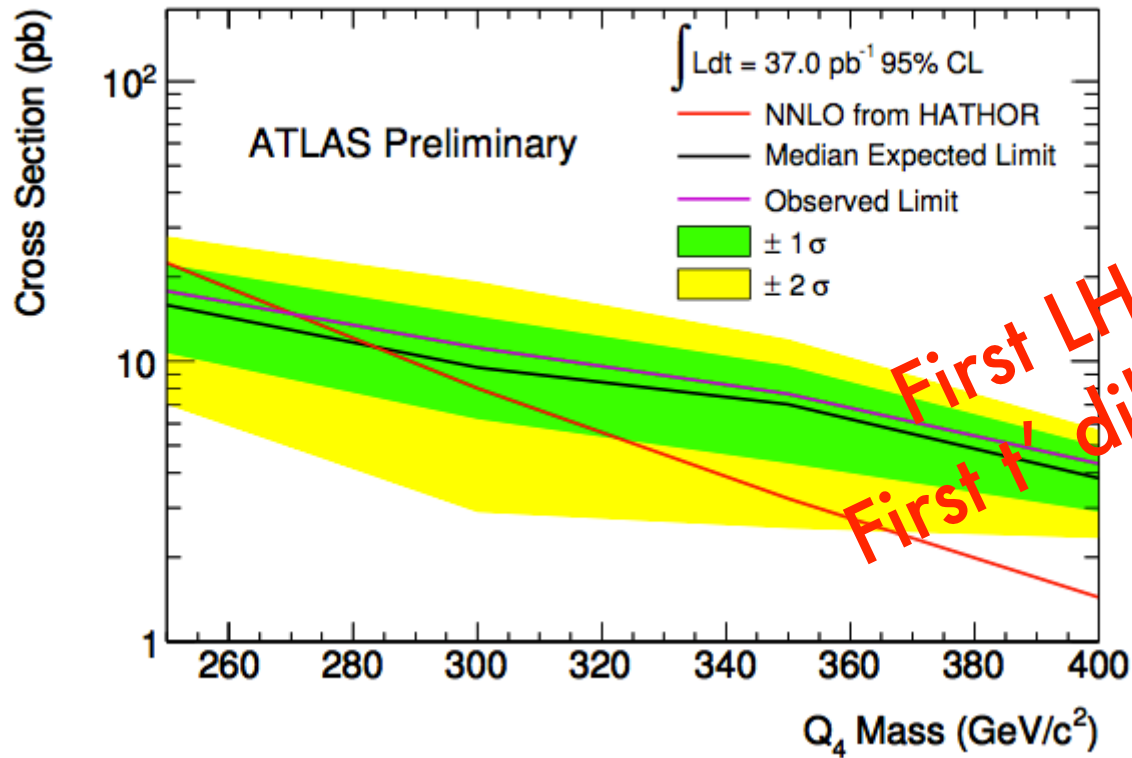
# Limit



Limit

$$m_{t'} > 275 \text{ GeV}$$

# Limit



Limit

$$m_{t'} > 275 \text{ GeV}$$



# Dark Matter

Need long lived dark matter X

# Dark Matter

Need long lived dark matter X  
Give it some dark charge that is conserved  
(eg R-parity for susy LSP)

SM Particles

SM Charges

Dark Matter X

Dark Charge

# Dark Matter

Need long lived dark matter  $X$   
Give it some dark charge that is conserved  
(eg R-parity for susy LSP)

