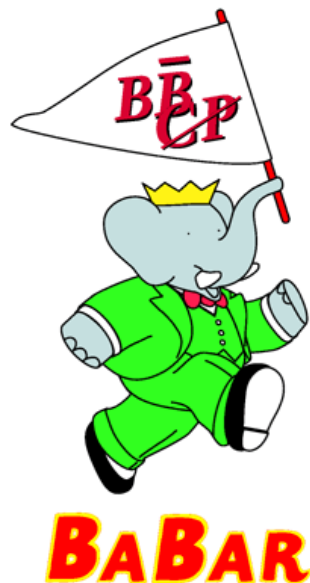


Recent progress in CP violation: New Physics under siege

*High Energy Seminar
University of Pennsylvania
Tuesday, January 29, 2008*

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Princeton University**

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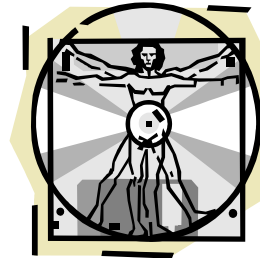
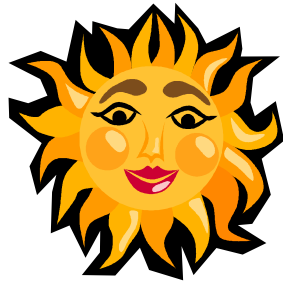
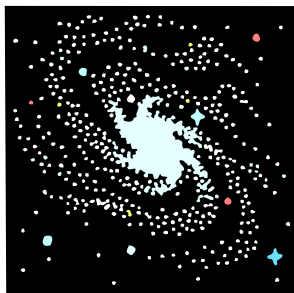




Matter-Antimatter Asymmetry the Universe

We are survivors of the post-Big Bang
mutual annihilation of matter and antimatter

Baryonic asymmetry of the Universe: $\frac{N_B - N_{\bar{B}}}{N_y} \simeq \frac{N_B}{N_y} \simeq 5 \times 10^{-10}$





Sakharov's three conditions

The three conditions necessary to produce the baryonic asymmetry of the Universe:



(photo circa 1943)

А. Д. Сахаров, *Письма в ЖЭТФ*, **5**, № 1, 32-35, 1 января 1967
A. D. Sakharov, *Soviet Journal of Experimental and Theoretical Physics, Letters to the Editor*, **5**, No. 1, 24-27, 1st January 1967

1. Baryon number violation

2. C and CP violation

This is by far the strongest argument for existence of sources of *CP* violation in nature

3. Departure from thermal equilibrium

Through the *CPT* Theorem, *CP* violation implies the existence of *T* violation



The Quark Mixing Matrix

The only Standard-Model source of CP violation in the quark sector

The Cabibbo-Kobayashi-Maskawa matrix relates the electroweak (q') and the mass (q) quark eigenstates:

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$V_{\text{CKM}} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4) + iO(\lambda^6)$$

$$V^\dagger V = 1 \Rightarrow V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$

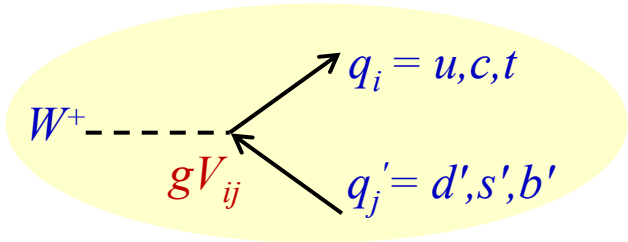
← the "unitarity triangle"

