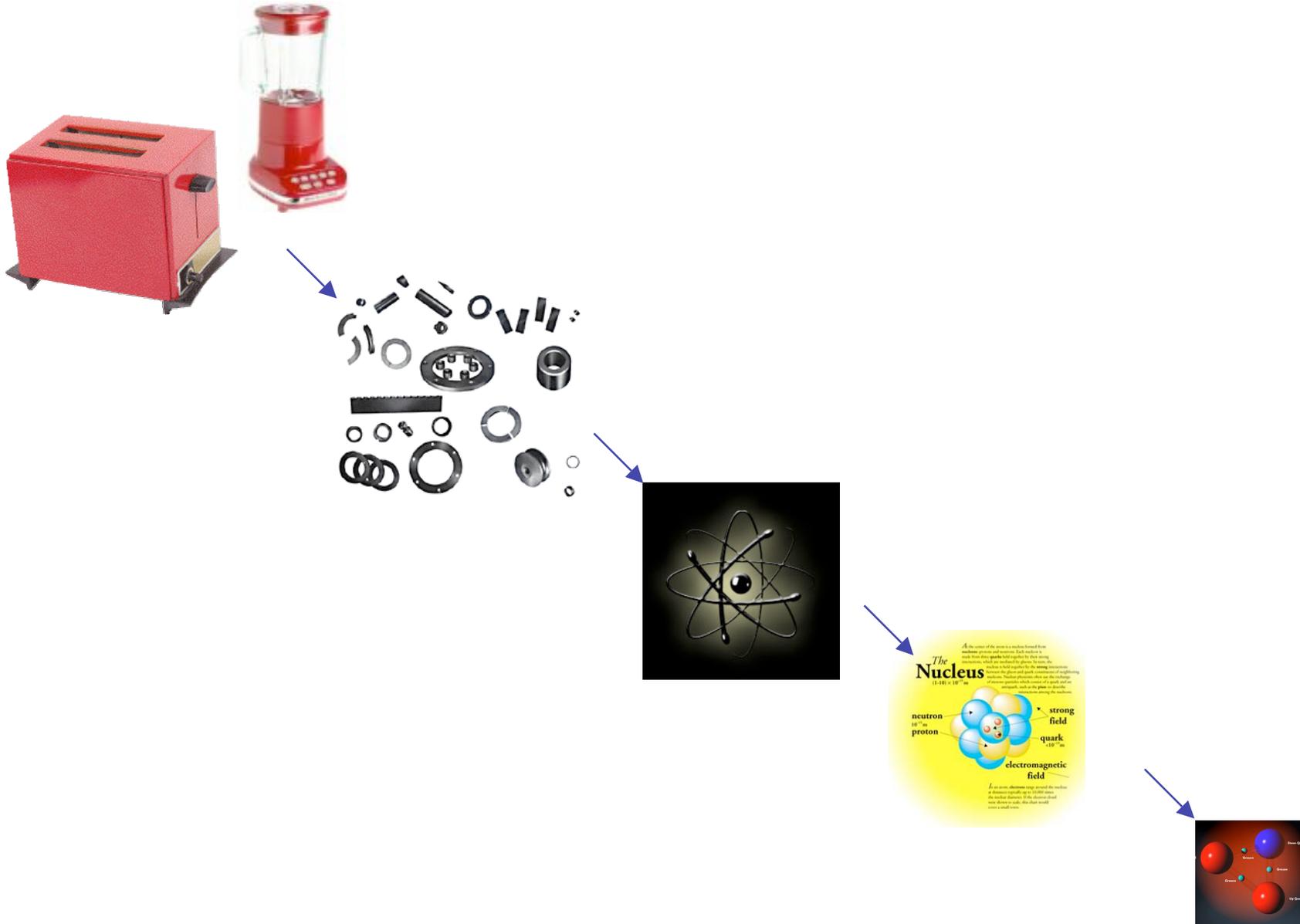


*Precision Measurement
of the Top Quark Mass
in Dilepton Events at CDF*



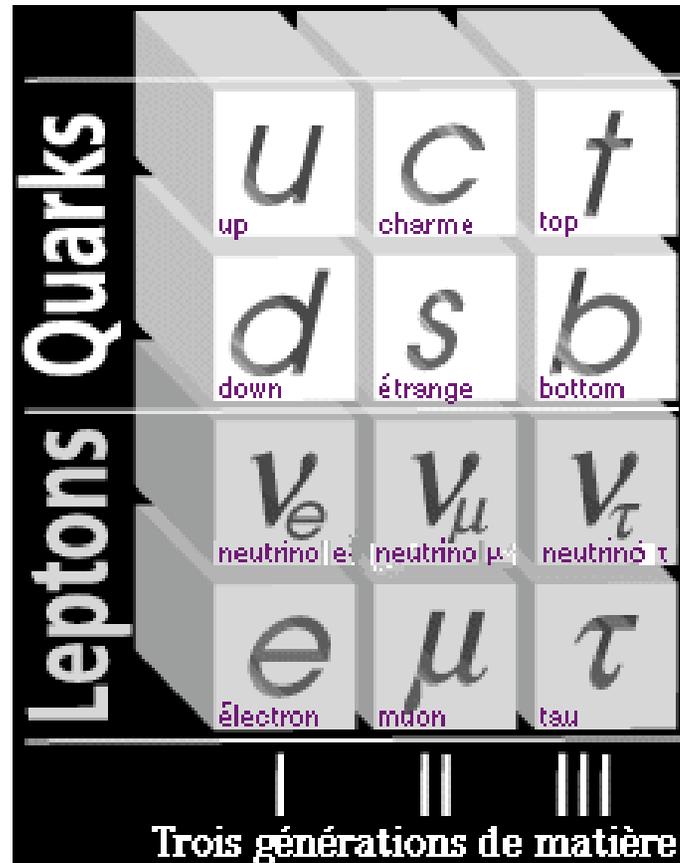
DANIEL WHITESON
UNIV. OF PENNSYLVANIA

Probing the structure of our environment



State of the Art

The Standard Model



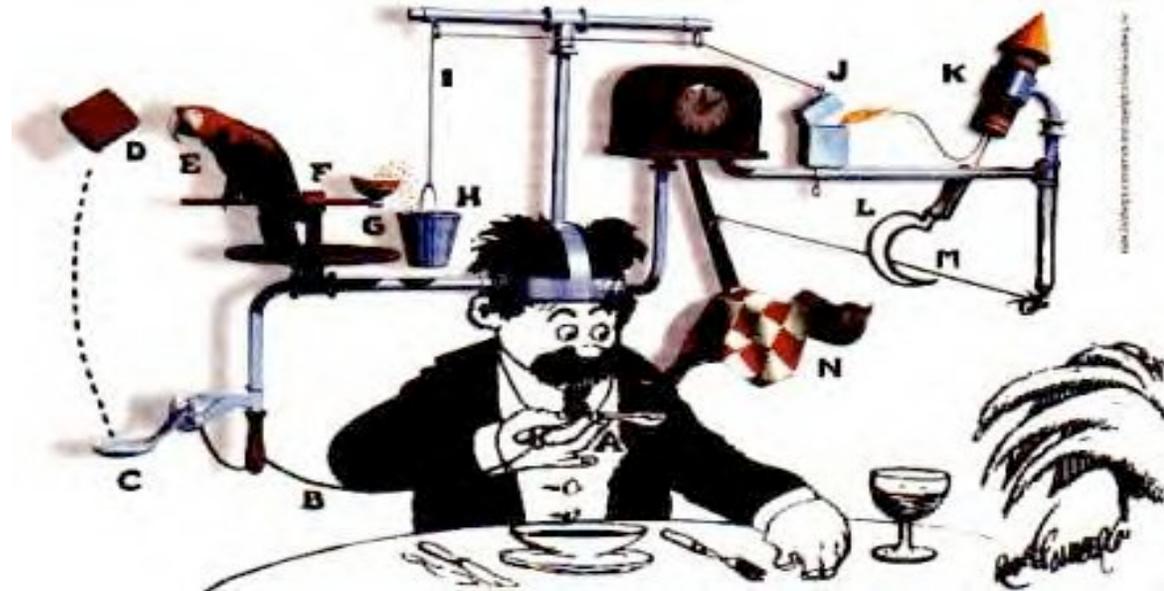
State of the Art

Exercise 1.1.1.1a: Given locality, causality, Lorentz invariance, and known physical data since 1860, show that the Lagrangian describing all observed physical processes (sans gravity) can be written:

$$\begin{aligned}
& -\frac{1}{2}\partial_\nu g_\mu^a \partial_\nu g_\mu^a - g_s f^{abc} \partial_\mu g_\nu^a g_\mu^b g_\nu^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_\mu^b g_\nu^c g_\mu^d g_\nu^e + \\
& \frac{1}{2}ig_s^2(\bar{q}_i^\sigma \gamma^\mu q_j^\sigma)g_\mu^a + \bar{G}^a \partial^2 G^a + g_s f^{abc} \partial_\mu \bar{G}^a G^b g_\mu^c - \partial_\nu W_\mu^+ \partial_\nu W_\mu^- - \\
& M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2c_w^2} M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\mu H - \\
& \frac{1}{2}m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2c_w^2} M \phi^0 \phi^0 - \beta_h \left[\frac{2M^2}{g^2} + \right. \\
& \left. \frac{2M}{g} H + \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right] + \frac{2M^4}{g^2} \alpha_h - igc_w [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
& W_\nu^+ W_\mu^-) - Z_\nu^0 (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - \\
& W_\nu^- \partial_\nu W_\mu^+)] - igs_w [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - \\
& W_\mu^- \partial_\nu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)] - \frac{1}{2}g^2 W_\mu^+ W_\mu^- W_\nu^+ W_\nu^- + \\
& \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^- - Z_\mu^0 Z_\nu^0 W_\nu^+ W_\nu^-) + \\
& g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\mu W_\nu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - \\
& W_\nu^+ W_\mu^-) - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-] - g\alpha [H^3 + H\phi^0 \phi^0 + 2H\phi^+ \phi^-] - \\
& \frac{1}{8}g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2] - \\
& gMW_\mu^+ W_\mu^- H - \frac{1}{2}g \frac{M}{c_w^2} Z_\mu^0 Z_\mu^0 H - \frac{1}{2}ig[W_\mu^+ (\phi^0 \partial_\mu \phi^- - \phi^- \partial_\mu \phi^0) - \\
& W_\mu^- (\phi^0 \partial_\mu \phi^+ - \phi^+ \partial_\mu \phi^0)] + \frac{1}{2}g[W_\mu^+ (H \partial_\mu \phi^- - \phi^- \partial_\mu H) - W_\mu^- (H \partial_\mu \phi^+ - \\
& \phi^+ \partial_\mu H)] + \frac{1}{2}g \frac{1}{c_w} (Z_\mu^0 (H \partial_\mu \phi^0 - \phi^0 \partial_\mu H) - ig \frac{s_w^2}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \dots
\end{aligned}$$

The Standard Model

The theoretical machinery is very effective and often beautiful...



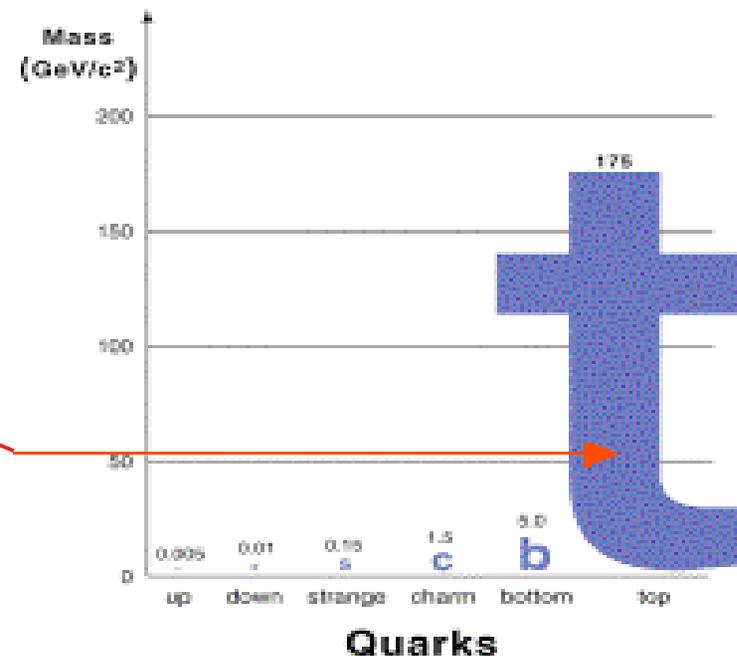
...but its rather unlikely to be the simplest description.

Top Quark

What are the clues?

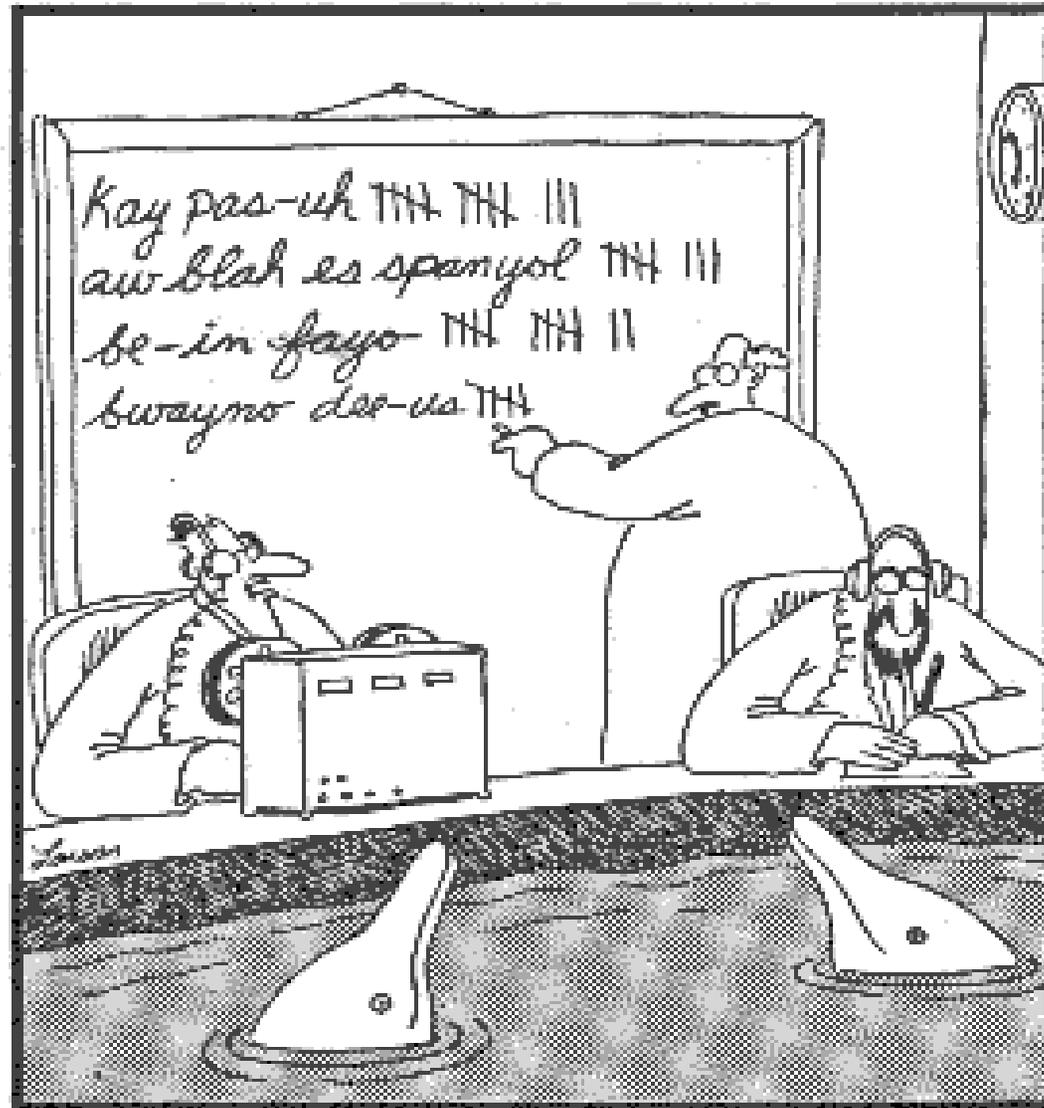
*Where does this mass structure come from?
What does it reveal?
Why is the top so heavy?
Why is its Yukawa coupling ~ 1 ?*

QUARK MASSES



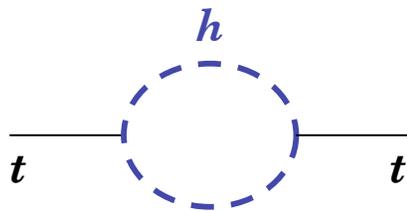
New Theories?

Is there some
simple but
foreign idea
which could
explain many of
these mysteries?



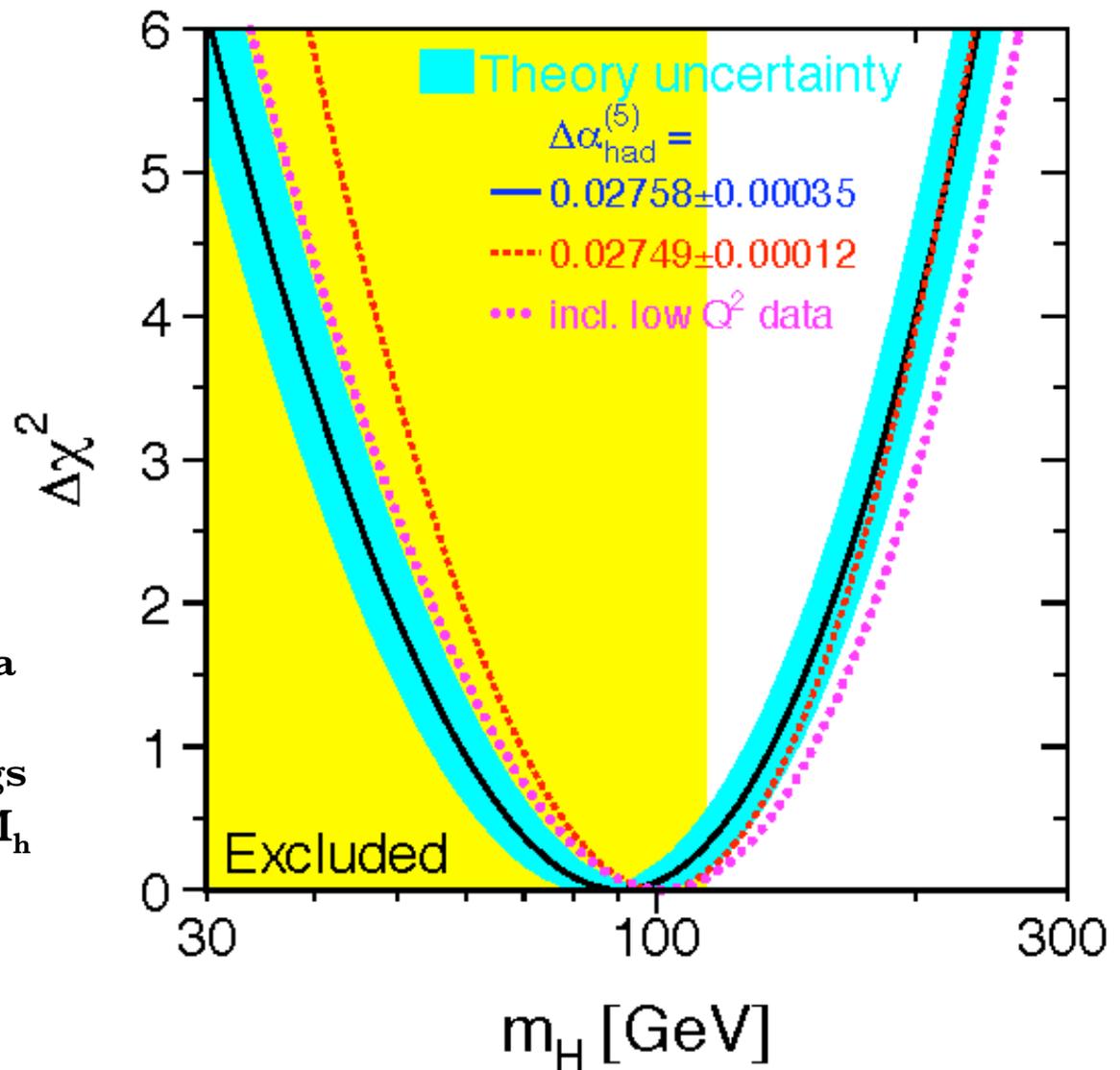
“Matthews... we’re getting another one of those strange ‘aw blah es span yol’ sounds”

Higgs



Higgs connection

Radiative correction to M_t via
Higgs loop
Heavy top means heavy Higgs
 M_t provides constraints on M_h