SM Particles

SM Charges

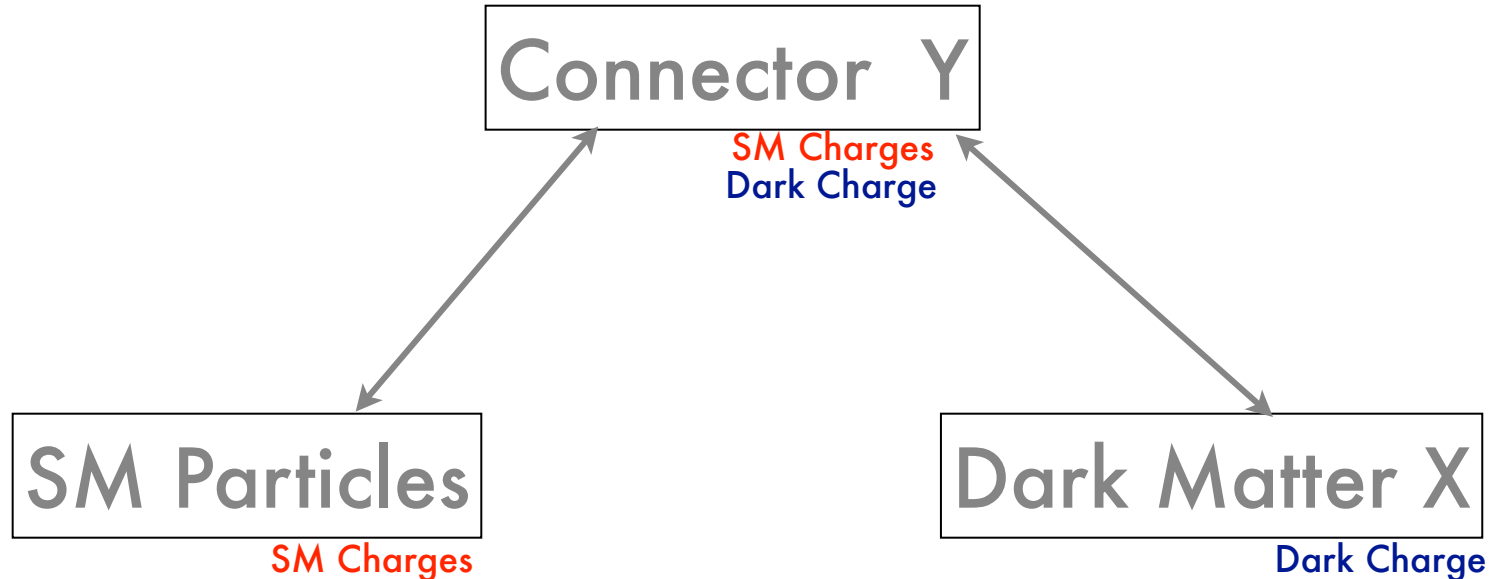
Dark Matter  $X$

Dark Charge

$X$  can't be light ( $\tilde{m} < 10$  GeV) and carry SM charges to be consistent with relic density.