$$\alpha \equiv \arg \left[-\frac{V_{td} V_{tb}^*}{V_{ud} V_{ub}^*} \right]$$

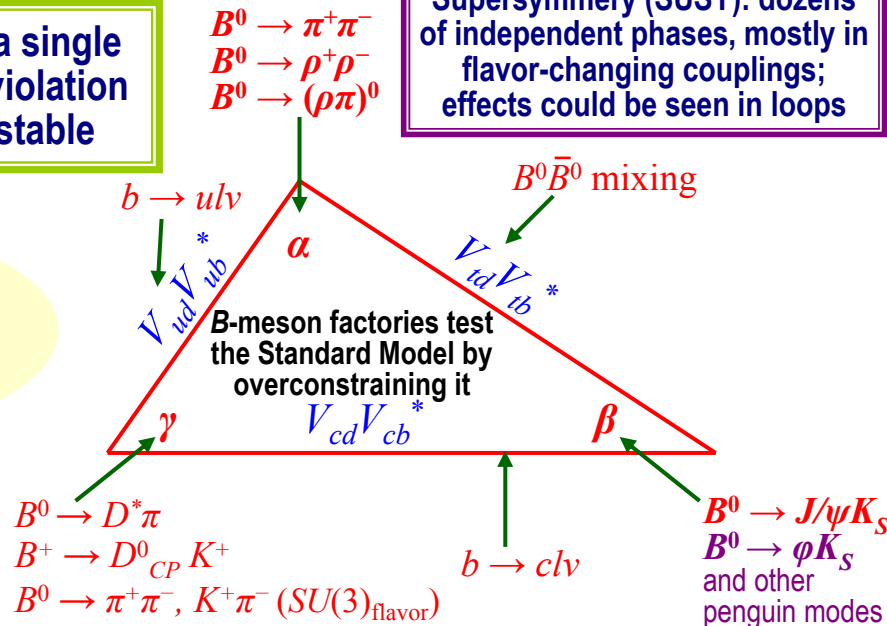
$$\beta \equiv \arg \left[-\frac{V_{cd} V_{cb}^*}{V_{td} V_{tb}^*} \right]$$

$$\gamma \equiv \arg \left[-\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*} \right]$$

In Kobayashi-Maskawa theory, a single phase is responsible for all CP violation in meson decays, making it testable



Supersymmetry (SUSY): dozens of independent phases, mostly in flavor-changing couplings; effects could be seen in loops



CKM matrix overview: Review of Particle Physics, W.-M. Yao et al. [Particle Data Group], J. Phys. G **33**, 1 (2006)