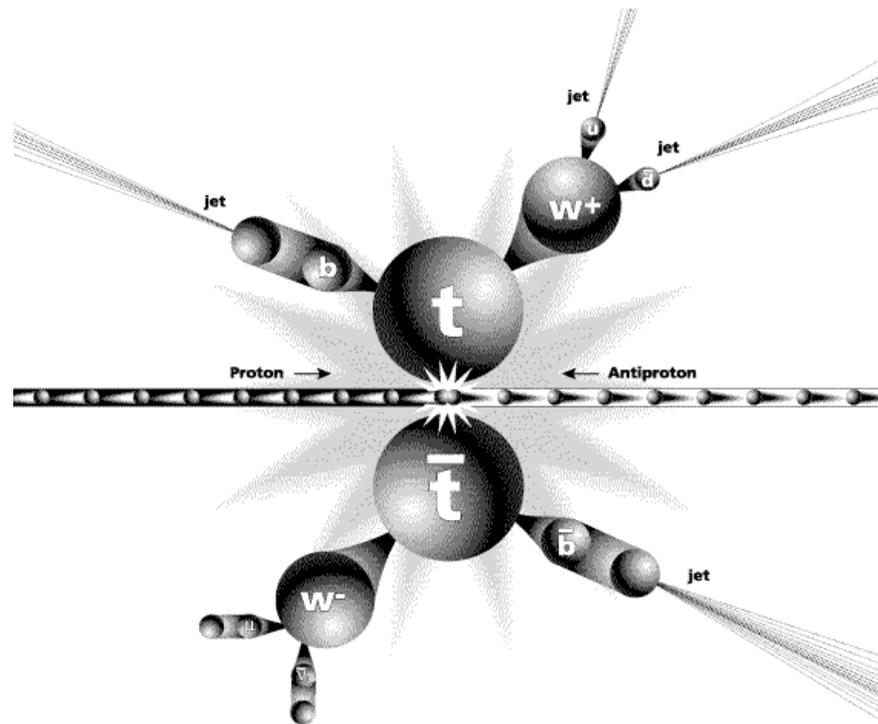


Outline

- Top quark pair production
 - Decay
 - Detector signature
- Dilepton mass measurement problem
 - Template strategy
 - Matrix-element strategy
- Measurements
 - Current results
 - Combination of measurements

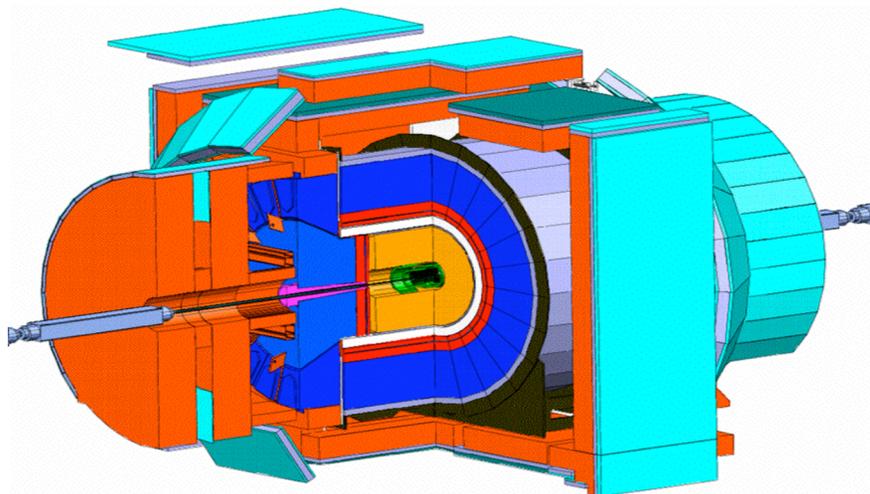
Top Quark Production

The FermiLab Tevatron is the only facility with enough energy to produce and observe top quarks



CDF at Run2

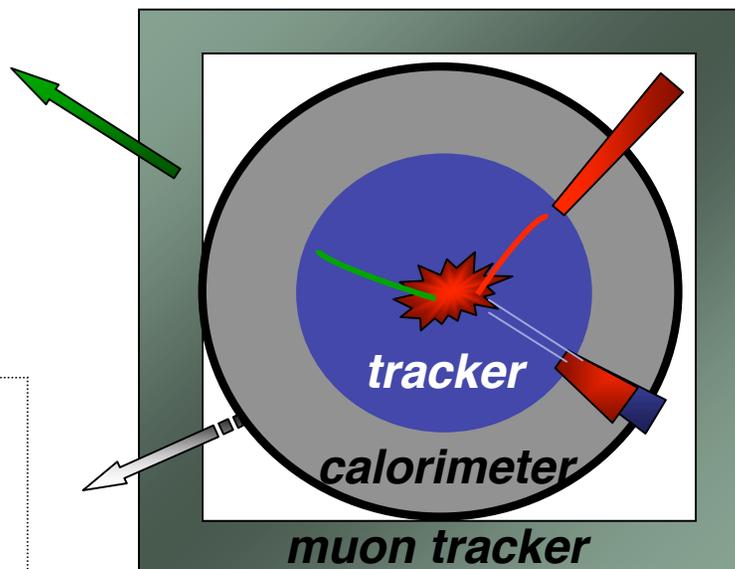
CDF



Cartoon Detector

muon
Muon track
Central track

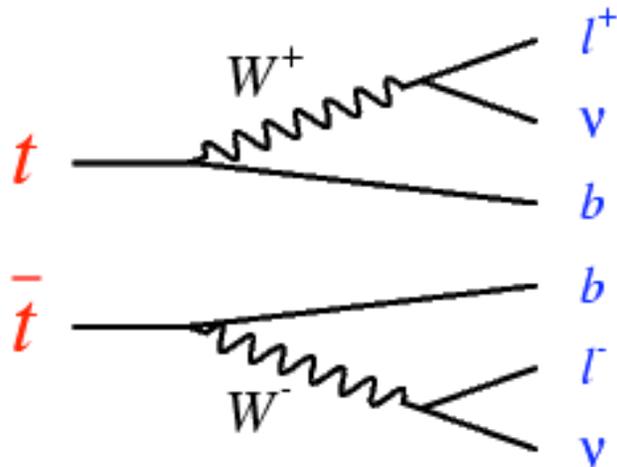
E_T
Neutrinos,
new physics



electron
EM shower
Central track

Jets
EM and hadronic
showers

Dilepton Signature



Dileptons are clean

More leptons!

Fewer jets

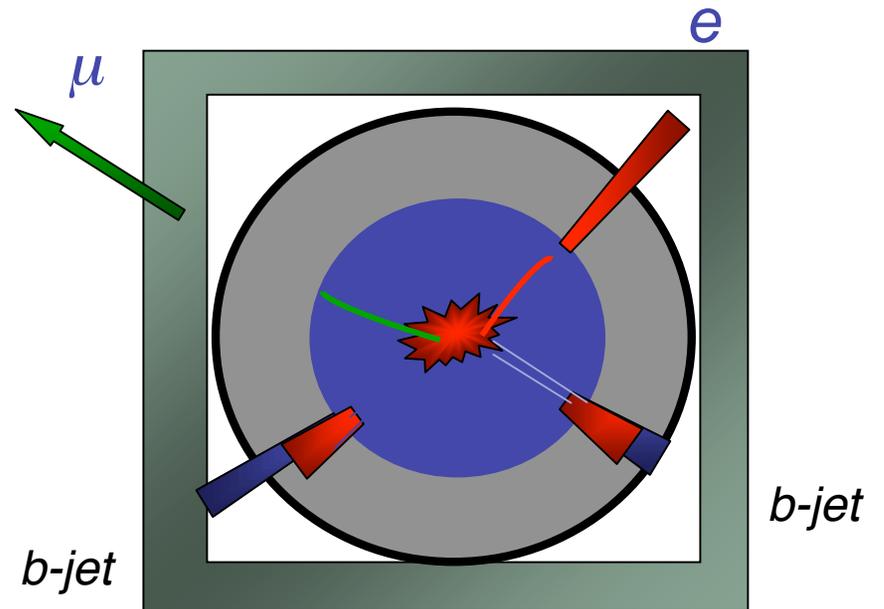
Smaller backgrounds

Small statistics

$bb+ee/e\mu/\mu\mu$: 5%

$bb+e/\mu+qq$: 30%

$bb+qq+qq$: 44%



Particle ID and measurement

Charged leptons are well-measured

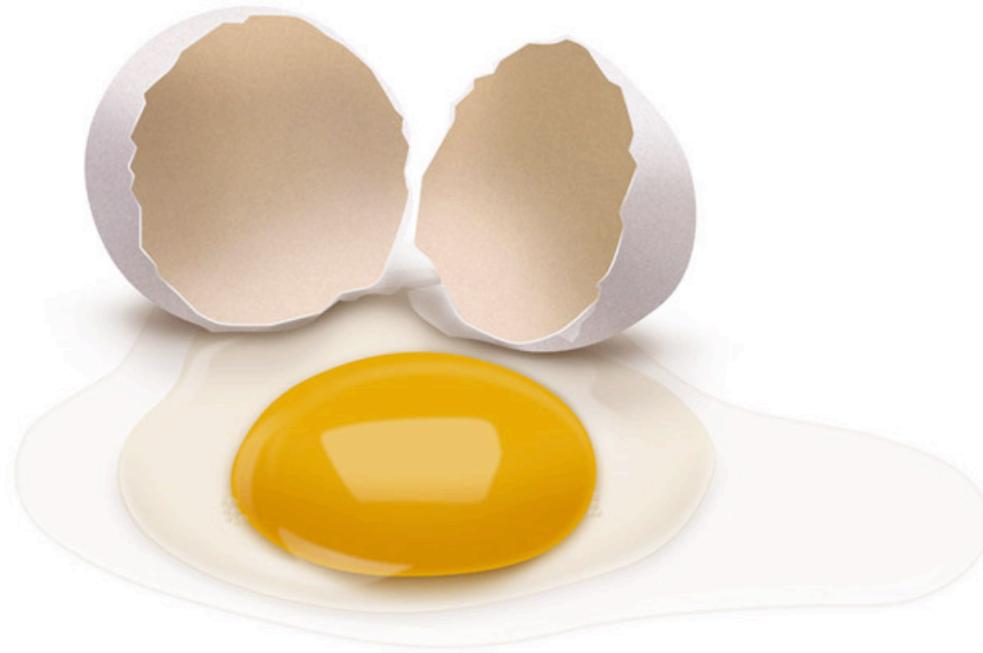
Jets measure quark energies

Neutrinos are undetected

How to measure mass

How do we know how much mass the top had?

We can't just put the pieces back together again



Information has been lost

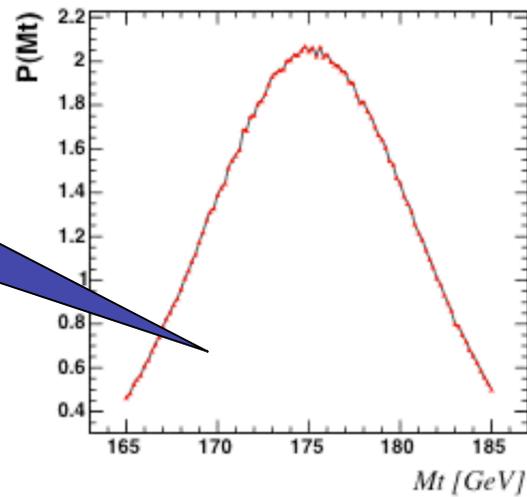
-Neutrinos have escaped undetected

-B-quarks have hadronized, showered, been clustered into *jets*

Solution

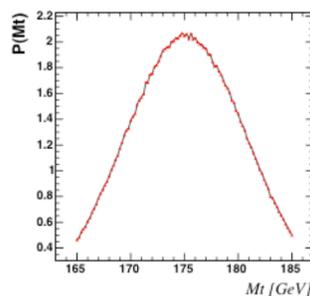
Integrate over unknown quantities.

*Each event has
a curve,
rather than
a single
mass value*



Construct probability curve.

Probability calculation



Template strategy

Construct parameterized probability; fit parameters to examples from simulated events

This sweeps difficult effects into the parameters.

Matrix-element strategy

Calculate probability directly from matrix-element

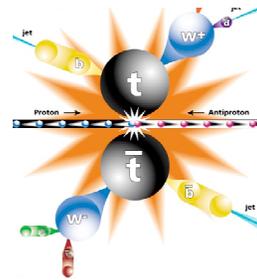
Make simplifying approximations

Measure any necessary corrections using simulated events.

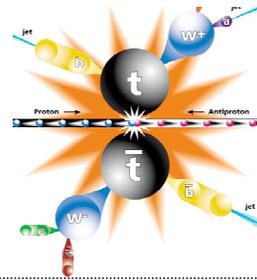
Template Method

$$P(\textit{partons} \mid M_t) \propto P(\textit{measurement} \mid \textit{partons}) = P(\textit{measurement} \mid M_t)$$

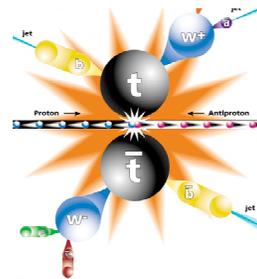
$M_t = 165$



$M_t = 175$



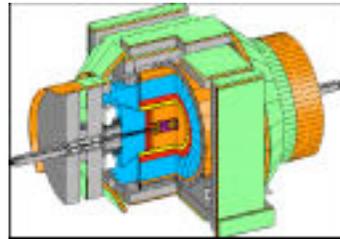
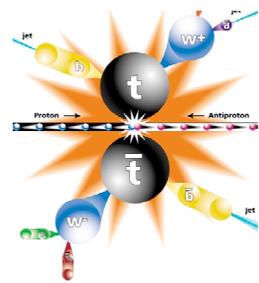
$M_t = 185$



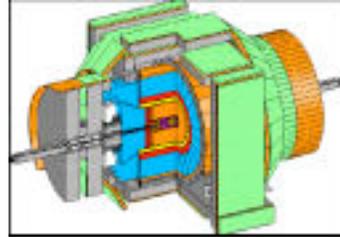
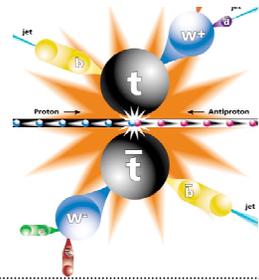
Template Method

$$P(\textit{partons} \mid M_t) \propto P(\textit{measurement} \mid \textit{partons}) = P(\textit{measurement} \mid M_t)$$

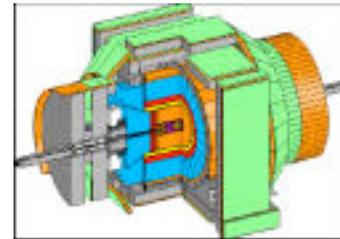
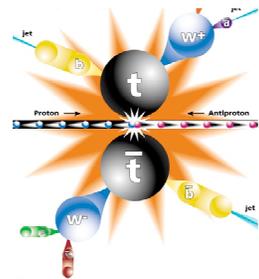
$M_t = 165$



$M_t = 175$



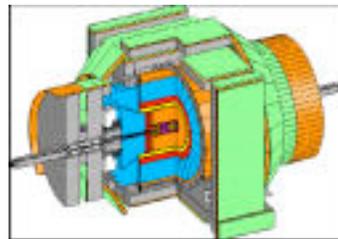
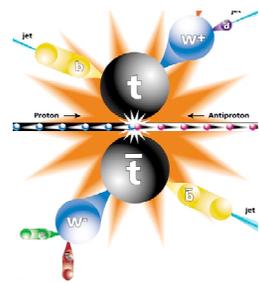
$M_t = 185$



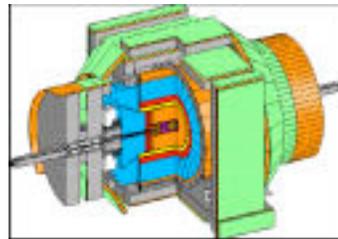
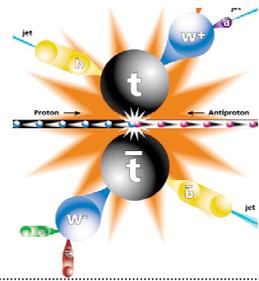
Template Method

$$P(\text{partons} \mid M_t) \propto P(\text{measurement} \mid \text{partons}) = P(\text{measurement} \mid M_t)$$

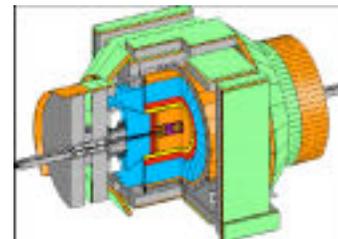
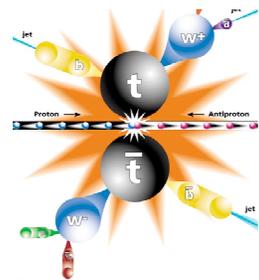
$M_t = 165$



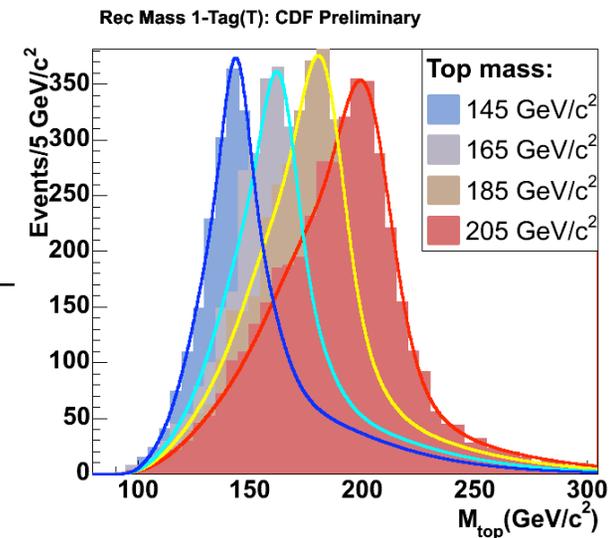
$M_t = 175$



$M_t = 185$



Given examples,
deduce rule for P .

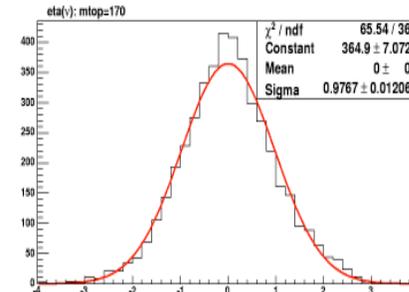


Apply rule to data,
extract mass.