# Dark Matter

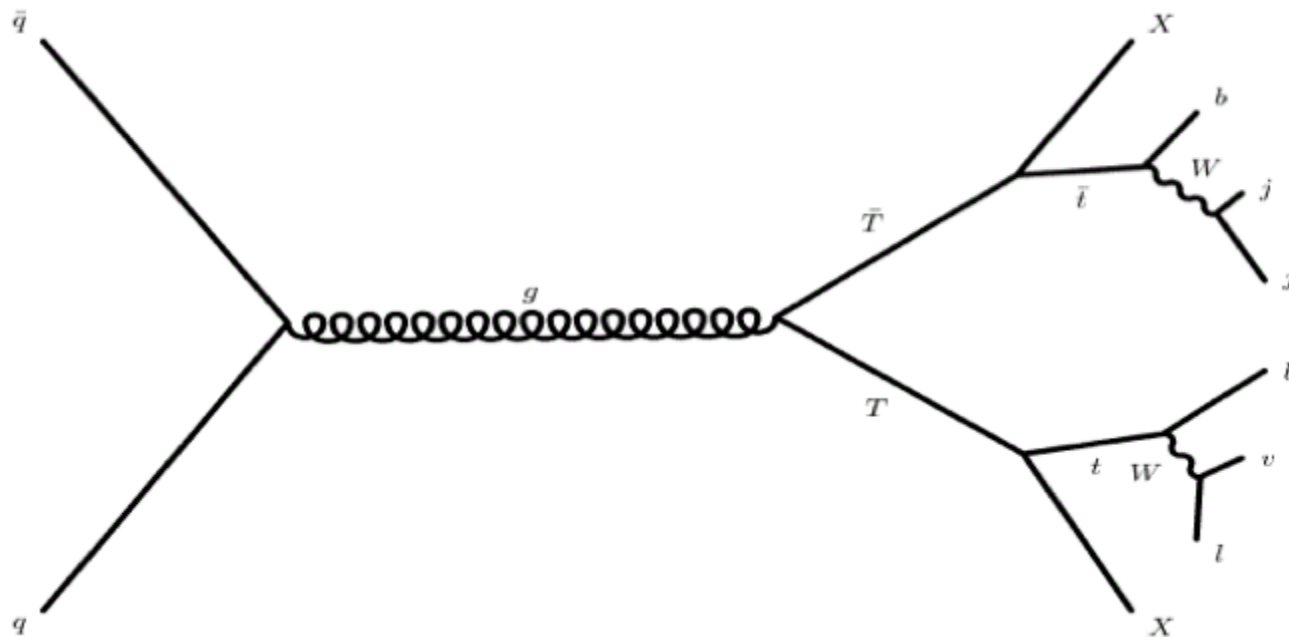
Need long lived dark matter X



Produce Y, decay as  $Y \rightarrow f X$

# Dark Matter+4th gen

UCI grad student  
Kanishka Rao

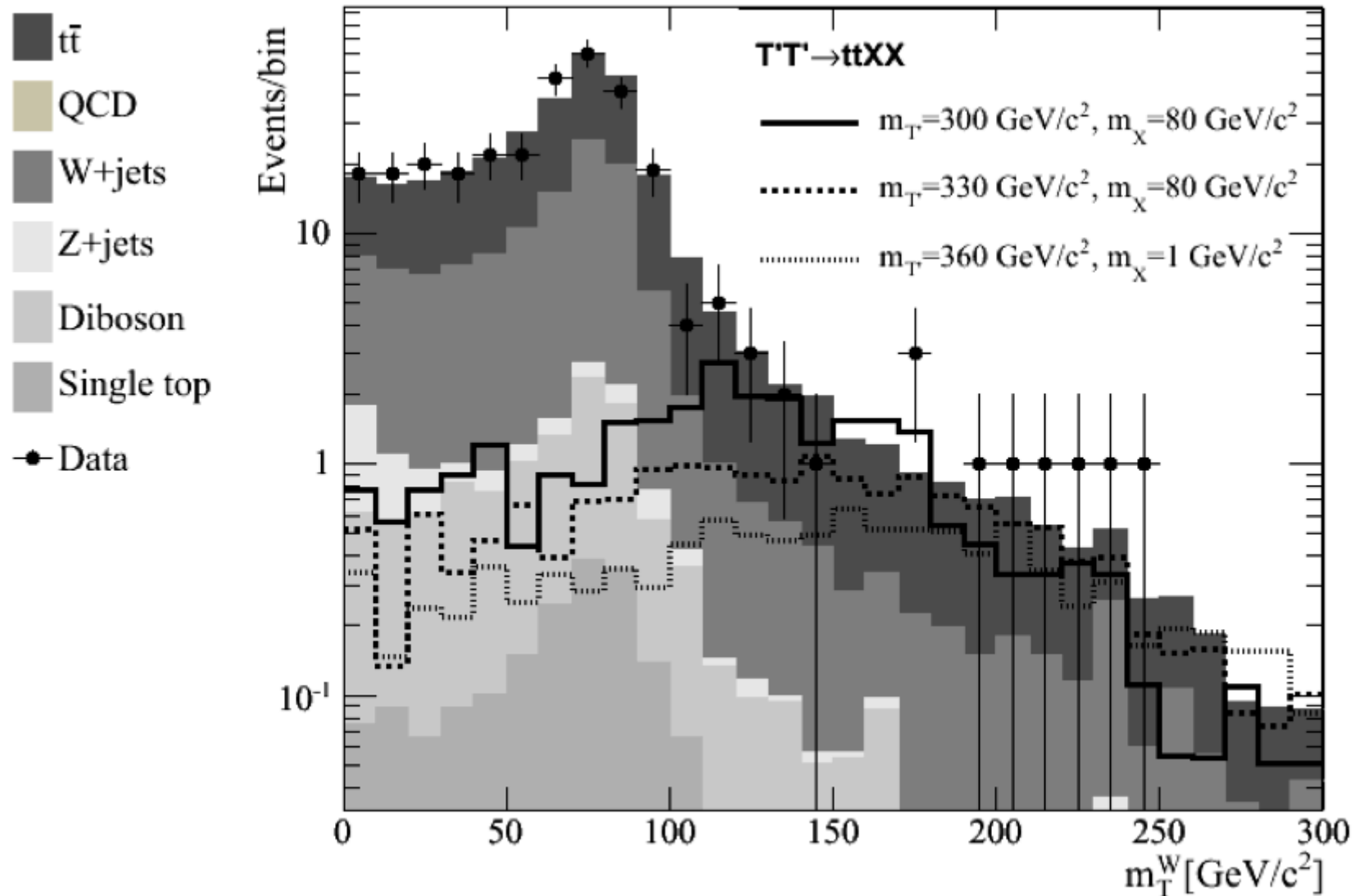


Look for  $t\bar{t}$  + invisible  $X$

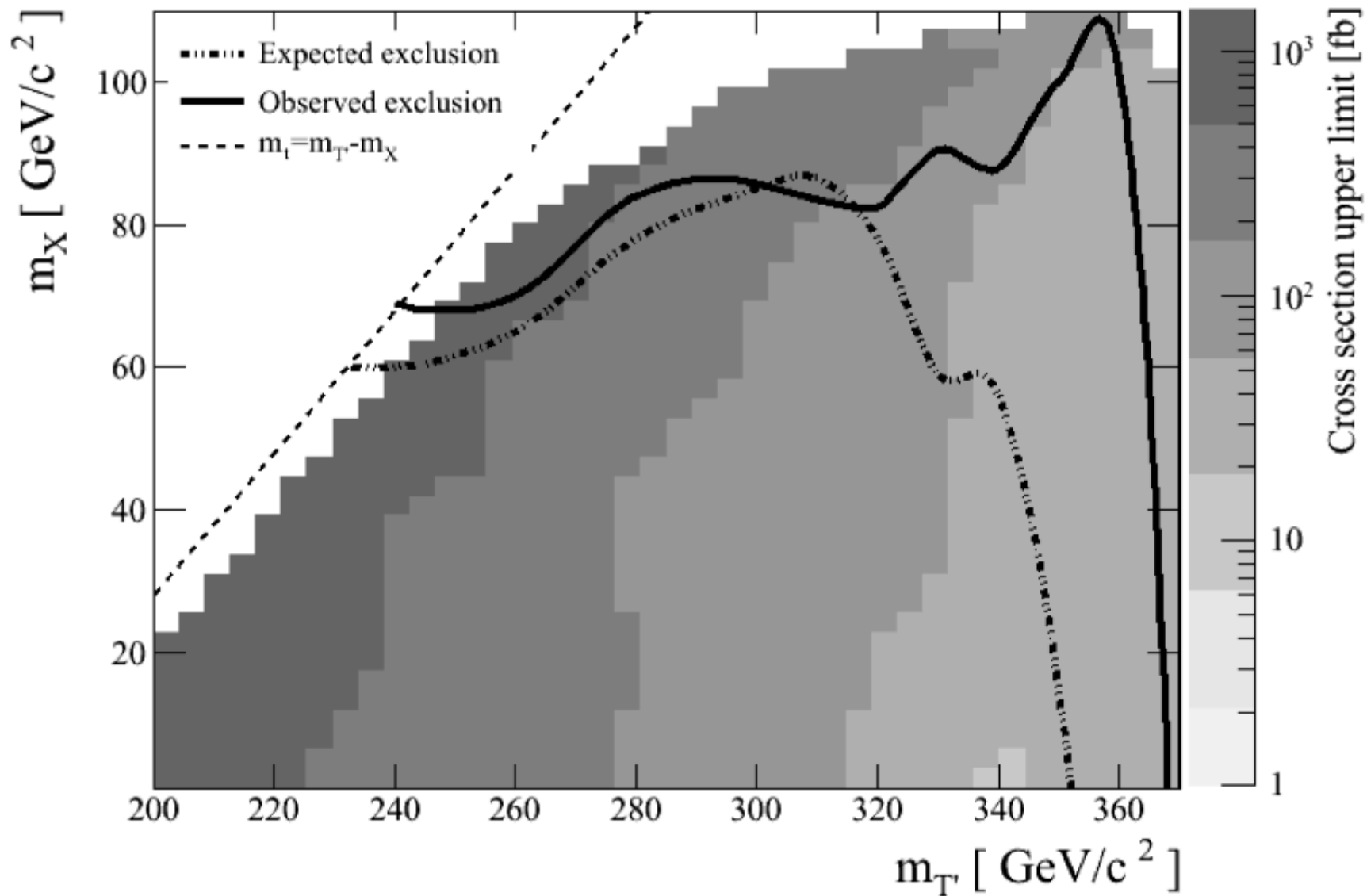
$T' \rightarrow t + X$

stop  $\rightarrow t + \text{LSP}$

# Transverse mass



# Limits



# Outline

I. Motivation

II. Strategy

III. Results

*a. Heavy resonances ( $Z'$ )*

*b. Heavy quarks ( $b'$ ,  $t'$ )*

*c. Simplified SUSY*



# SUSY

## Goal

Set limits on SUSY-like processes  
in as general a fashion as possible

## Approach

Use effective lagrangian, explicitly set particle masses (EW scale):  
simple to handle, easy to interpret

Set limits as functions of these masses, not parameters of specific models:  
can be easily translated into arbitrary models

# How?

How many particles & parameters needed?