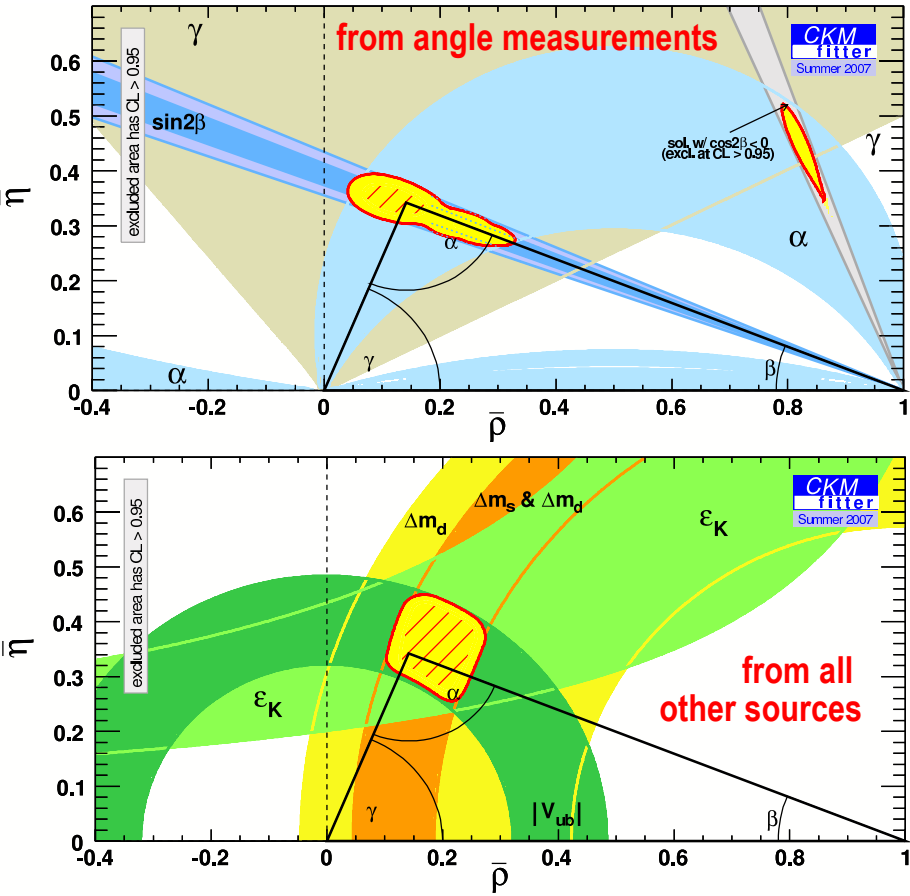


Current knowledge of the Unitarity Triangle

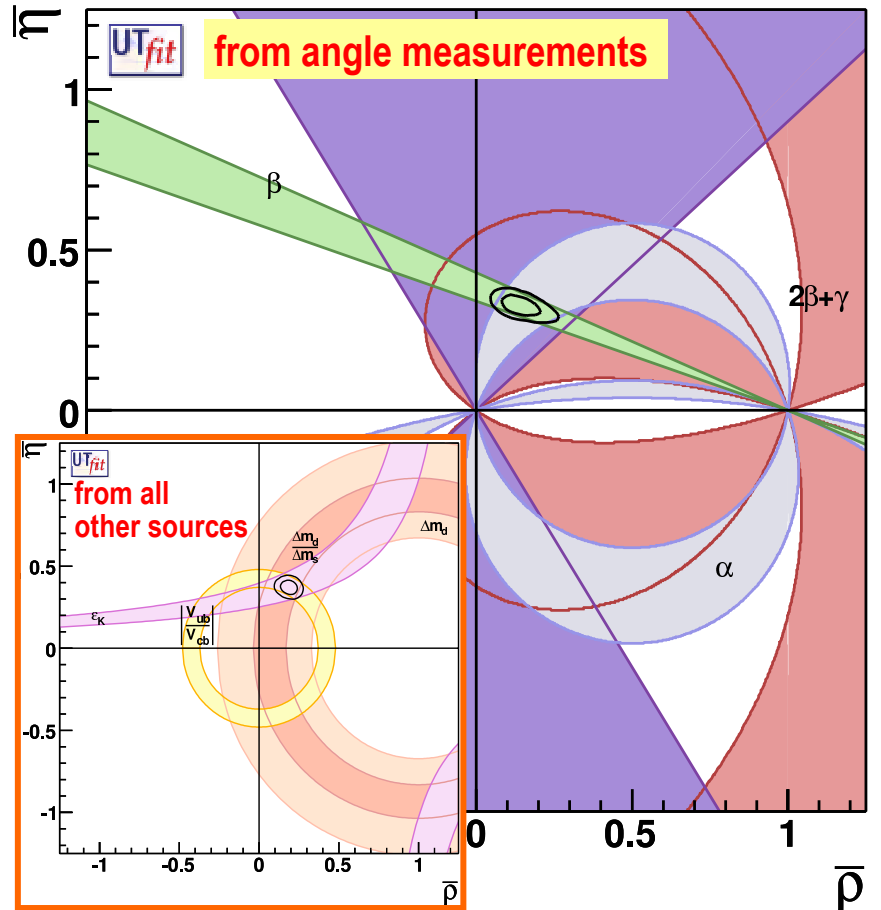
Important: all measurements are still statistics-limited

A frequentist interpretation

A Bayesian interpretation, with model-dependent choices of priors



CKMfitter Group (J. Charles et al.), Eur. Phys. J. C41, 1-131 (2005), [hep-ph/0406184], updated results and plots available at <http://ckmfitter.in2p3.fr>



M. Ciuchini, G. D'Agostini, E. Franco, V. Lubicz, G. Martinelli, F. Parodi, P. Roudeau, A. Stocchi, JHEP 0107 (2001) 013 [hep-ph/0012308], updated results and plots available at <http://utfit.roma1.infn.it>



Time evolution of the B^0 meson

The time-dependent rate for \bar{B}^0 (f_+) or B^0 (f_-) decays to a final state f (neglecting the lifetime difference between the mass eigenstates B_H and B_L):

$$f_{\pm}(\Delta t) = \frac{e^{-|\Delta t|/\tau_B}}{4\tau_B} [1 \mp C_f \cos(\Delta m \Delta t) \mp S_f \sin(\Delta m \Delta t)]$$

where

S and C is what we measure

$$|B_{L/H}\rangle = p|B^0\rangle \pm q|\bar{B}^0\rangle$$

$$\lambda_f = \frac{q}{p} \frac{\bar{A}_f}{A_f}$$

from mixing, $\approx e^{-2i\beta}$

$$S_f = \frac{-2 \operatorname{Im} \lambda_f}{1 + |\lambda_f|^2}, \quad C_f = \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2}$$

a_f is the time-evolution asymmetry:

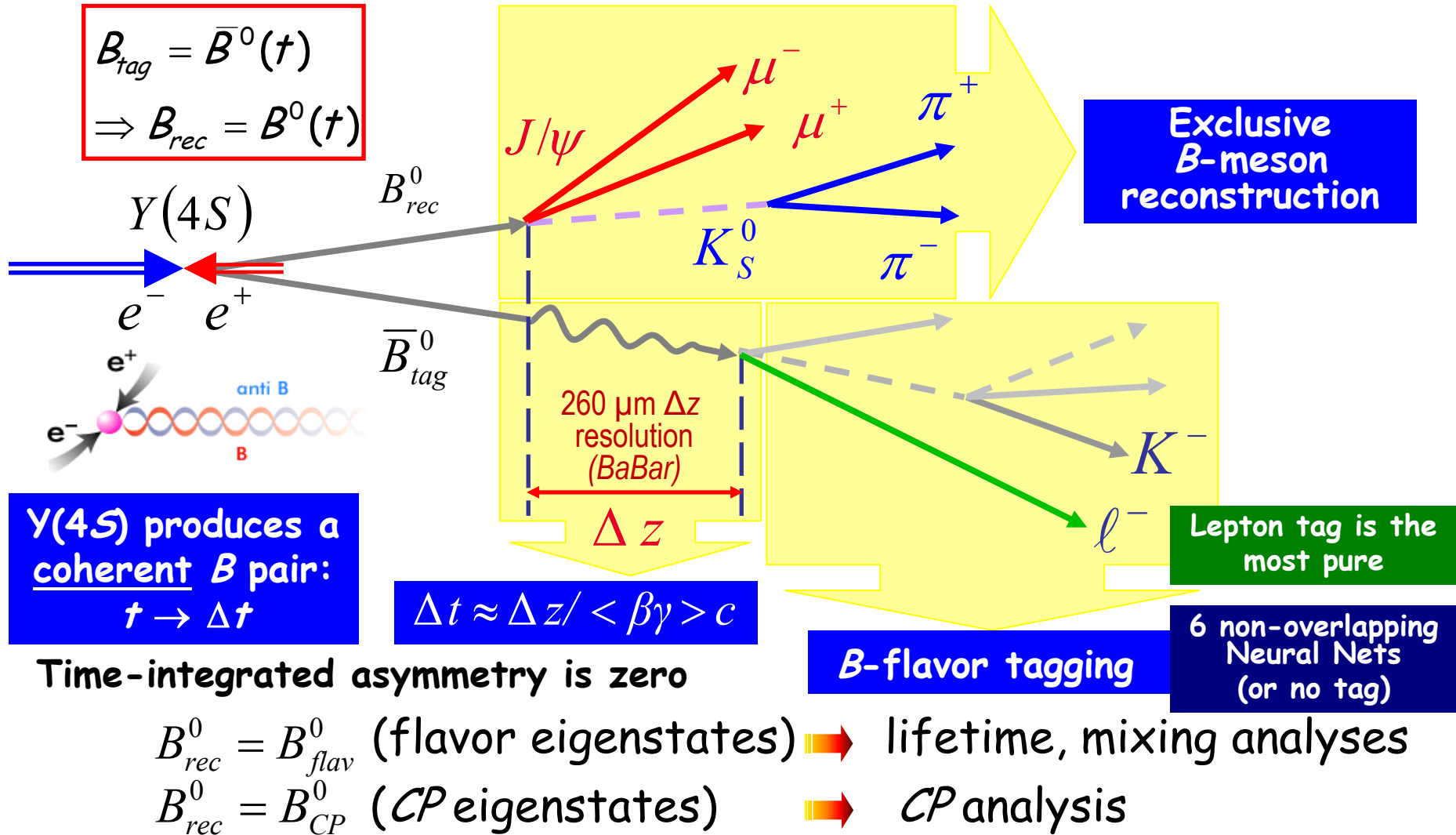
$$a_f(\Delta t) = \frac{f_+(\Delta t) - f_-(\Delta t)}{f_+(\Delta t) + f_-(\Delta t)}$$

If f is a CP eigenstate, f_{CP} , we have CP violation if $\lambda_f \neq \pm 1$:

- $|q/p| \neq 1$ (CP violation in mixing, very small) ← *B. Aubert et al. (BaBar Collaboration) Phys. Rev. Lett. 96, 251802 (2006)*
- $|\bar{A}_f/A_f| \neq 1$ (direct CP violation, small in $b \rightarrow c\bar{c}s$)
- $\operatorname{Im}(\lambda_f) \neq 0$ (interference between mixing and decay)



Time-dependent CP analysis at a B -meson factory

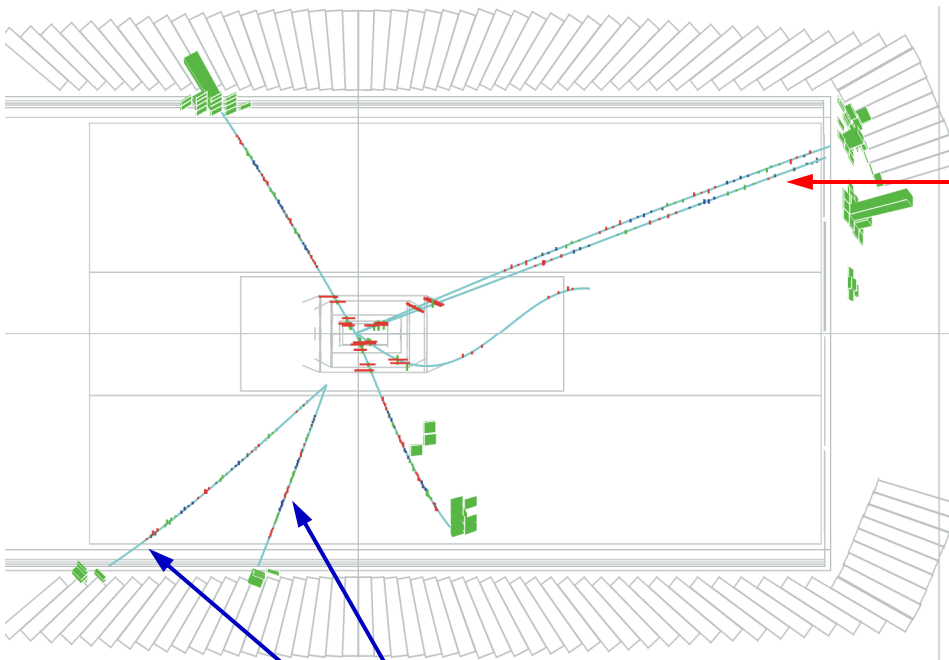




Example of a B -decay event in data:

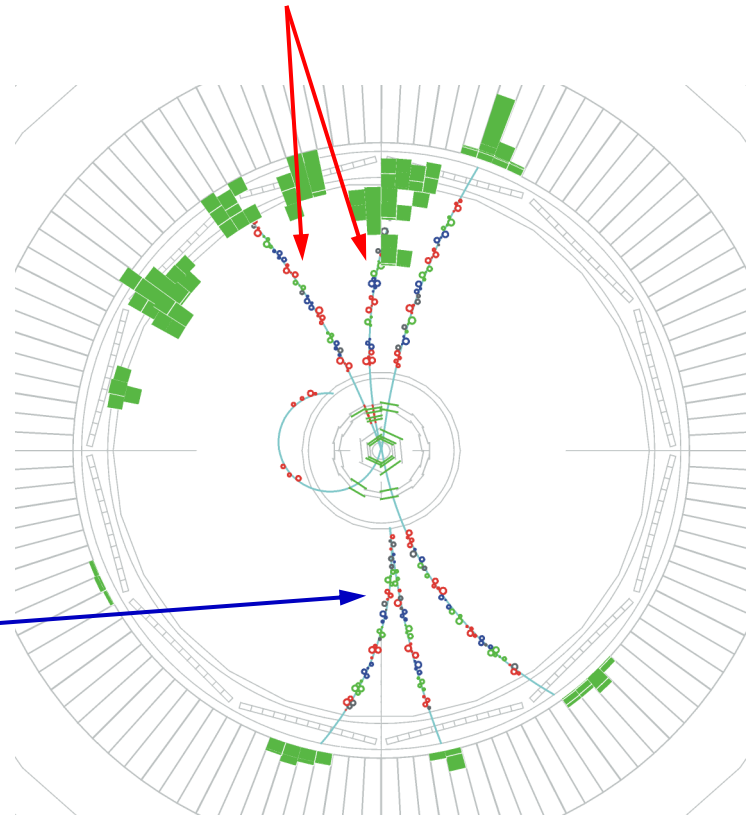
$$B^0 \rightarrow \varphi K_S \text{ at } BaBar$$

only the hits that correspond to reconstructed tracks and neutral candidates are shown (noise hits are removed)



$$\varphi \rightarrow K^+ K^-$$

$$K_S \rightarrow \pi^+ \pi^-$$



Run 29368, event hexID 249a4b/d610dd73 (June 27, 2002)