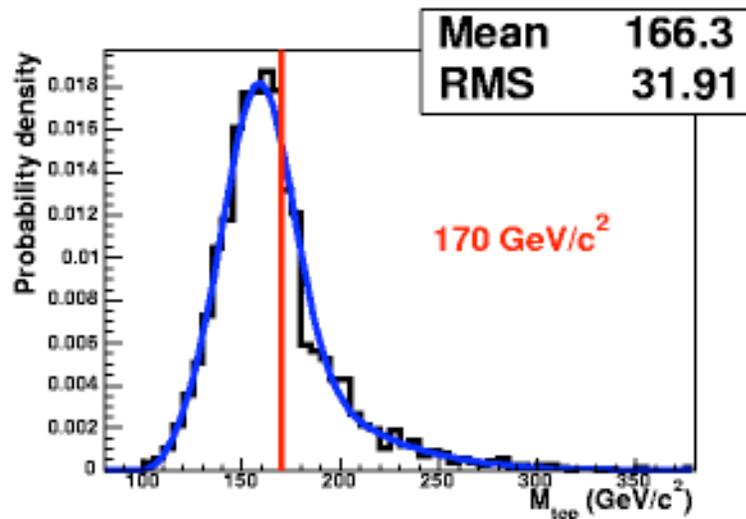
Neutrino Weighting

Inspired by Run1 neutrino weighting analyses

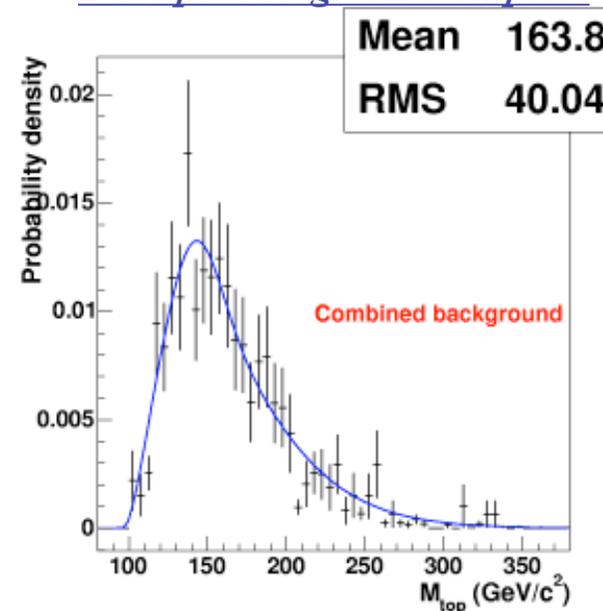
- Constrains M_w , $M_t = M_t$
- Integrates over 2 neutrino eta distributions
- At specific M_t , weight factor derived from missing energy resolution



Example Signal template



Example background template



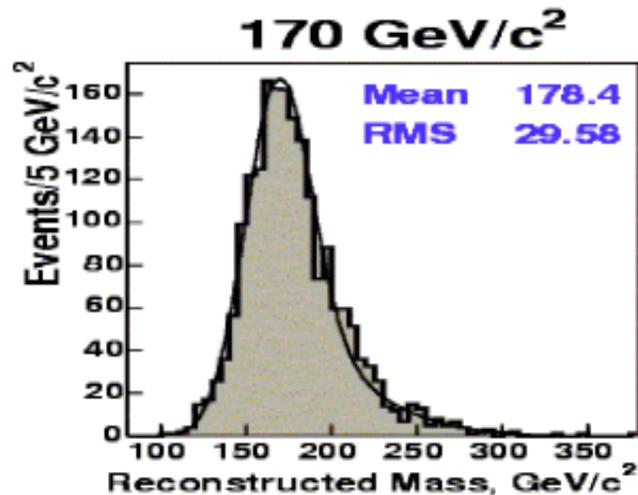
Neutrino phi

Neutrino angle method

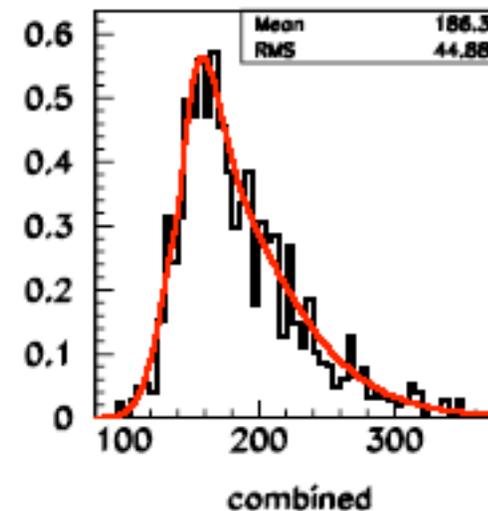
- Integration over two neutrino angles; flat distribution requires no MC prior
- Constrain $M_t = M_t$
- Find most probable top mass which minimizes:

$$\chi^2 = \sum_{i=1,2} \frac{(P_T^i - \tilde{P}_T^i)^2}{\sigma_i^2} + \sum_{j=1,2} \frac{(P_T^j - \tilde{P}_T^j)^2}{\sigma_j^2} + \sum_{i=x,y} \frac{(UE^i - \tilde{UE}^i)^2}{UE^{i^2}} + \frac{(M_{l\nu 1} - M_W)^2}{\Gamma_{M_W}^2} + \frac{(M_{l\nu 2} - M_W)^2}{\Gamma_{M_W}^2} + \frac{(M_{j1l\nu 1} - \tilde{m}_t)^2}{\Gamma_{M_t}^2} + \frac{(M_{j2l\nu 2} - \tilde{m}_t)^2}{\Gamma_{M_t}^2}$$

Example Signal template



Example background template



$t\bar{t} p_z$

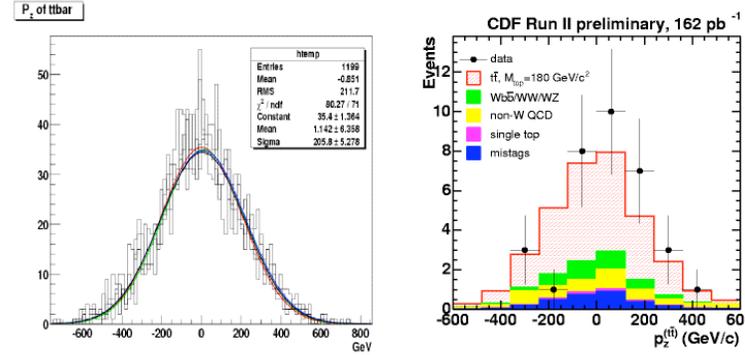
Method using $t\bar{t} p_z$

Assume $M_w, M_t = M_t$

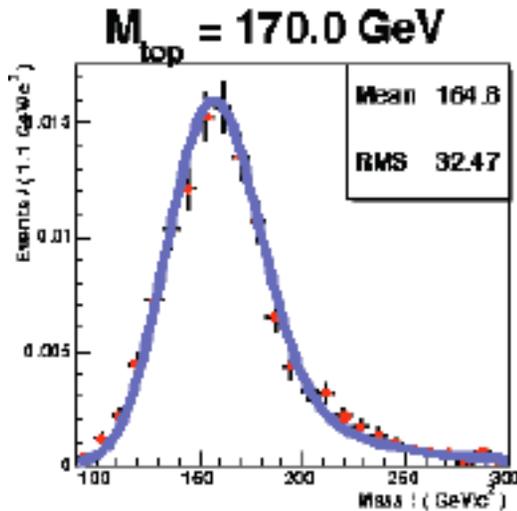
Scan over P_z of $t\bar{t}$ system, prior from MC

Smear jet E_t , $t\bar{t}$ P_t , and missing E_t

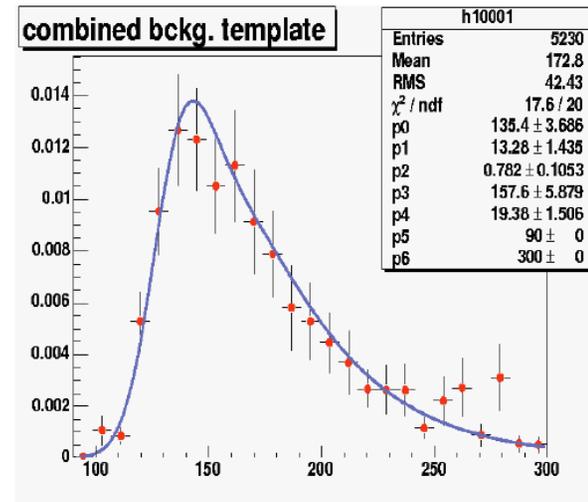
Choose most probable M_t



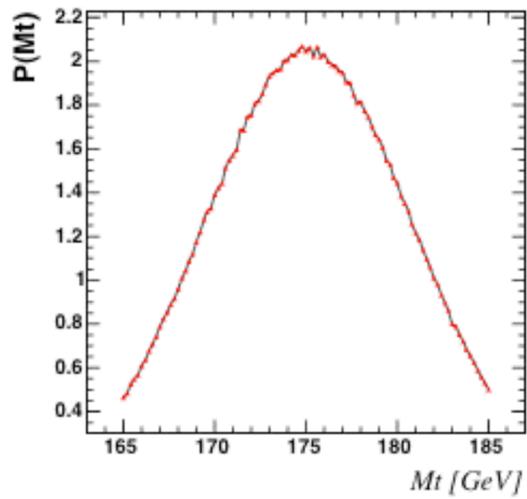
Example Signal template



Example background template

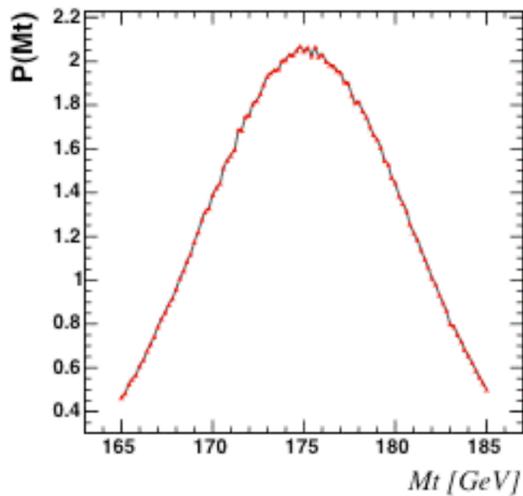


Matrix Element Method



Method

$$P(\text{event } \mathbf{x} \mid M_t) :$$



$$= \frac{1}{\sigma(M_t)} \frac{d\sigma(M_t)}{d\mathbf{x}}$$

History

First ideas by K. Kondo, *et. al.*

First M_t measurement in Run1 by D0 in **+jets**

Applied in Run2 by CDF and D0 in **+jets**

[J. Phys. Soc. Japan **57** 4136 (1988)]

[Nature **429** 636 (2004)]

[W&C 6/11/04 and 7/22/05]

Direct Calculation

Differential cross-section calculation:

$$\frac{d\sigma(M_t)}{d\mathbf{x}} = \int d\Phi_6 |\mathcal{M}_{t\bar{t}}(p; M_t)|^2 f_{PDF}(q_1) f_{PDF}(q_2)$$

Phase-space
Integral

Matrix
Element

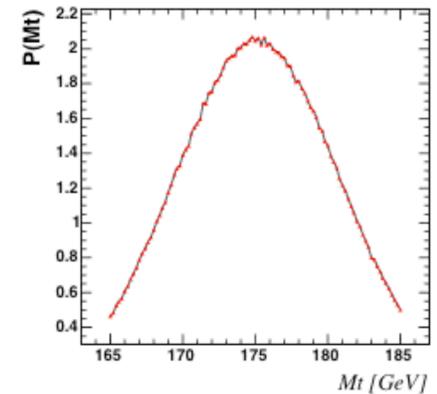
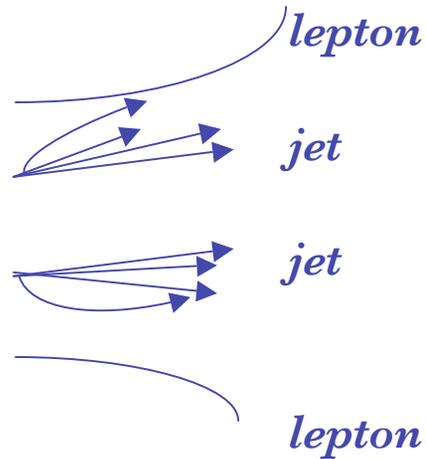
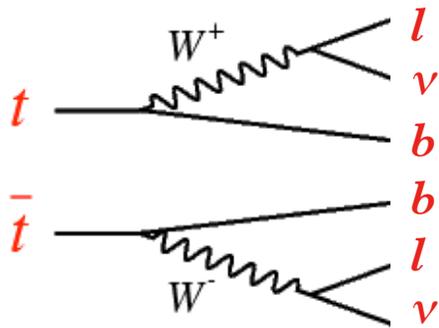
p 4-momentum of final partons
 q 4-momentum of initial partons
 \mathbf{x} measured event variables

But there is no reference to our measured quantities \mathbf{x} !

Method

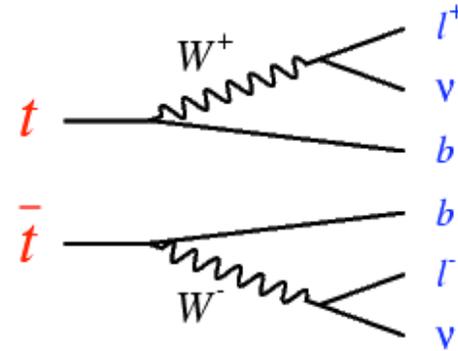
Break into two pieces: parton-level process and showering/resolution effects

$$P(\textit{partons} \mid M_t) \quad \times \quad P(\textit{event } x \mid \textit{partons}) \quad = \quad P(\textit{event } x \mid M_t)$$



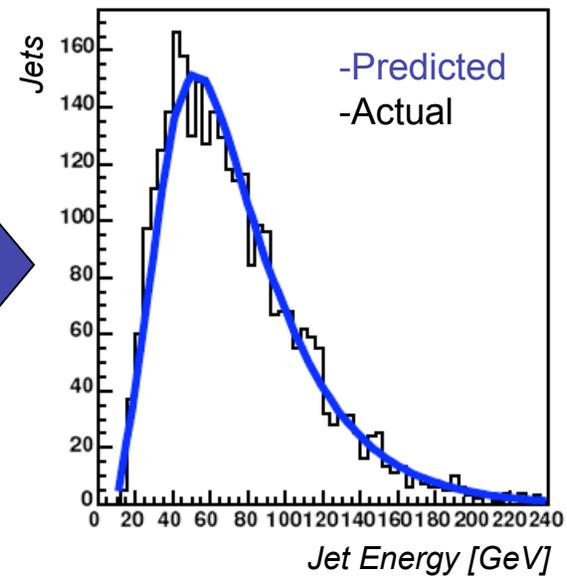
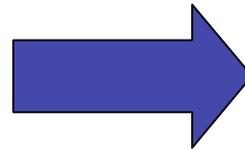
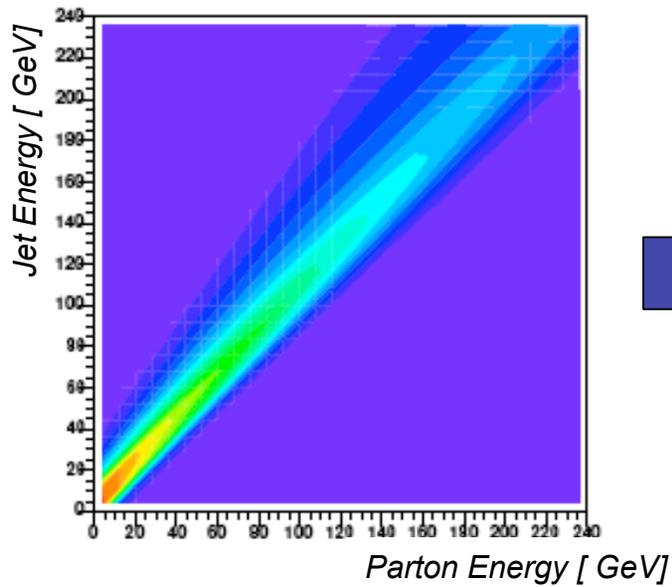
Direct Calculation

We know $P(\text{partons} \mid M_t)$



We can parameterize parts of $P(\text{event } x \mid \text{partons})$:

For jets:



Assumptions

- Initial state
 - No initial state radiation
 - Transverse energy of system is negligible
- Final state
 - Leptons
 - Energy well-measured
 - Direction well-measured
 - Jets
 - Jets arise from b -quarks
 - Direction well-measured
 - Energy can be parameterized from parton energy
- Assumptions make calculation tractable
 - balance sensitivity with computation time

Calculation

p 4-momentum of final partons
 q 4-momentum of initial partons
 \mathbf{x} measured event variables

For each event, calculate differential cross-section:

$$P(\mathbf{x}|M_t) = \frac{1}{N} \int d\Phi_6 |\mathcal{M}_{t\bar{t}}(p; M_t)|^2 \prod_{jets} f(p_i, \mathbf{x}) f_{PDF}(q_1) f_{PDF}(q_2)$$

Phase-space
Integral

Matrix
Element

Transfer
Functions

Only partial information available

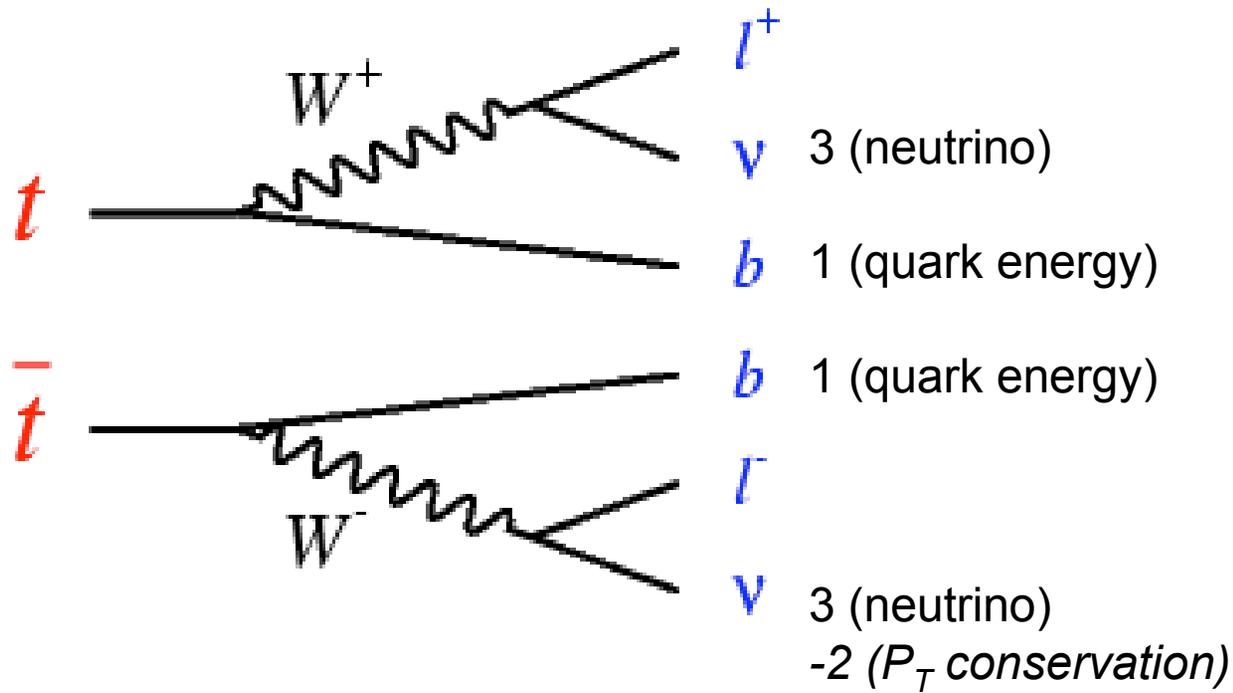
Fix measured quantities

Integrate over unmeasured parton quantities

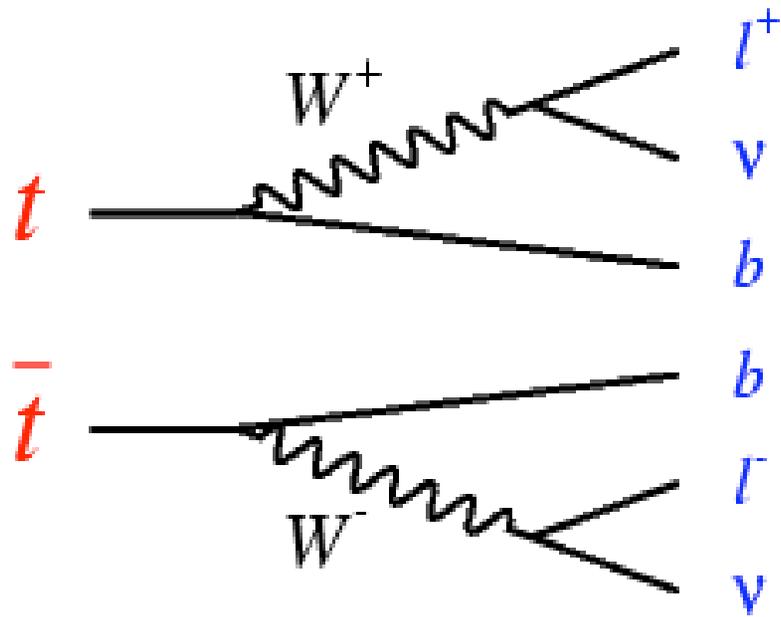
consistent with $t\bar{t}$ production and measured event.