Want leptons

*needs Ws and Zs, so **chargino/neutralinos** and **sleptons***

Want strong production

*so **squarks** and **gluinos***

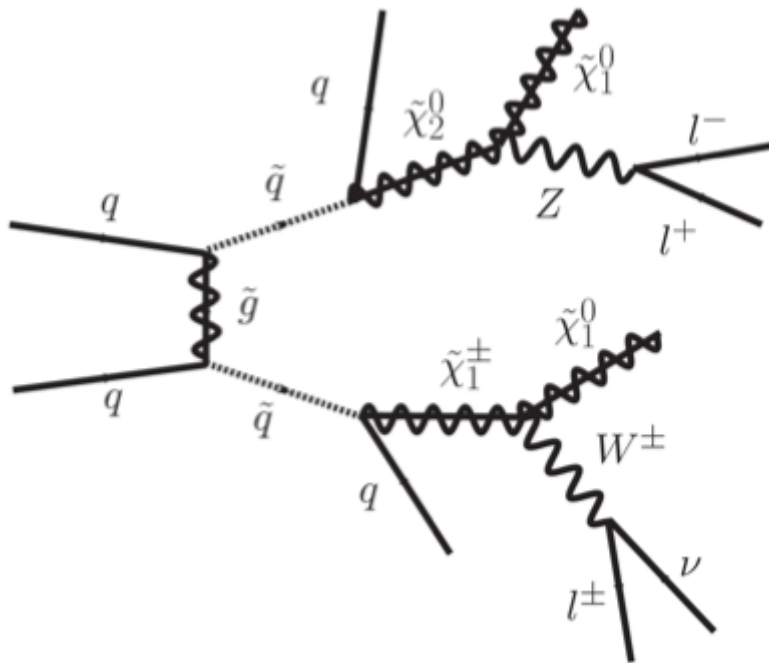
R-Parity conserving

*need **LSP***

Large sections of this space are 3 or 4-dimensional

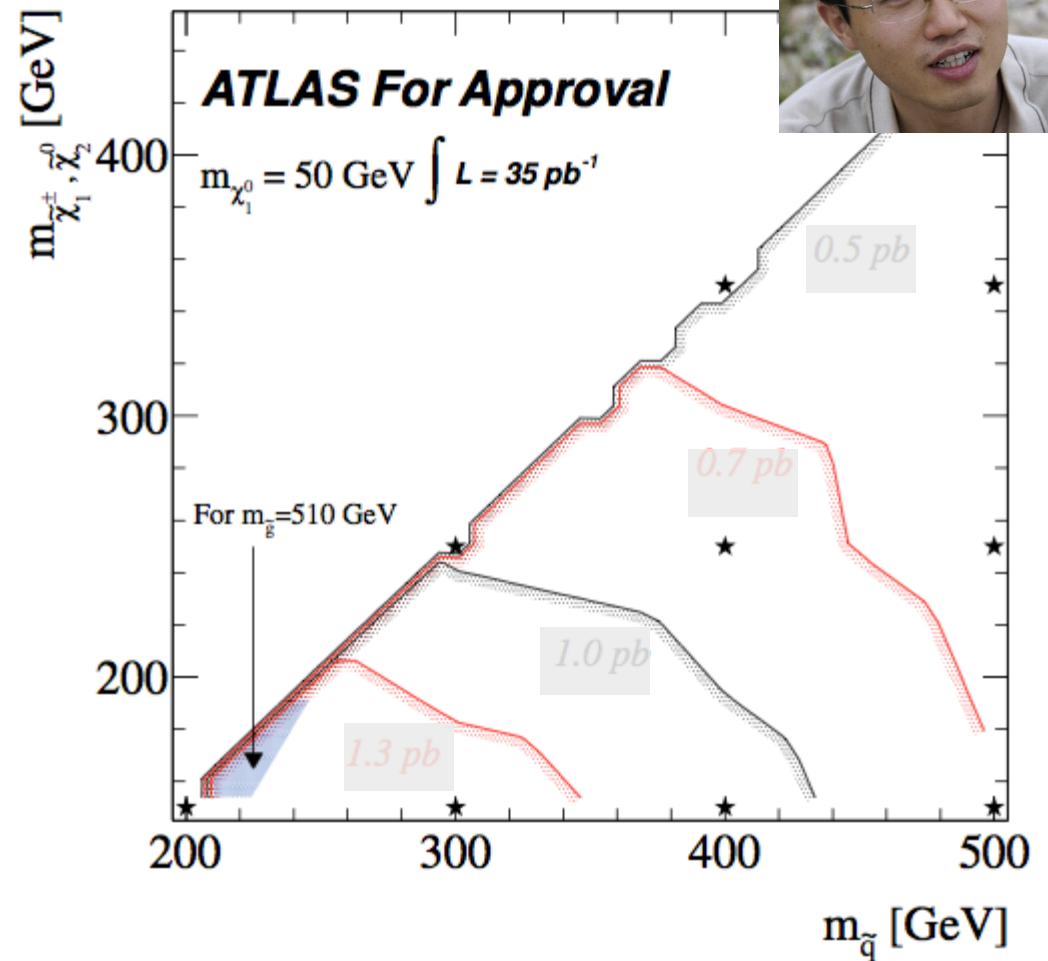
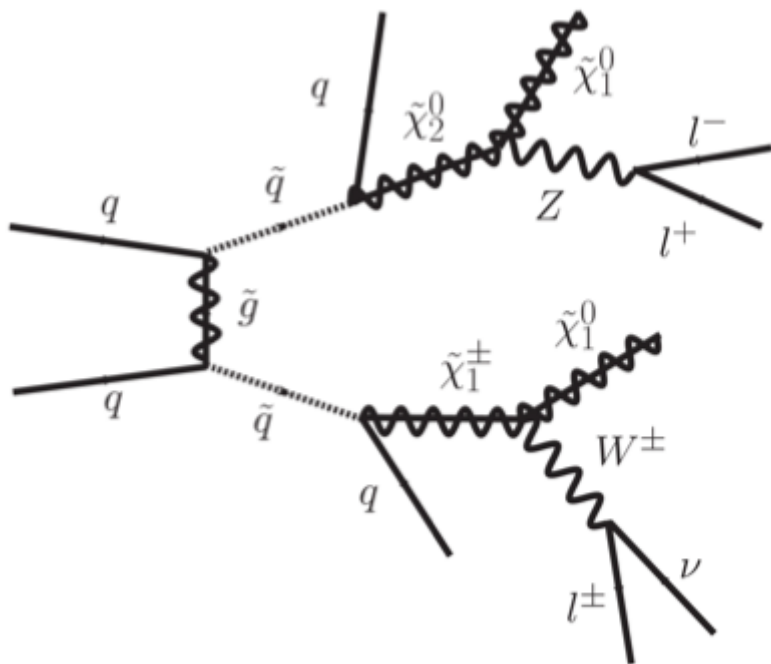
# SUSY simplified

UCI postdoc  
Ning Zhou



# SS SUSY simplified

UCI postdoc  
Ning Zhou



## CDF

Still producing  
world-class physics

## ATLAS

Working well, much  
more to come