Unknowns

6 unknowns



Integration



6 unknowns

l^+
 ν 3 (neutrino)
 b 1 (quark energy)
 b 1 (quark energy)
 l^-
 ν 3 (neutrino)
 -2 (P_T conservation)

6 integration variables

W, t invariant masses

quark energy

quark energy

W, t invariant masses

Choice of variables

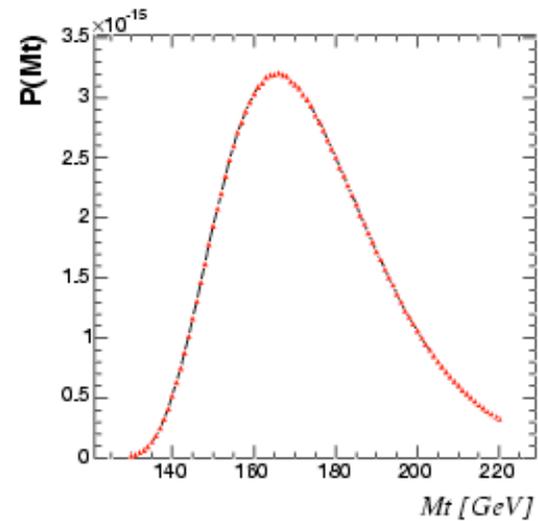
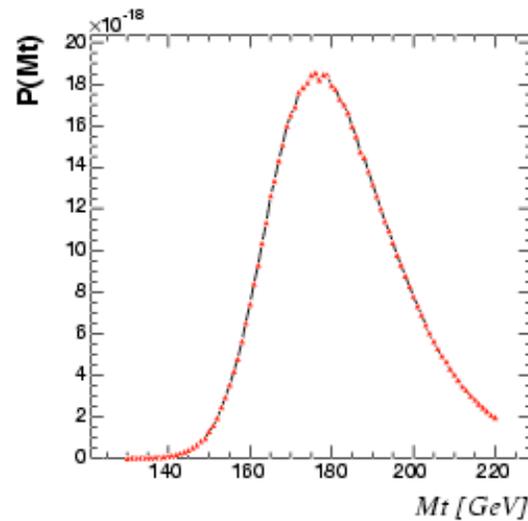
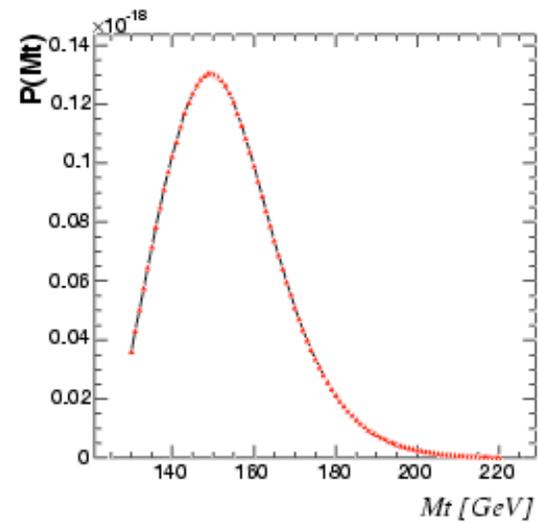
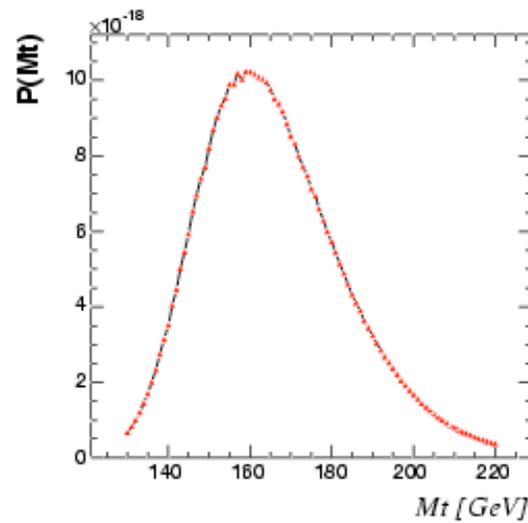
Transform phase space to exchange variables for those which are more efficient.

Requires numerical solution at each integration point.

Integration done with VEGAS.

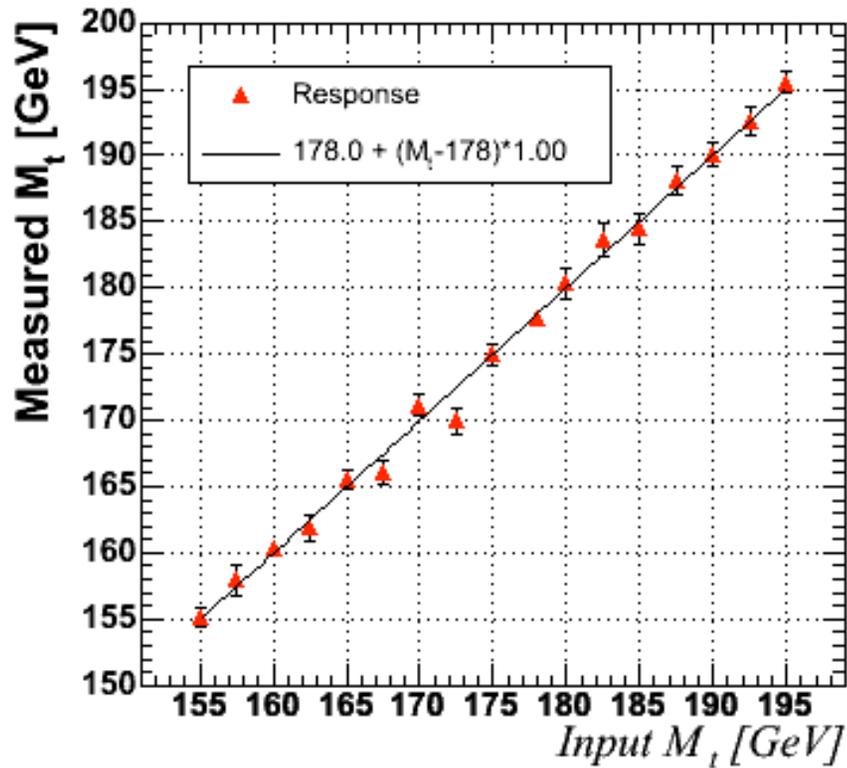
Examples

Four simulated events (Herwig, $M_t=180$ GeV)



To test performance, build joint probability in pseudo-experiments

Signal only results: full simulation



$$\text{Response} = \langle M_{\text{meas}} \rangle$$

$$\text{Pull} = \frac{M_{\text{meas}} - M_{\text{true}}}{\sigma_{\text{meas}}}$$

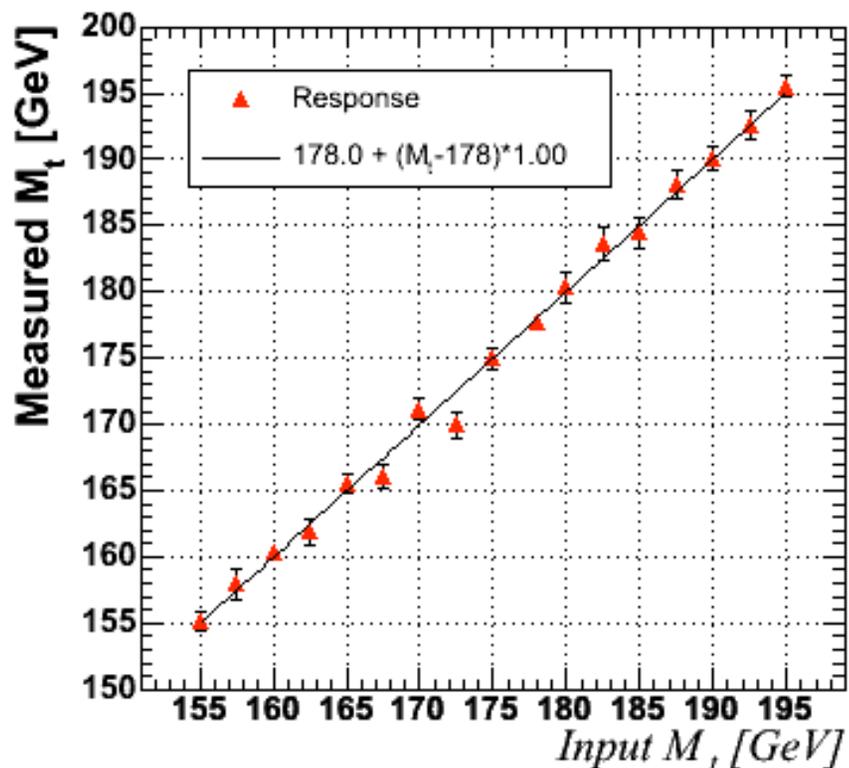
Response

Linear.

Unbiased.

Signal only results: full simulation

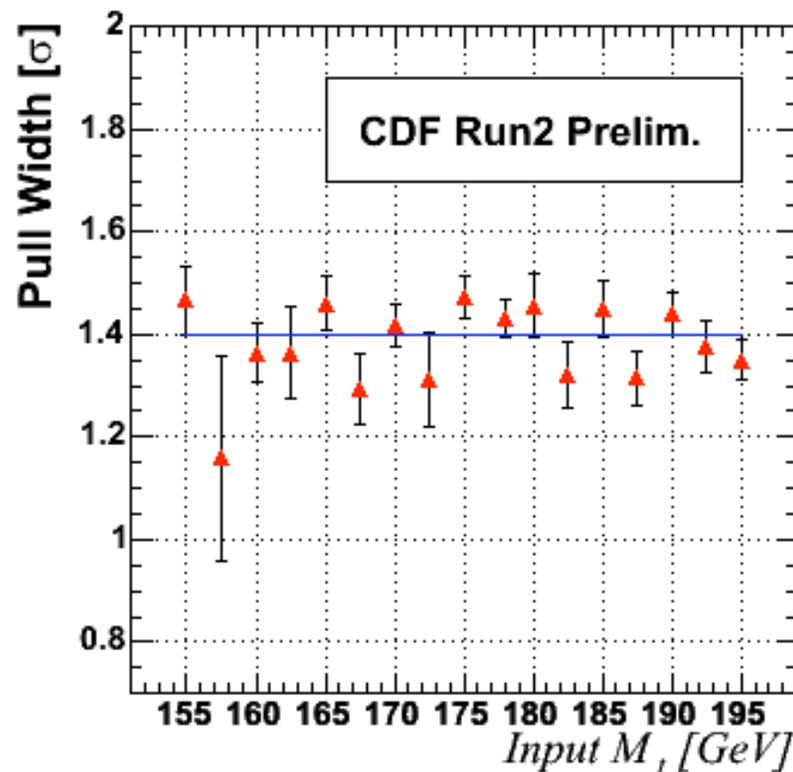
Joint probability in signal-only pseudo-experiments



Response

Linear.

Unbiased.



Pull width

Flat.

Width is ~ 1.4 .

In parton-level tests, width is ~ 1.0 .

(*template methods have $\sim 1.05-1.1$*)

Pull width

PW inflated because probability contains assumptions broken by a subset of events.

Assumptions held

Parton events: $width = 1.0$

Assumptions broken

Full simulation : $width = 1.4$

- + Two matched jets: $width \rightarrow 1.2$
- + Well measured leptons: $width \rightarrow 1.1$
- + Small parton-jet angle: $width \rightarrow 1.0$

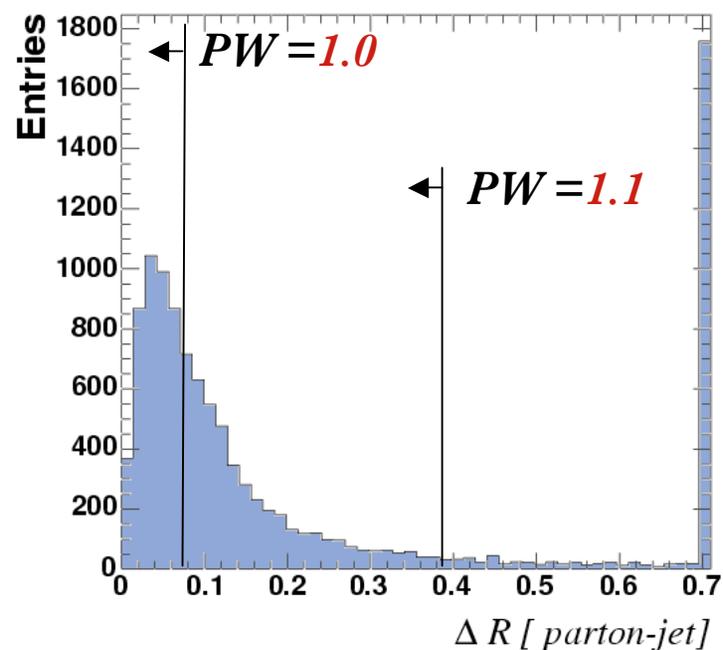
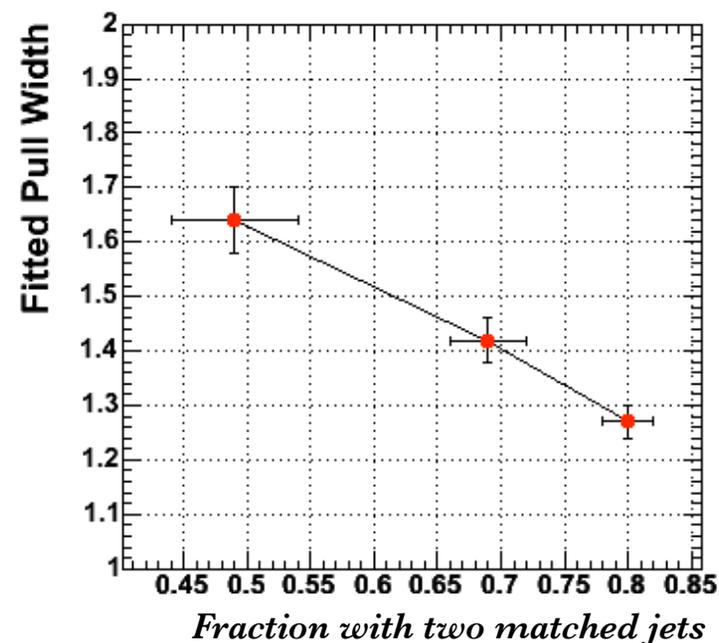
Scale factor for error

Flat in top mass

Flat in measured statistical error

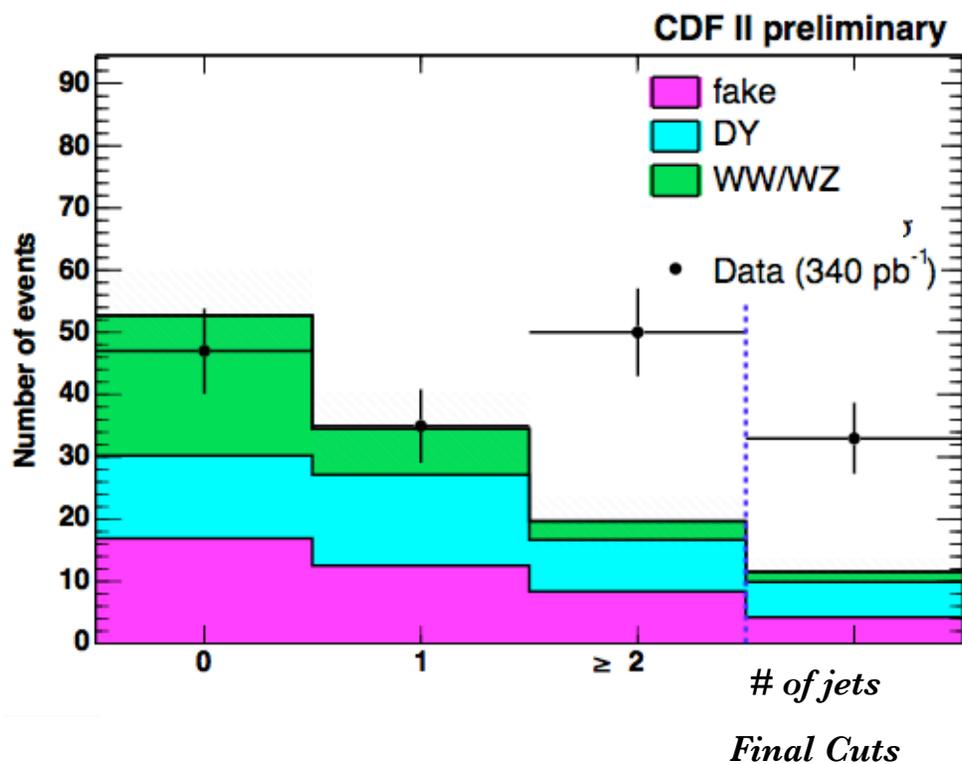
Insensitive to systematic variations

Error on scale factor is ~ 0.03



DIL Sample composition

DIL Sample is expected to be **3:2** signal to background



<i>Sample</i>	N_{events}
$t\bar{t}$ ($M_t=175\text{GeV}$, $\sigma = 6.7 \text{ pb}$)	17.2
$Z \rightarrow e\bar{e}/\mu\bar{\mu}$	4.9
MisID lepton	4.2
WW, WZ	1.6
$Z \rightarrow \tau\bar{\tau}$	0.8
Total	28.7 ± 2.1
Observed	33

Probability expression needs to reflect presence of background

Background Likelihood

We generalize the probability to be a weighted sum of **signal** & **bg** probabilities

$$P(\mathbf{x}|M_t) = P_s(\mathbf{x}|M_t)p_s + P_{bg1}(\mathbf{x})p_{bg1} + P_{bg2}(\mathbf{x})p_{bg2}\dots$$

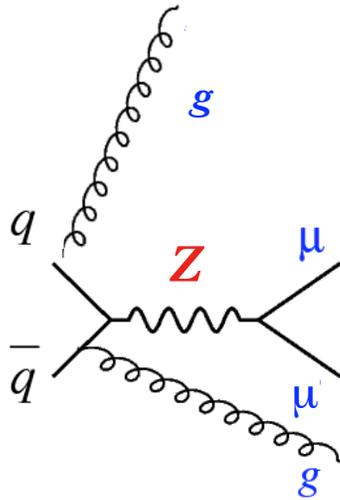
Where the **weights** are the **expected sample fraction**:

$$p_s(M_t) = \frac{\lambda_s(M_t)}{\lambda_s(M_t) + \lambda_b}$$

$$p_b = \frac{\lambda_b}{\lambda_s(M_t) + \lambda_b}$$

$\lambda =$ expected number of events
--

Z+jets



0 unknowns

1 parton energy

1 parton energy

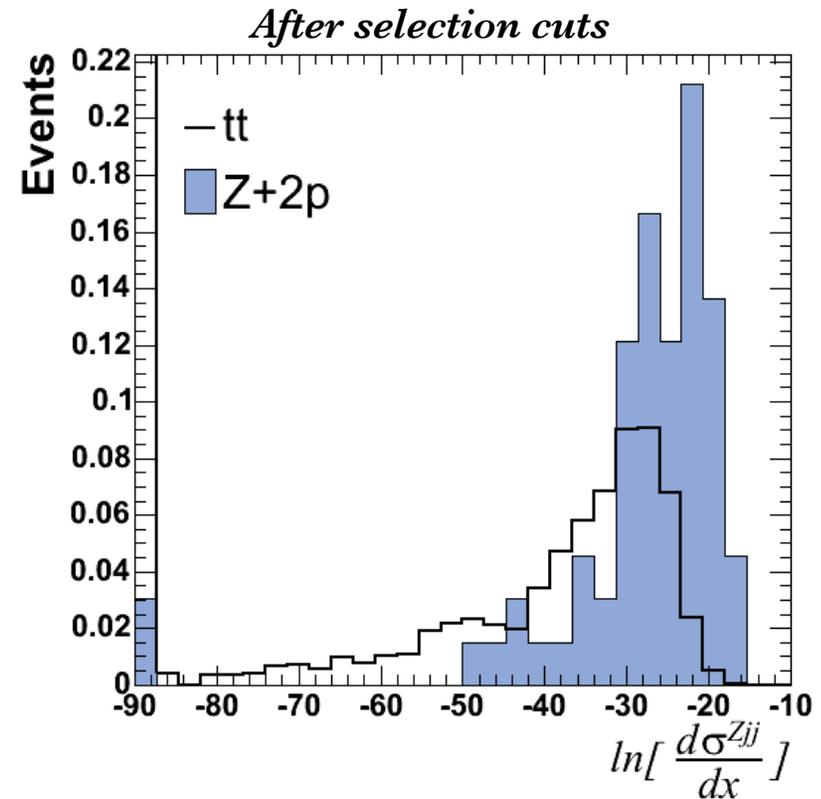
-2 (P_T conservation)

Matrix Element & Integrals

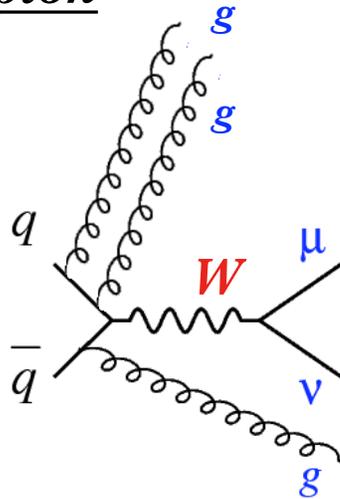
Add 2 integrals for P_T of Zjj system

Alpgen subroutine for Z+2p

Integration with Vegas



Mis-ID Lepton



Matrix Element & Integrals

AlpGen subroutine for W+3p

Integration with Vegas

3 unknowns

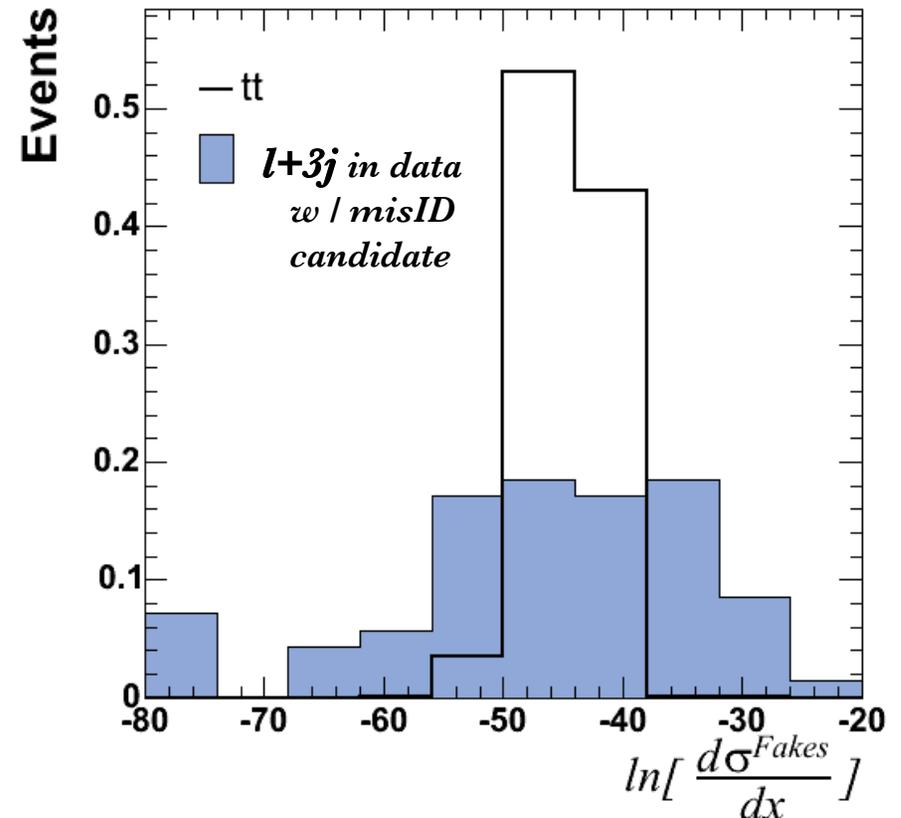
1 parton energy

1 parton energy

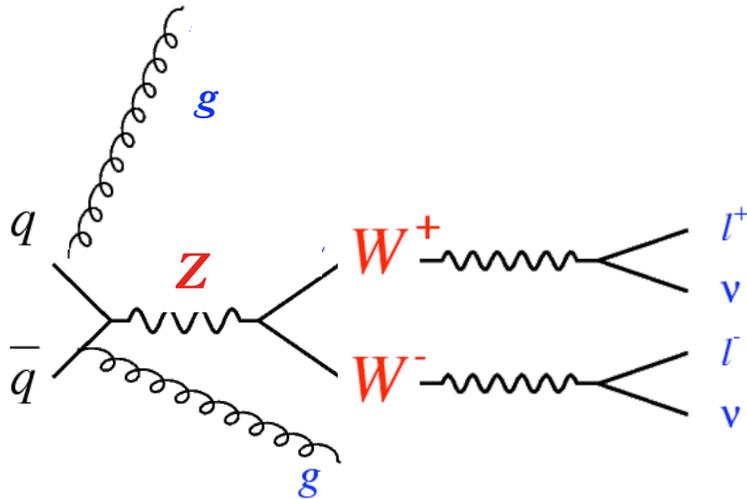
3 neutrino (*P* components)

-2 (*P_T* conservation)

After selection cuts



WW+jets



6 unknowns

Parton energy

Parton energy

3 neutrino (P components)

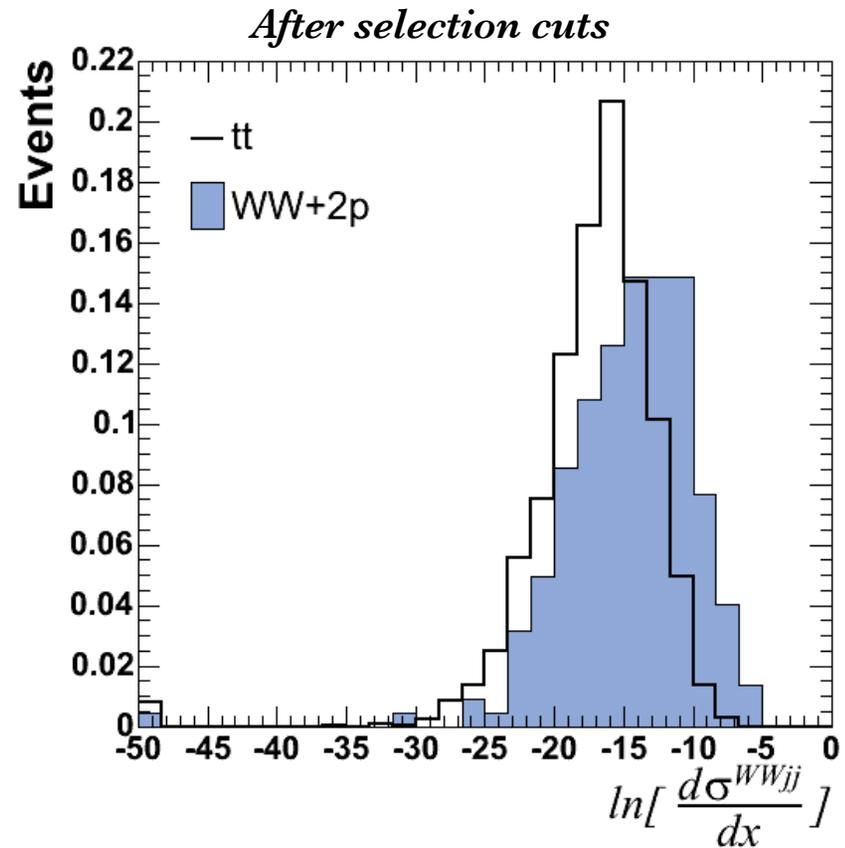
3 neutrino (P components)

-2 (P_T conservation)

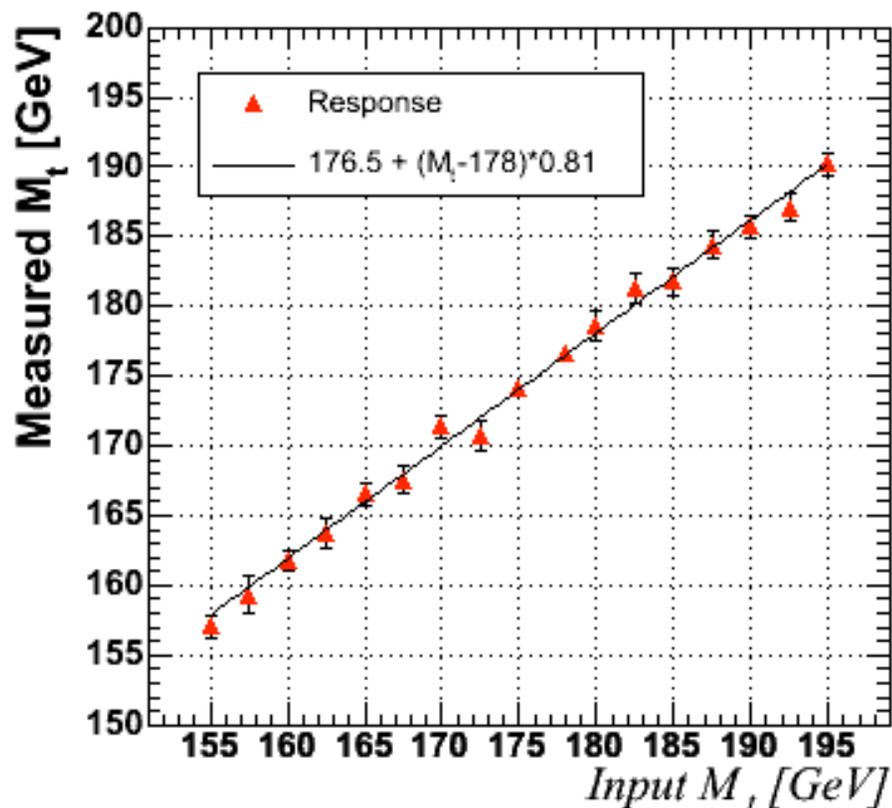
Matrix Element & Integrals

Alphen subroutine for WW+2p

Integration with Vegas



Response & Pulls



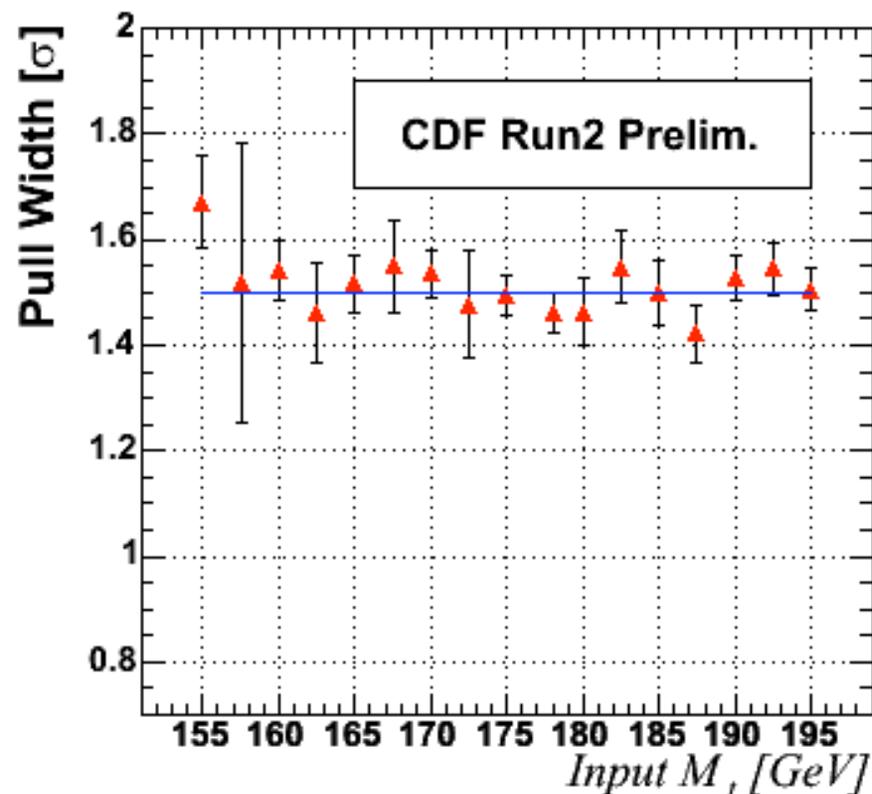
Response

Linear.

Slope < 1 due to backgrounds.

Error is **reduced 15%** by P_{bg}

Slope is **improved by 20%** by P_{bg}

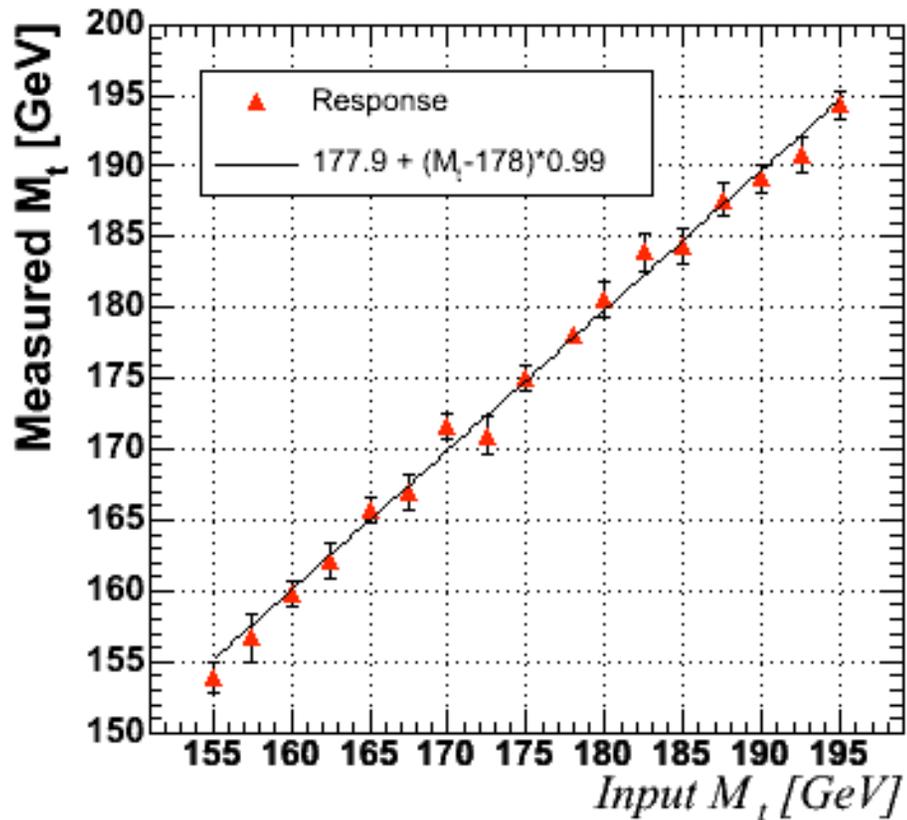


Pull width

Flat.

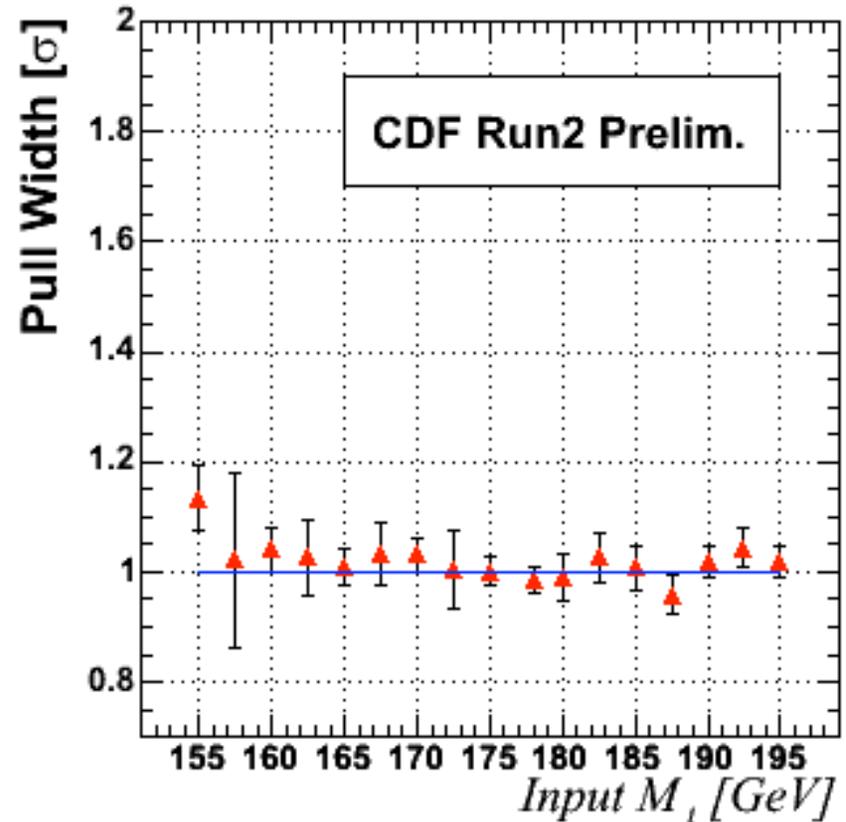
Width increased due to
unmodelled bg ($Z \rightarrow \tau\tau$, WZ , ZZ)

Response & Pulls



Response

After slope correction.

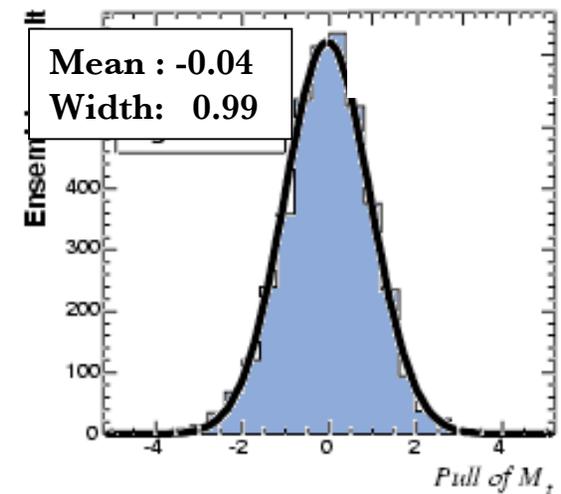
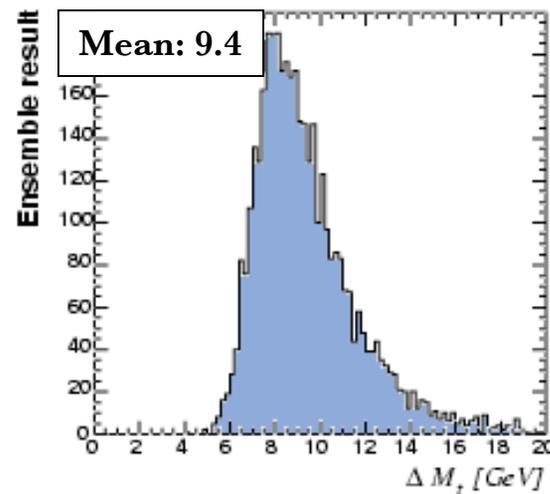
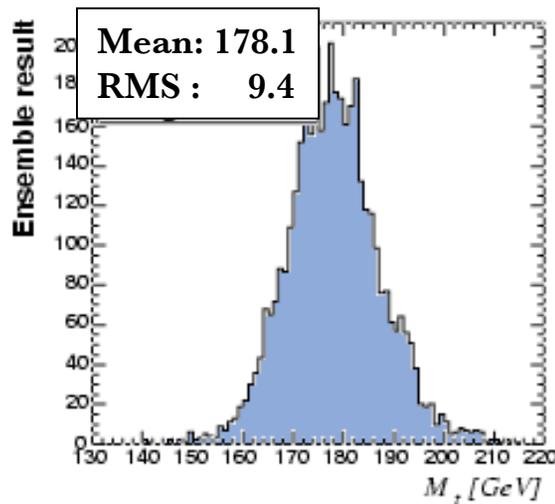


Pull width

After error correction.

Response at $M_t=178$ GeV

Residual, error, and pull distribution



<u>Method</u>	<u>Mean Error ($M_t=178$ GeV)</u>
Matrix Element	9.4 GeV
Neutrino Weighting	12.8 GeV
Kinematics	14.6 GeV
Neutrino Phi	14.9 GeV

Systematic Errors

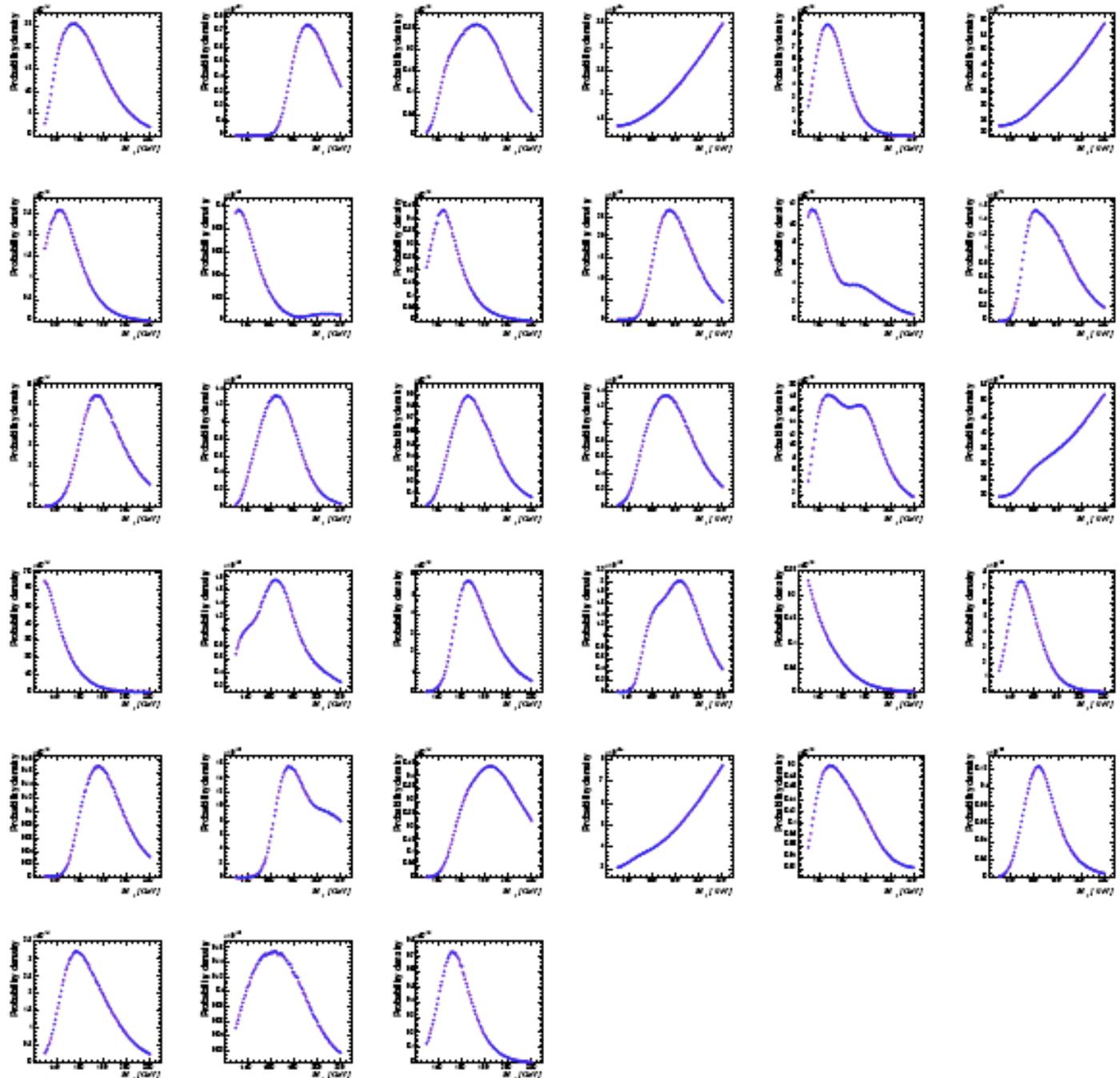
Source	Size
Jet Energy Scale	2.6 GeV
Generator	1.0 GeV
Method	0.6 GeV
Sample composition uncertainty	0.7 GeV
Background statistics	1.5 GeV
Background modelling	0.8 GeV
FSR modelling	0.5 GeV
ISR modelling	0.5 GeV
PDFs	1.1 GeV
Total	3.6 GeV

Data

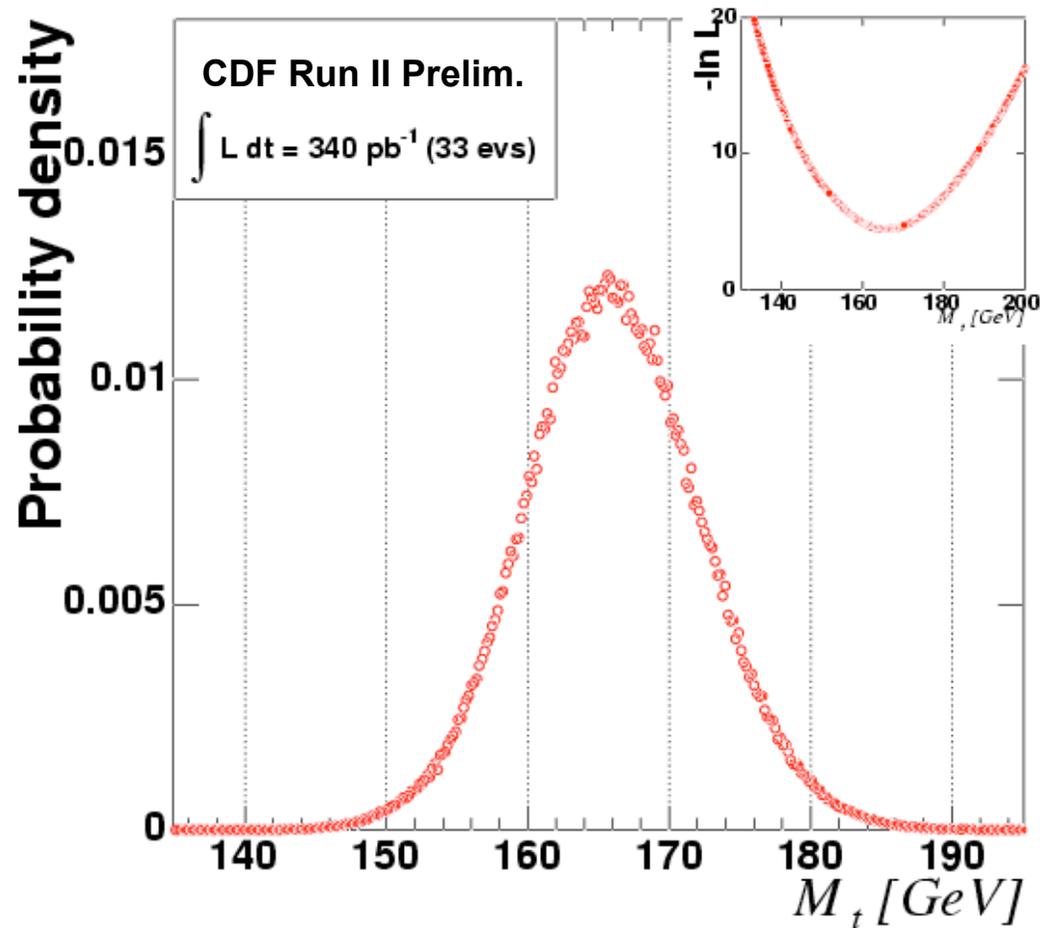
33 candidates
signal and bg
probabilities

Slope and error
corrections
not applied

Range is
130-220 GeV/c²



Measurement!



$$M_t = 165.3 \pm 6.3_{stat} \pm 3.6_{syst} \text{ GeV}/c^2$$

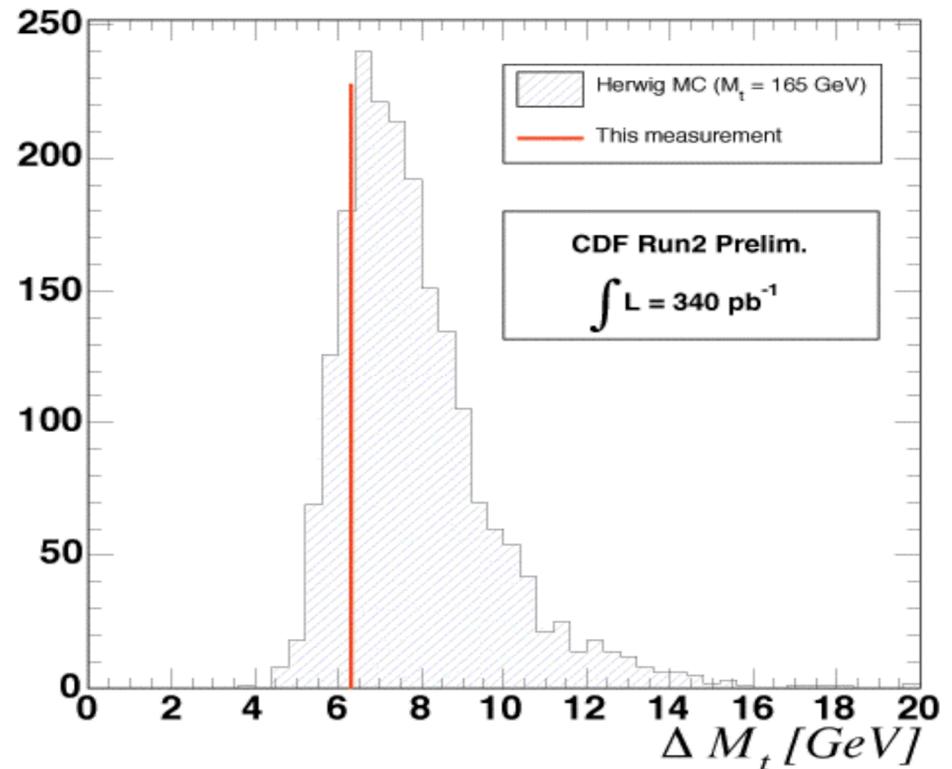
Most precise single dilepton measurement to date.

Expected Sensitivity

Sensitivity

If $M_t = 165 \text{ GeV}/c^2$,
mean stat error is **7.8 GeV/c²**

If $M_t = 178 \text{ GeV}/c^2$
mean stat error is **9.4 GeV/c²**



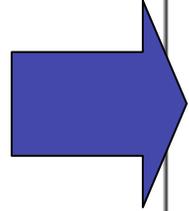
17% of pseudo-experiments
give smaller errors

Impact

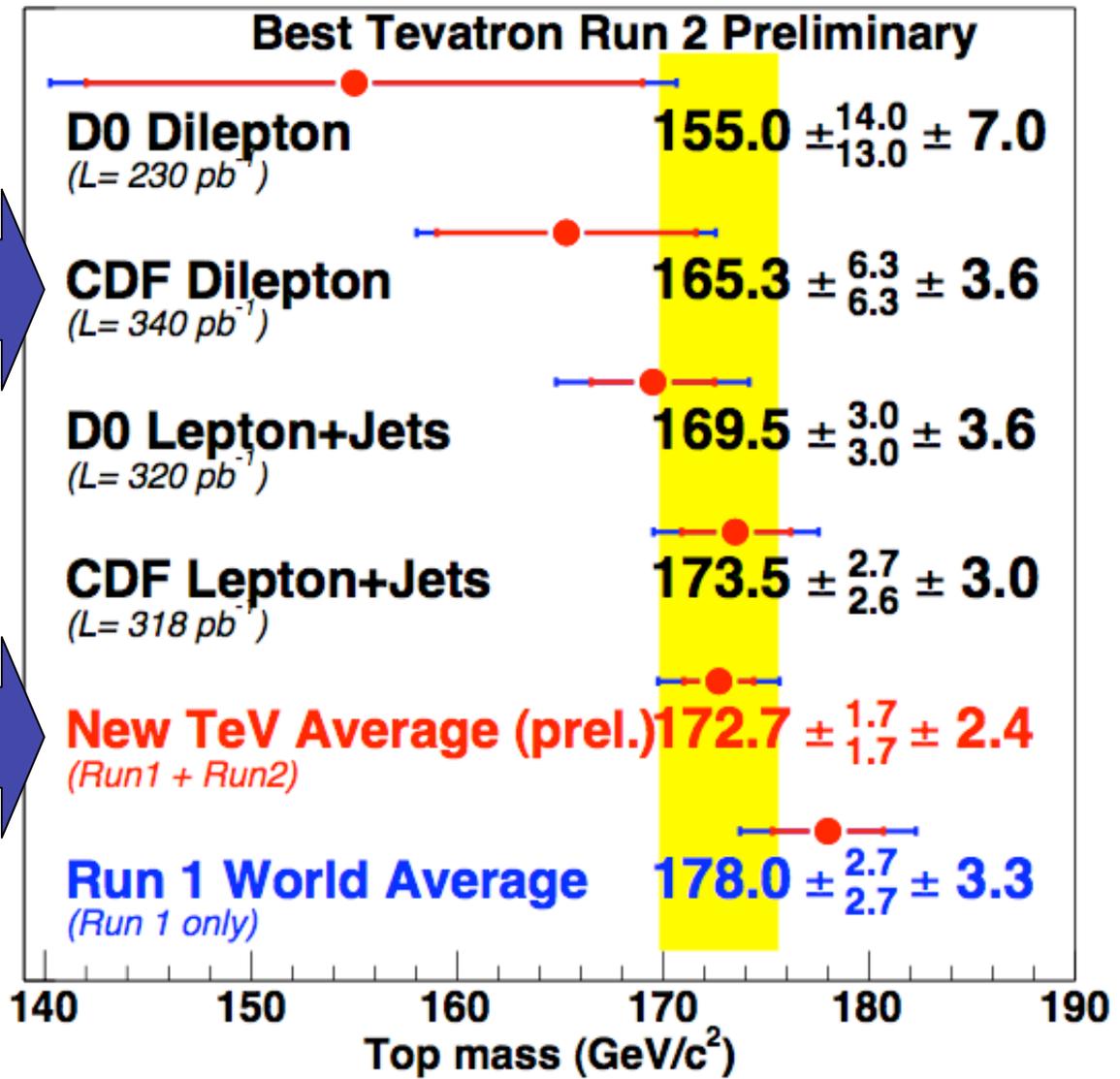
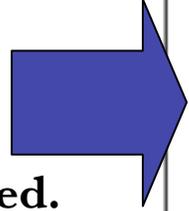
In World Average

Weight = 8%

Pull = -1.1



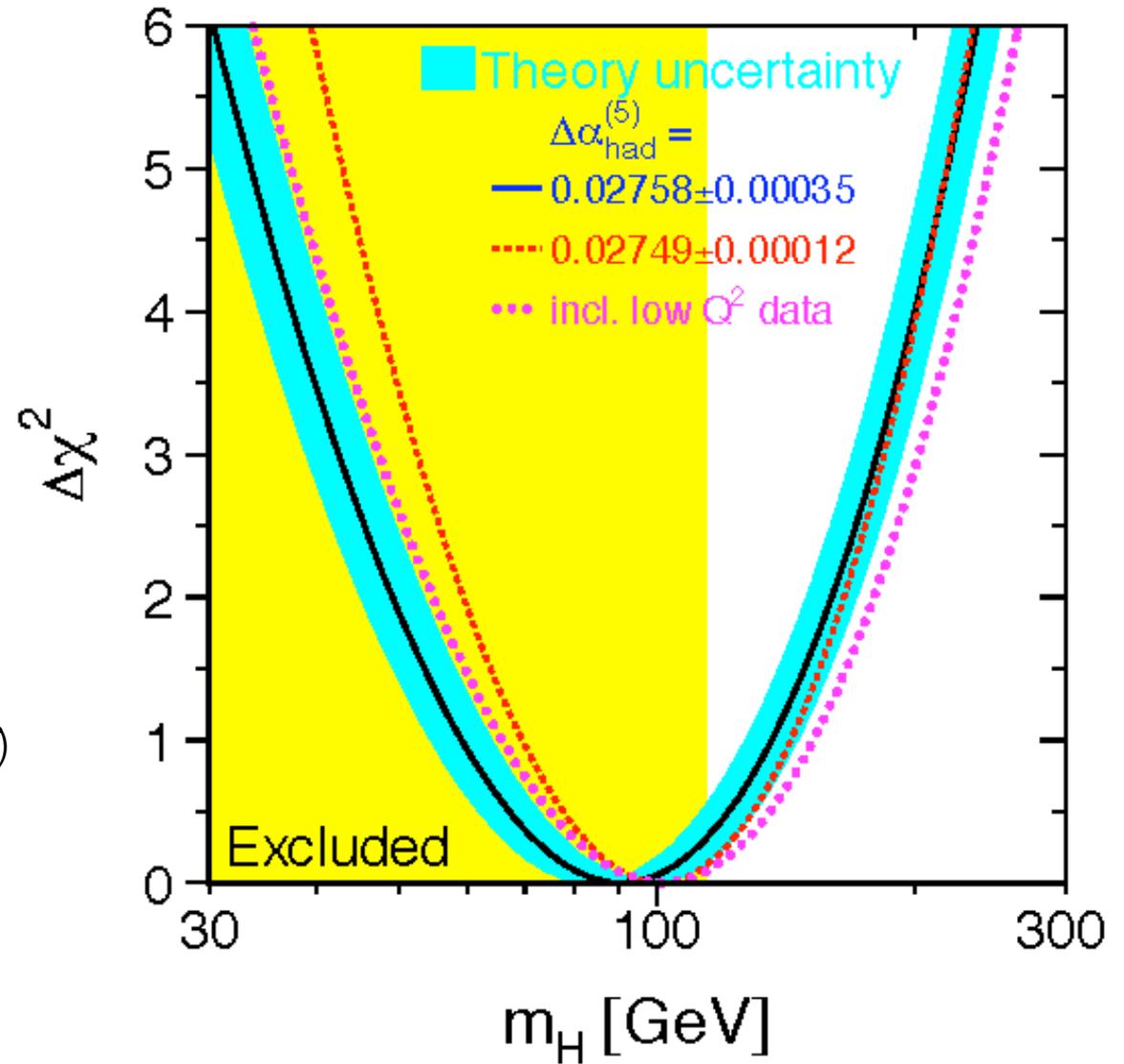
Total result has error of 2.9 GeV, systematics limited.



Higgs Impact

$$M_{\text{higgs}} = 91^{+45}_{-32} \text{ GeV}/c^2$$

$$M_{\text{higgs}} < 186 \text{ GeV}/c^2 \text{ (95\% conf)}$$



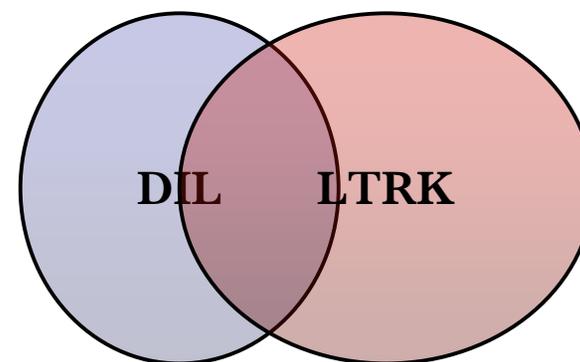
Combining dilepton results

Statistical Correlation

Data samples overlapping

Measure correlation in pseudo-experiments
which model common data

Two overlapping datasets



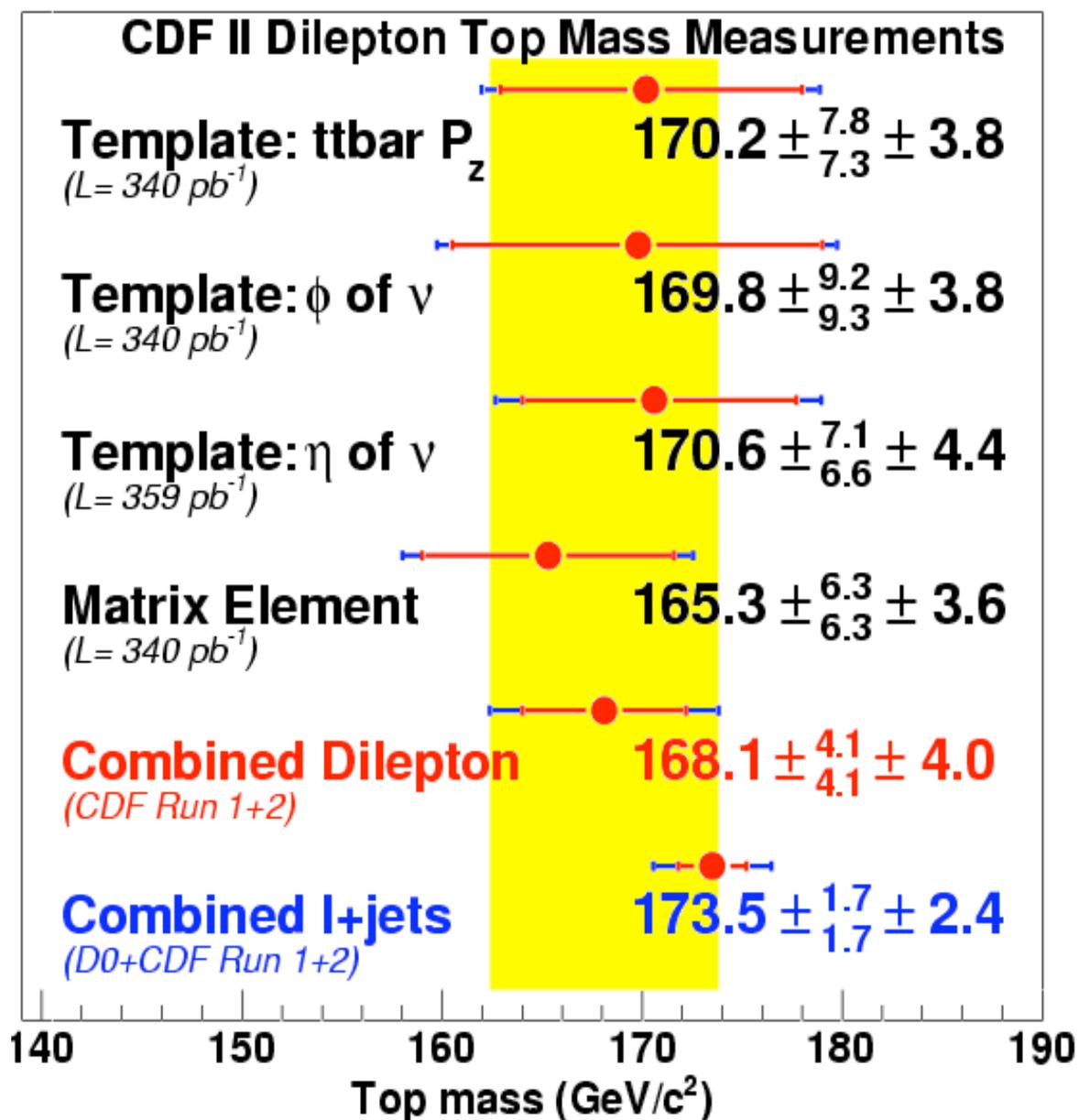
Statistical Correlation Matrix [CDF Run2 Preliminary]

	<u>ME</u>	<u>KIN</u>	<u>PHI</u>	<u>NWA</u>	<u>Run1</u>
<i>DIL</i> sample	Matrix Element	1.00			
	Kinematic	0.33	1.00		$\Delta\rho \sim 0.06$
<i>LTRK</i> sample	Neutrino Phi	0.15	0.18	1.00	
	Neutrino Weight	0.06	0.17	0.26	1.00
	Run1	0.00	0.00	0.00	0.00

Combined Dilepton Result

Preliminary

	<i>Weight Pull</i>	
Matrix El.	36%	-0.60
ν Weight	27%	0.42
Kinematic	13%	0.33
ν Phi	10%	0.20
Run1	14%	-0.07



Summary & Outlook

Precision Measurement

We have measured the top mass in dileptons with a stat. error of **4.1 GeV**

Results consistent with 1+jets

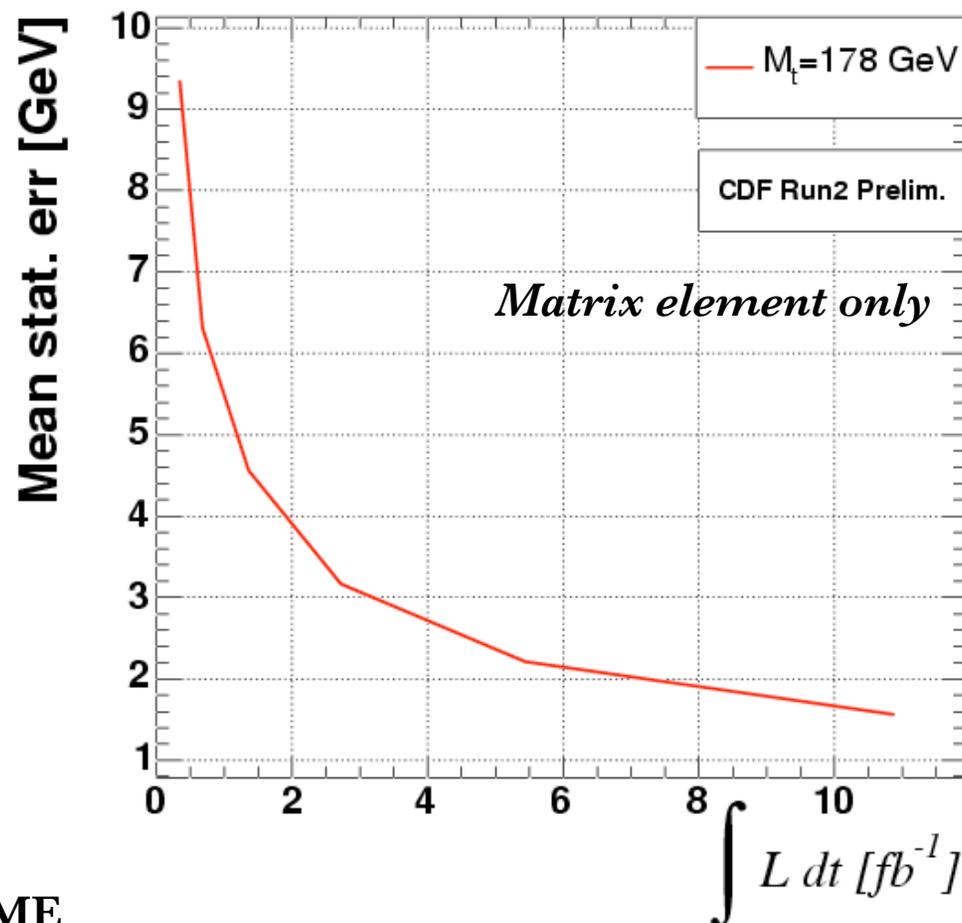
Room for Improvement

Identified assumptions which are critical to sensitivity

High Precision Future

With no improvements, in **4fb⁻¹**,

- expect **2.5 GeV** statistical error for ME
- Dileptons become systematics limited.
- 1+jets and dileptons approach equal weight
- Comparison of channels to test top hypothesis, look for new physics
- Systematic challenge: reduce jet energy scale (no W→jj decay for calibration)

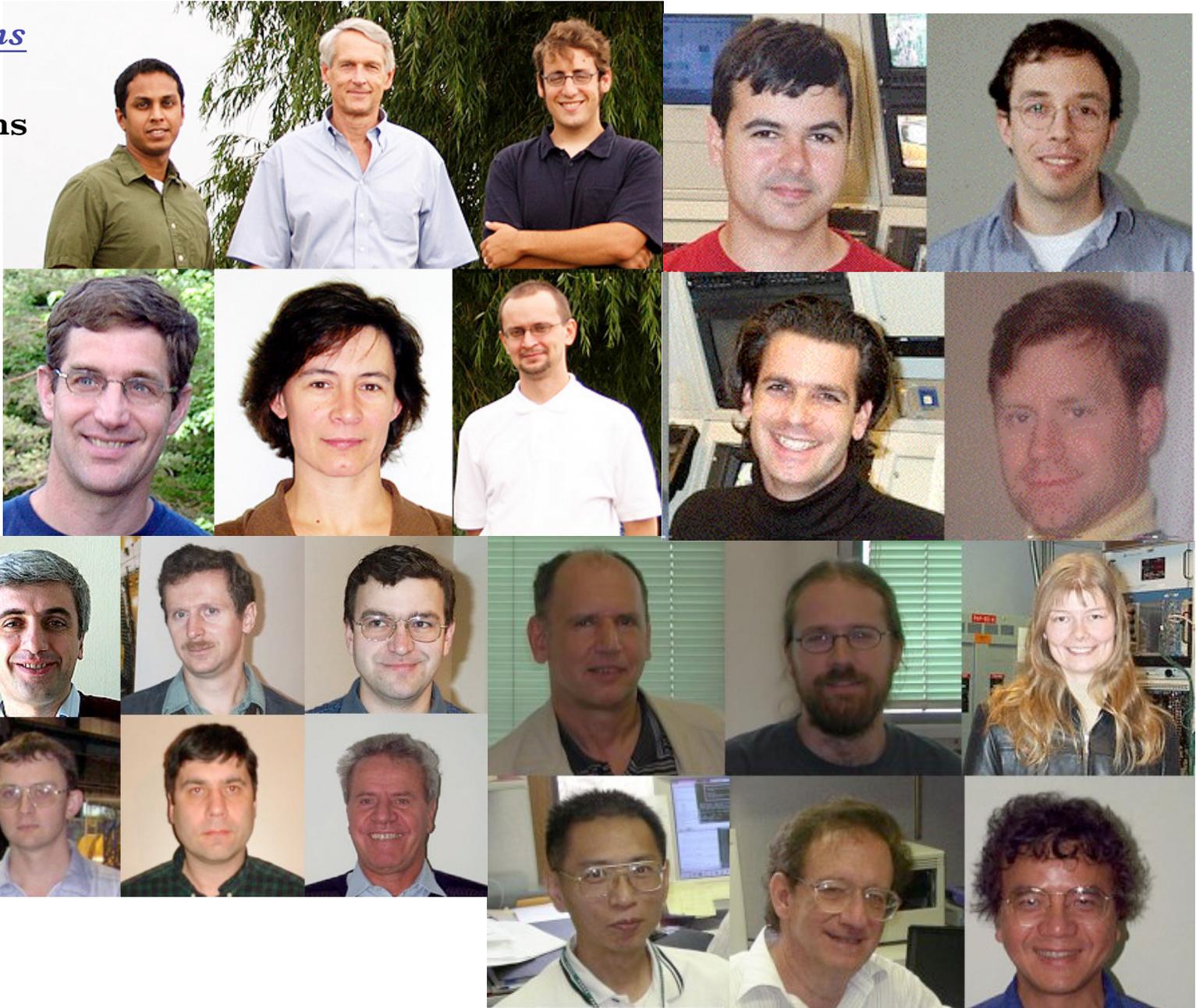


CDF Dileptons

24 people

9 institutions

7 countries

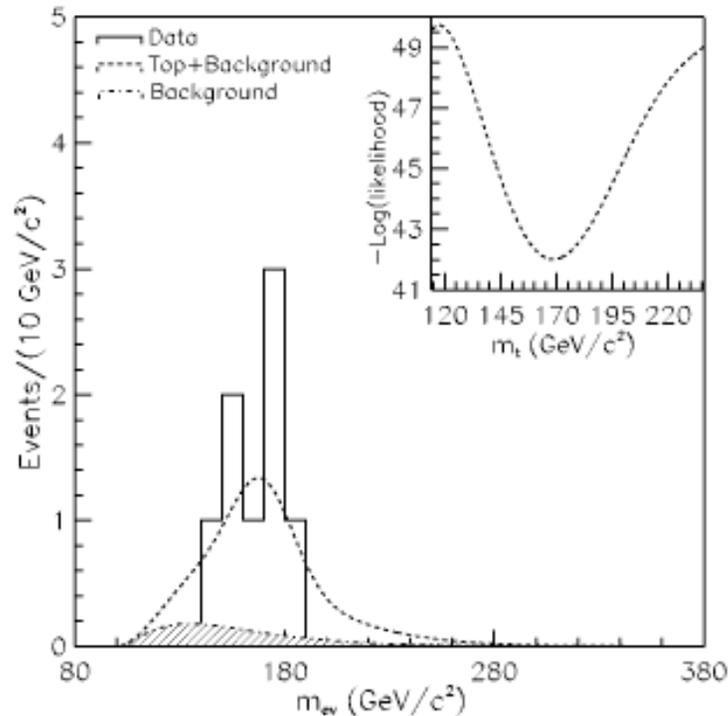


Backup material

Dilepton History: Run I

CDF: 8 events

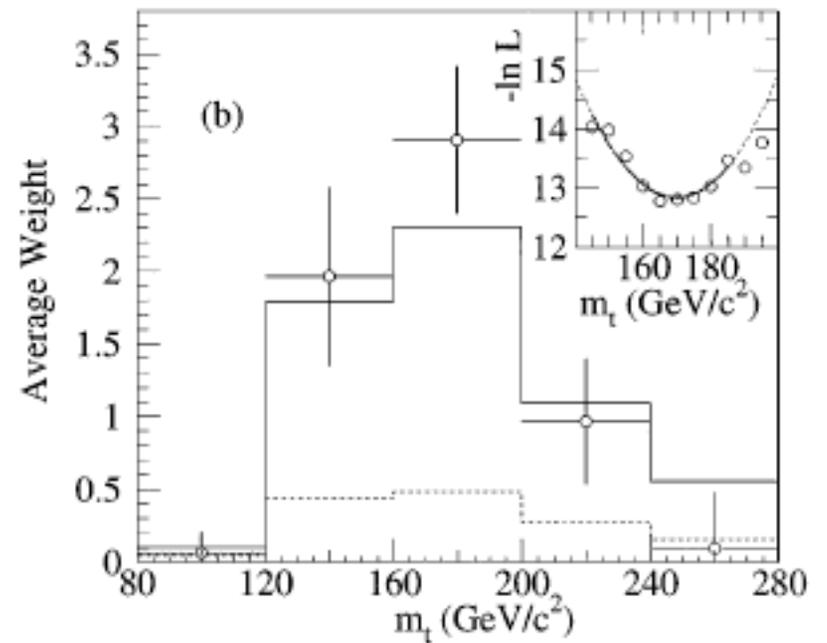
$M_t = 167.4 \pm 10.3 \pm 4.8 \text{ GeV}$



Phys. Rev. Lett. **82**, 271 (1999)

D0: 6 events

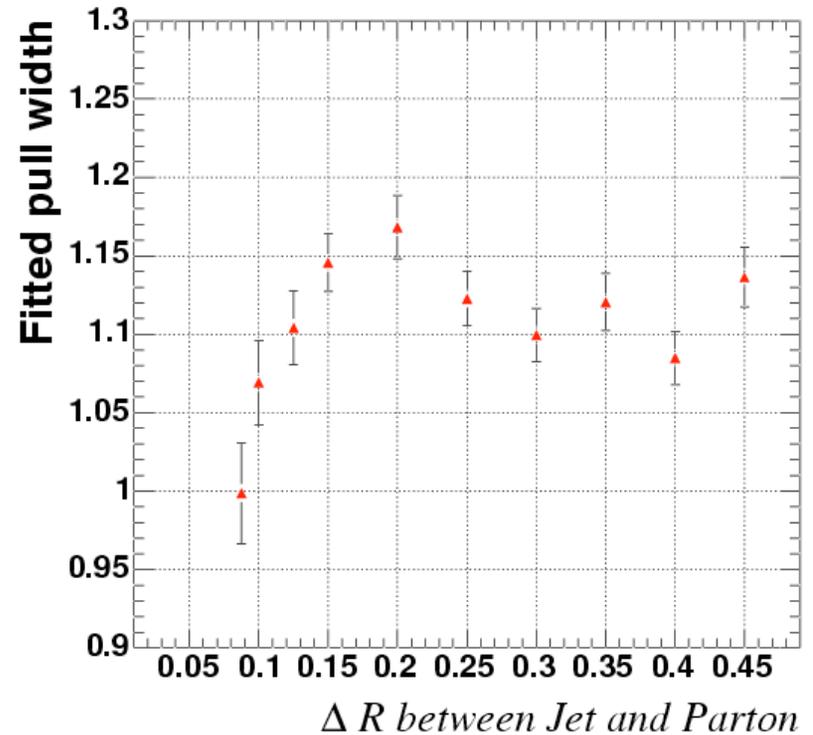
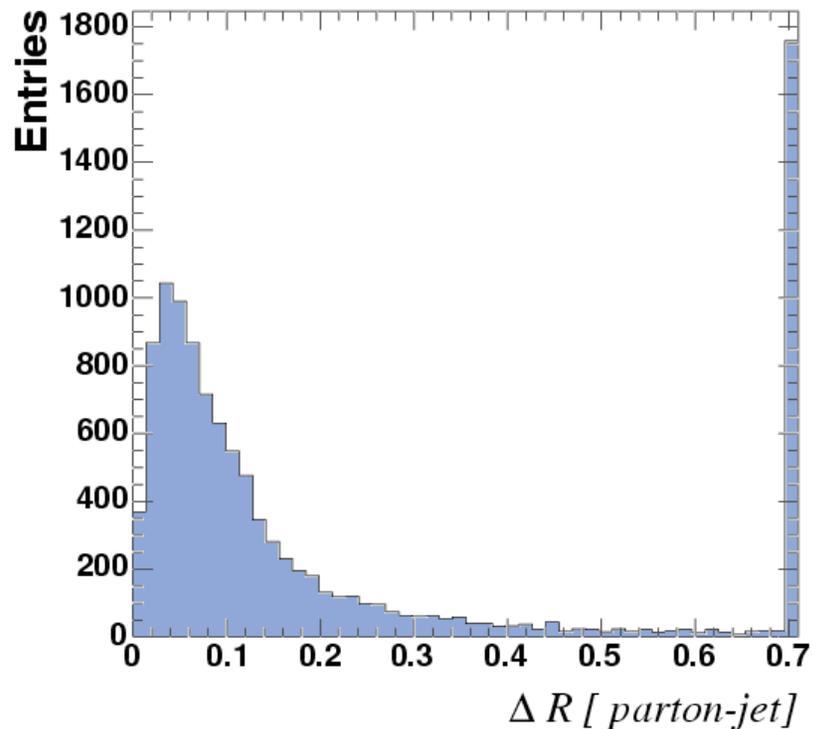
$M_t = 168.4 \pm 12.4 \pm 3.6 \text{ GeV}$



Phys Rev Lett **82**, 2063 (1998)

Pull Width: jet angles

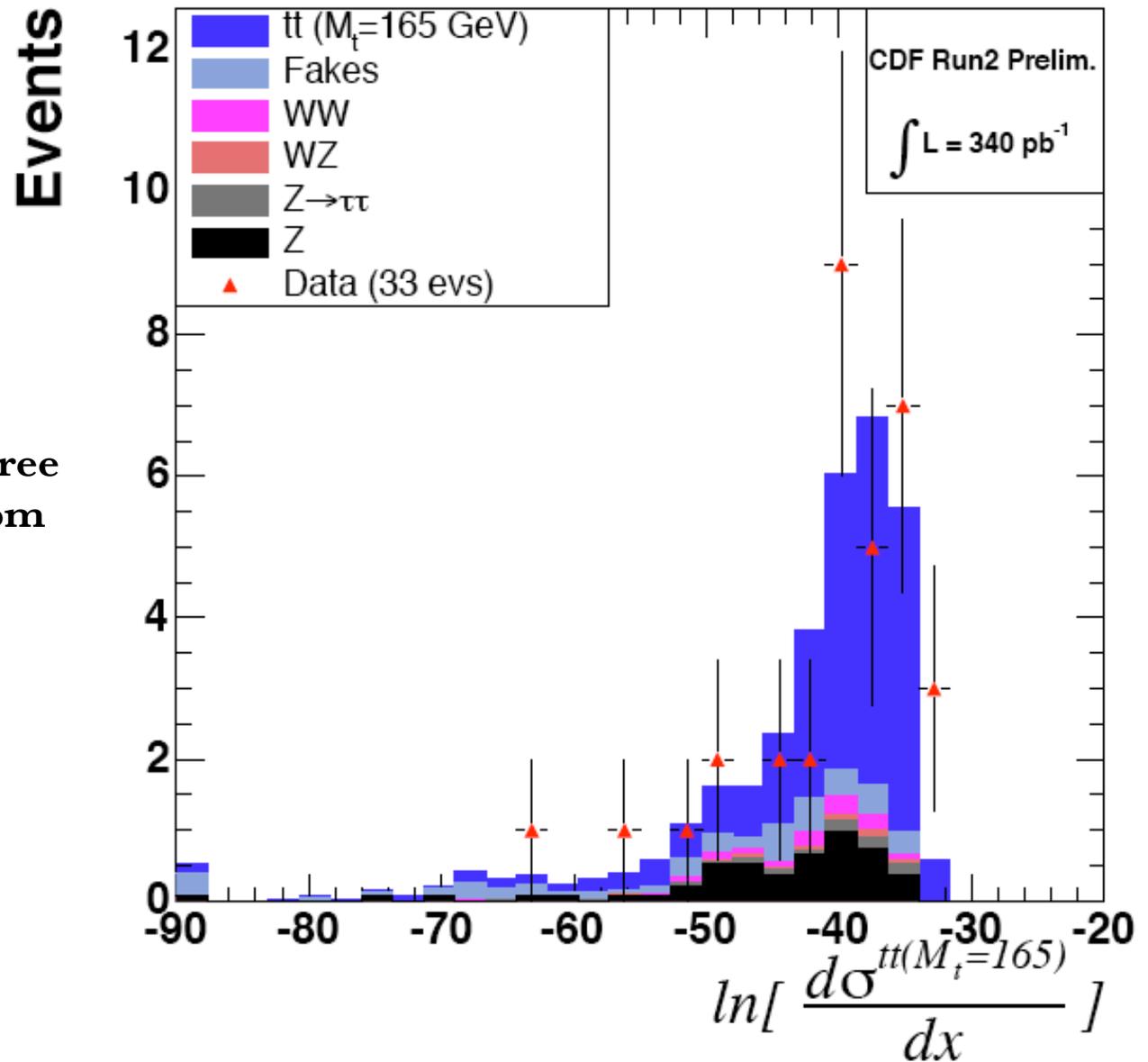
Requiring two matched b-jets
Requiring well-measured leptons



Pull width decreases to ~ 1 as angle improves

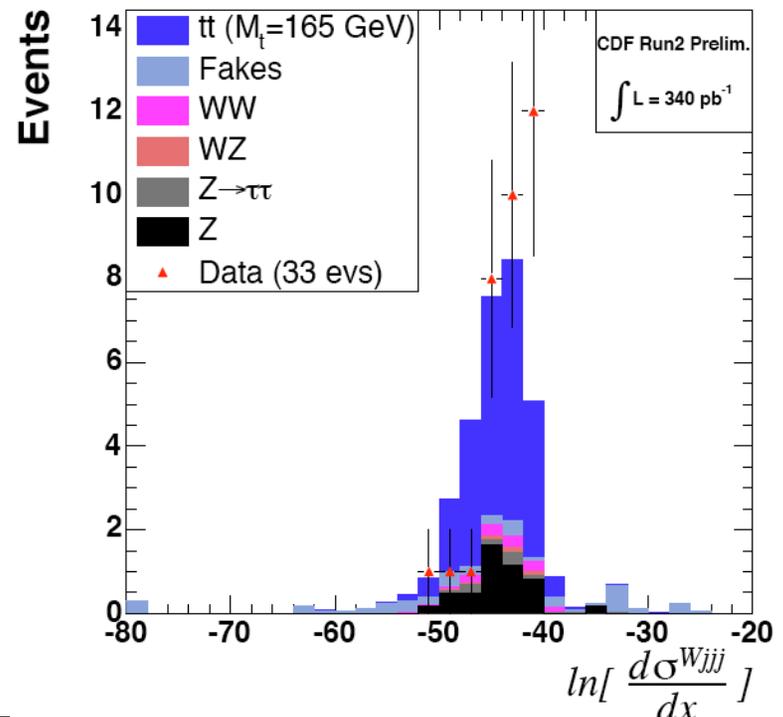
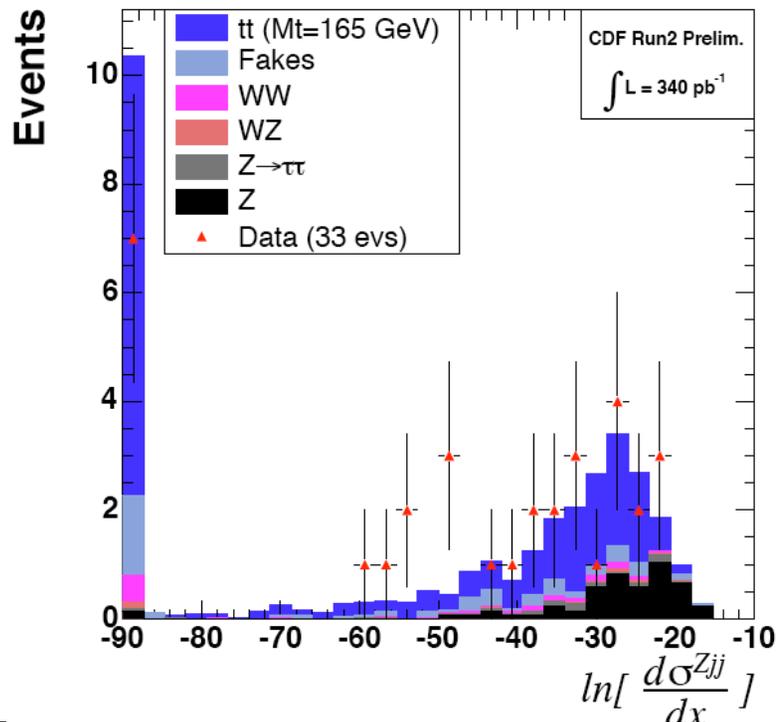
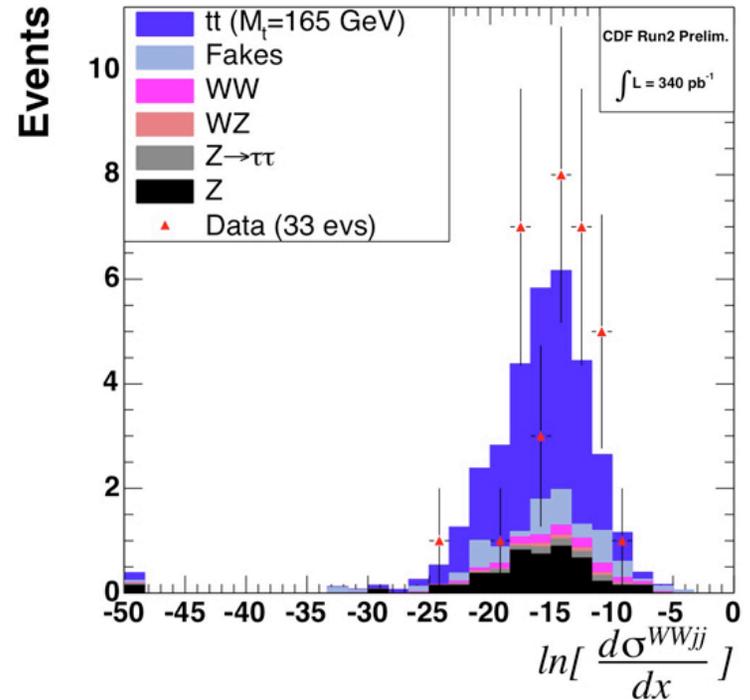
Signal probability in data

Distribution of un-normalized signal probability in data agree with that expected from MC



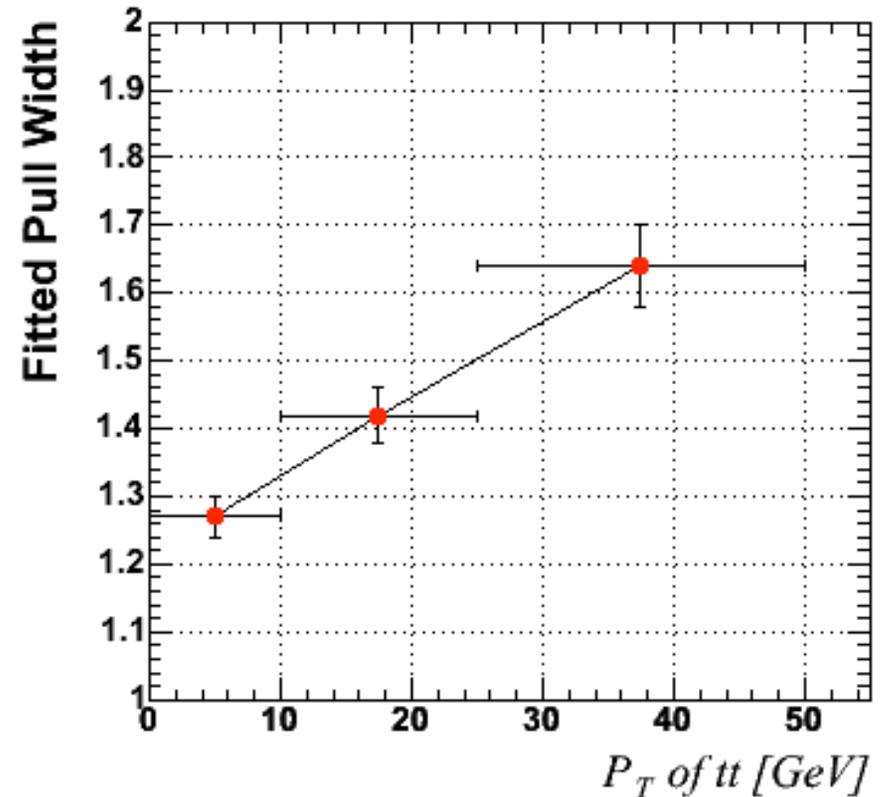
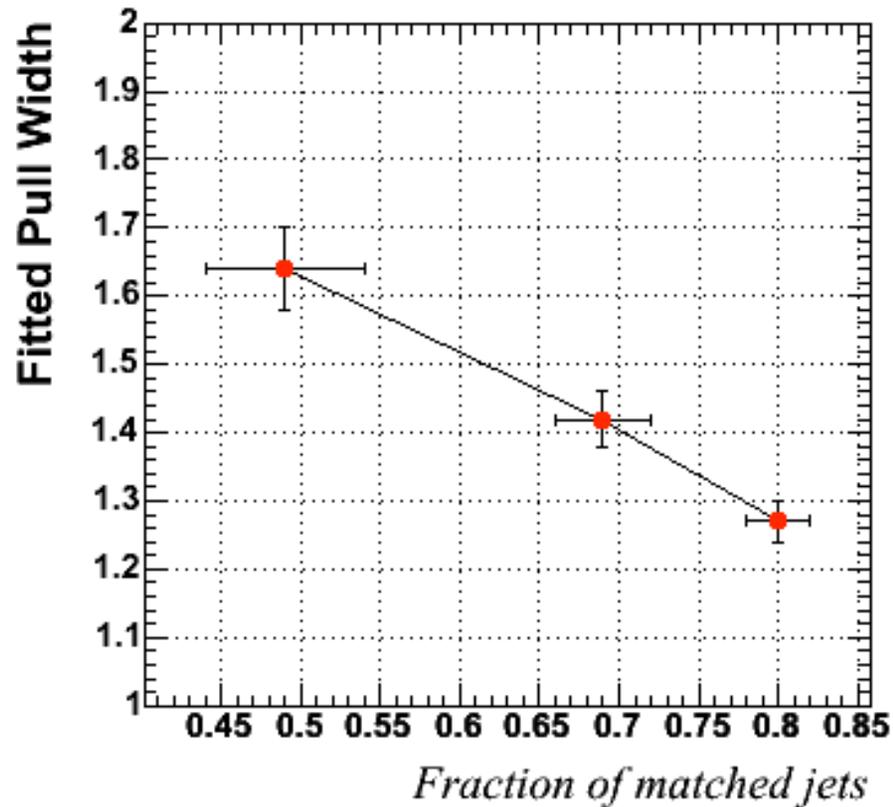
Background probabilities in data

Distribution of un-normalized background probabilities in data agree with that expected from MC



Pull width: mismatched jets

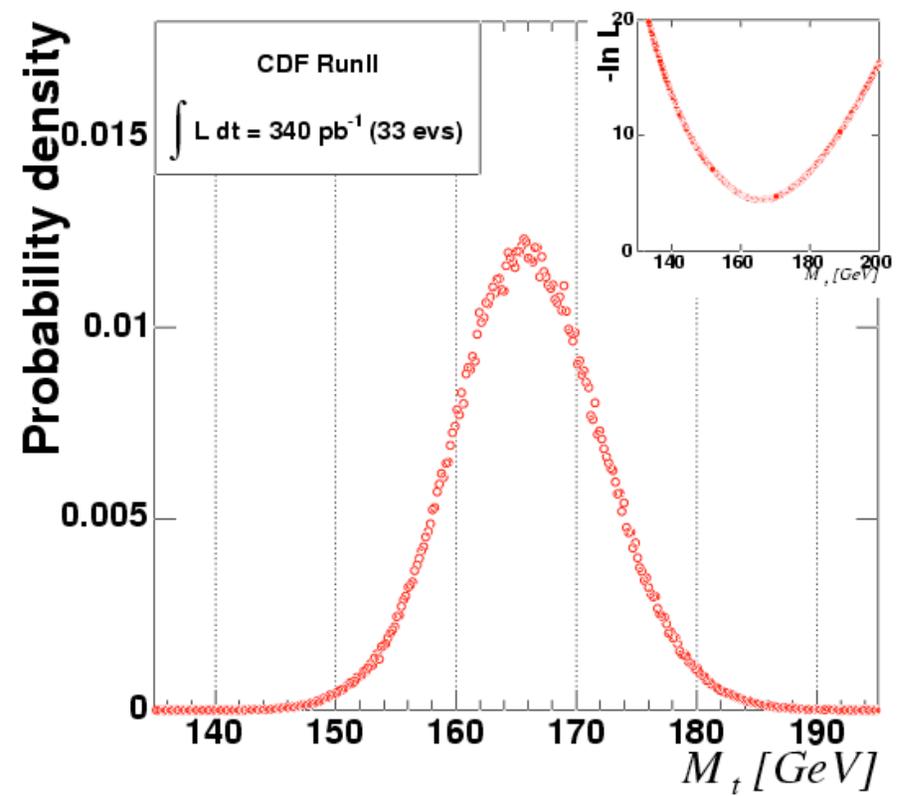
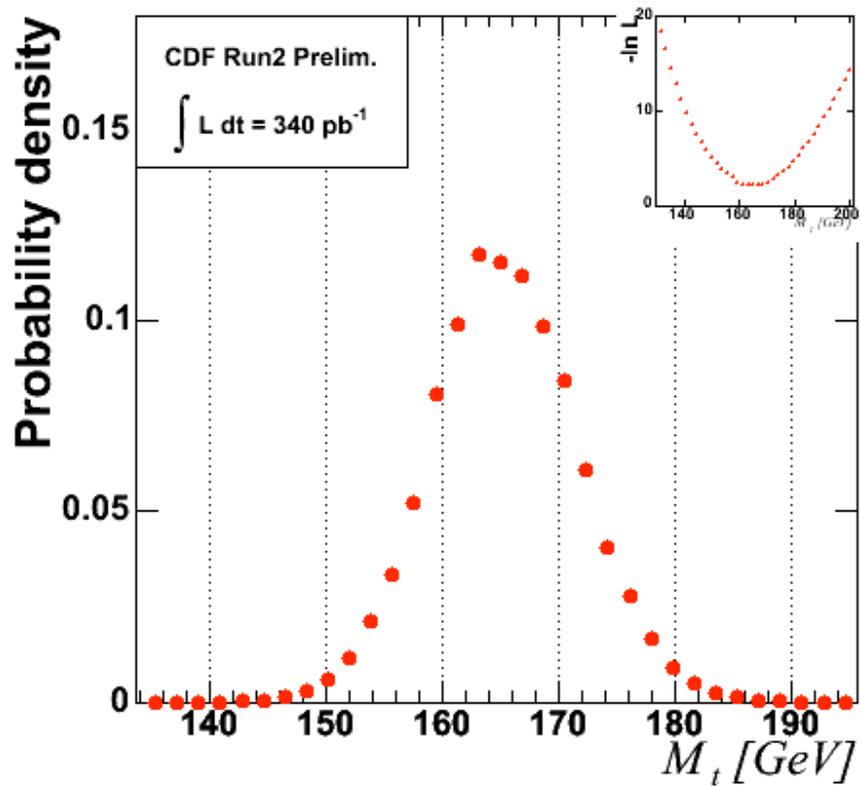
Pull width is affected by wrong jets



Most wrong jets come from initial state radiation, which can be probed by examining P_T of $t\bar{t}$

Mass steps refined

We scanned the space in M_t with finer steps to probe the shape:



Effect of SUSY events on dilepton mass measurement

Chargino/Neutralino

Topology is $ll+2j$ or $llqq$

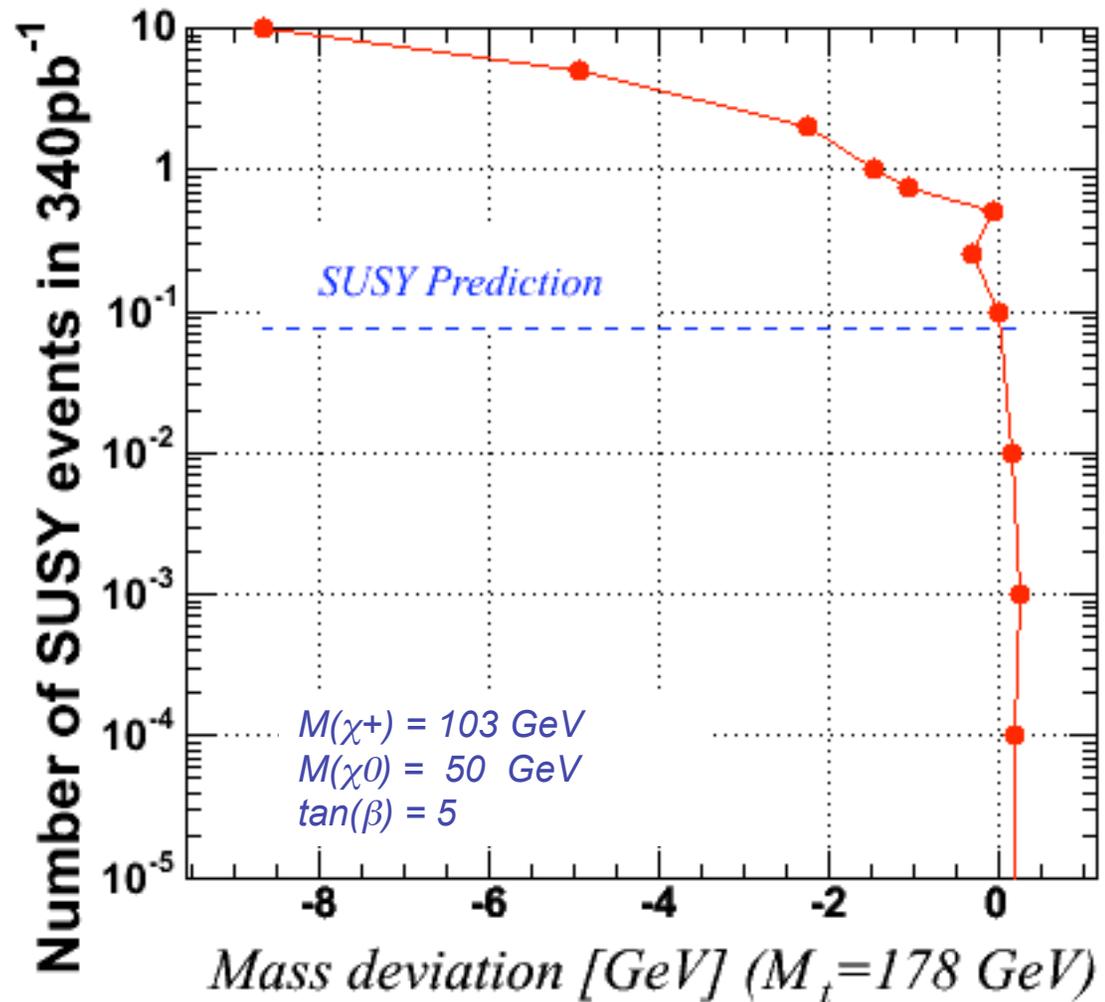
$$M(\chi^+) = 103 \text{ GeV}$$

$$M(\chi^0) = 50 \text{ GeV}$$

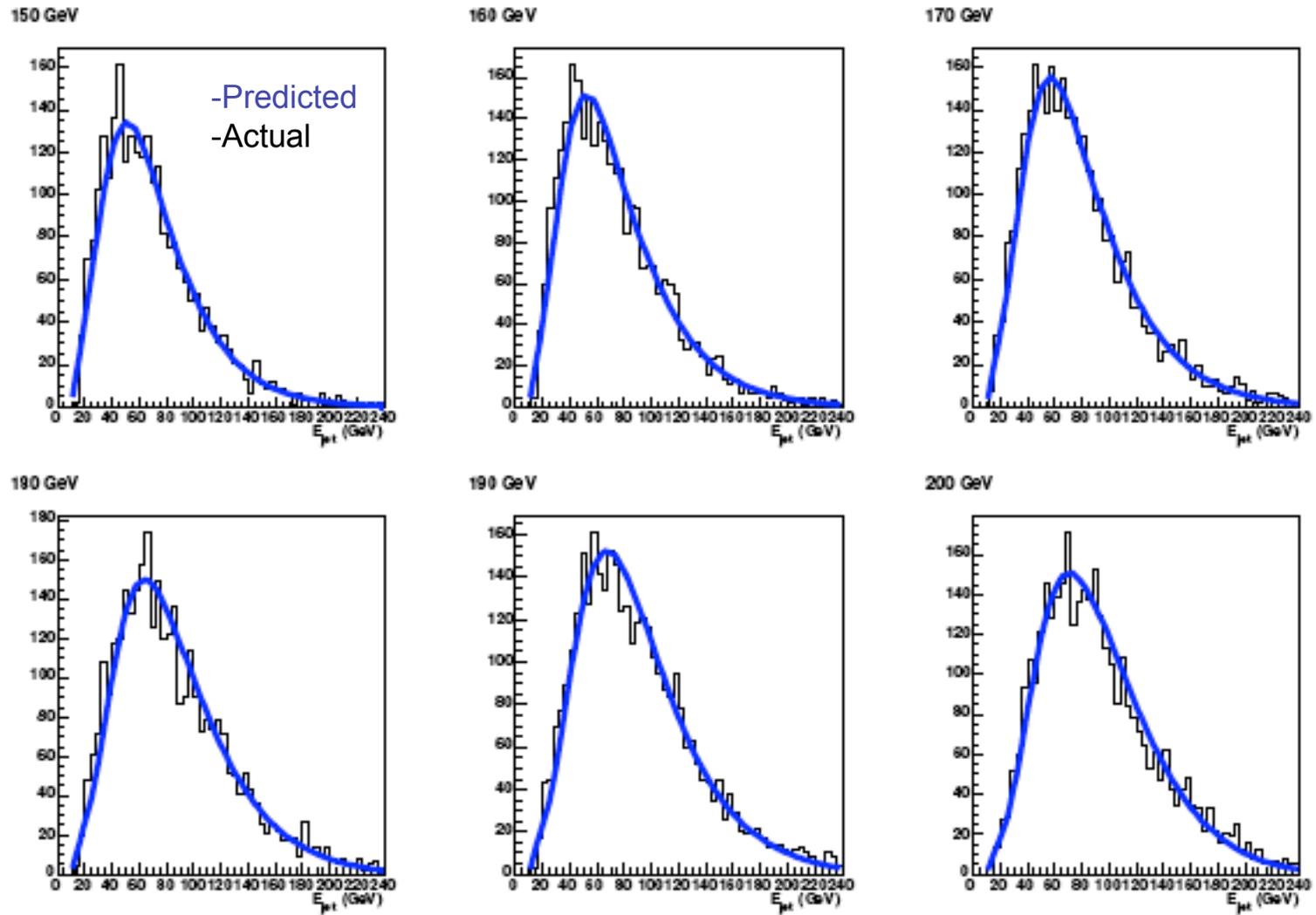
$$\tan(\beta) = 5$$

$$\text{Sigma} \cdot \text{br} = 150 \text{ fb}$$

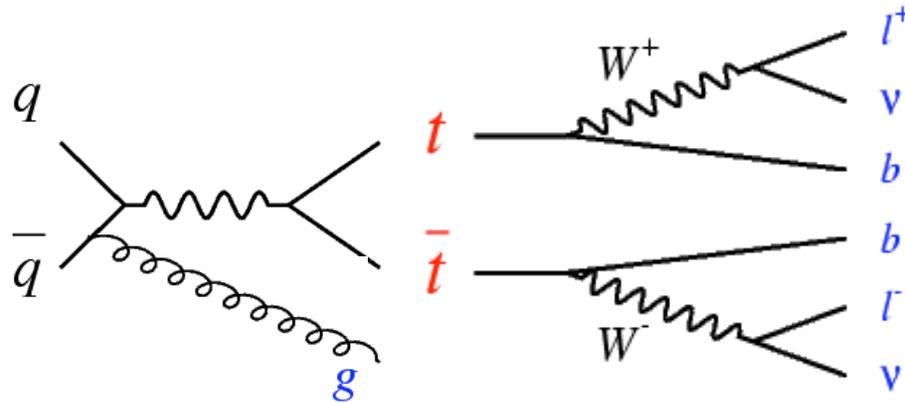
$$\text{Acc} = 0.15\%$$



Transfer functions predict jet energy spectrum at varying top masses.



Future work



Statistical error

- Improve handling of extra jets
(approximates NLO effect)

Systematic error

- Apply jet energy calibration from $Z \rightarrow b\bar{b}$
- Improve sophistication of background modelling