From the dataset used in *Phys.Rev.D* **69**:011102, 2004



BABAR

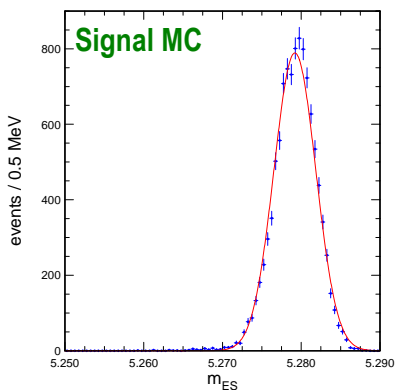


Common discriminating variables: kinematics

A pair of weakly correlated variables that reflect energy and momentum conservation: peaking for fully reconstructed B decays, smooth for combinatorial background

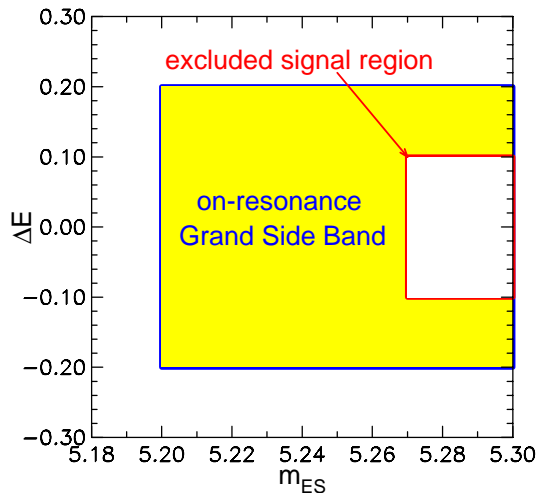
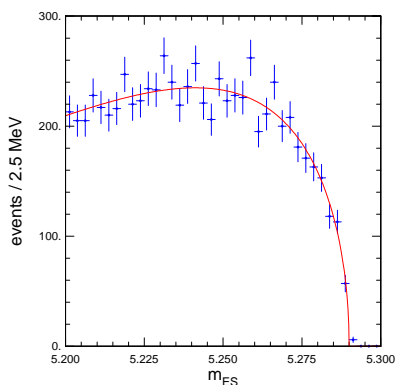
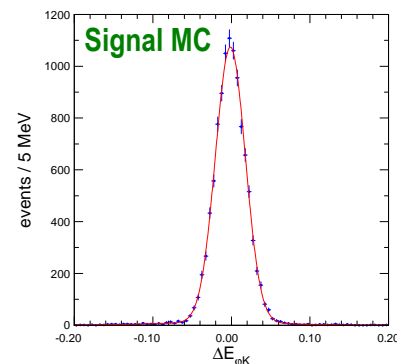
$$m_{ES} \equiv \sqrt{E_{CM \text{ beam}}^2 - p_{CM B}^2} = m_B$$

$$\Delta E \equiv E_{CM B} - E_{CM \text{ beam}} = 0$$

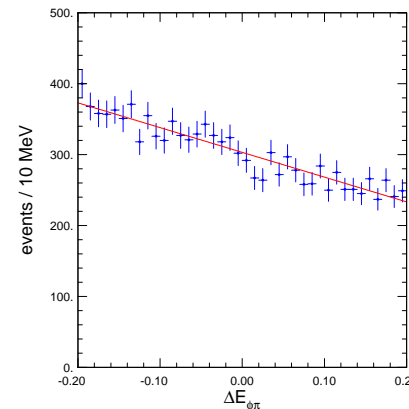


resolution $\sim 2.6 \text{ MeV}/c^2$
determined by the beam energy spread

resolution $\sim 15\text{--}80 \text{ MeV}$
depending on the number of tracks and presence of neutrals in the final state



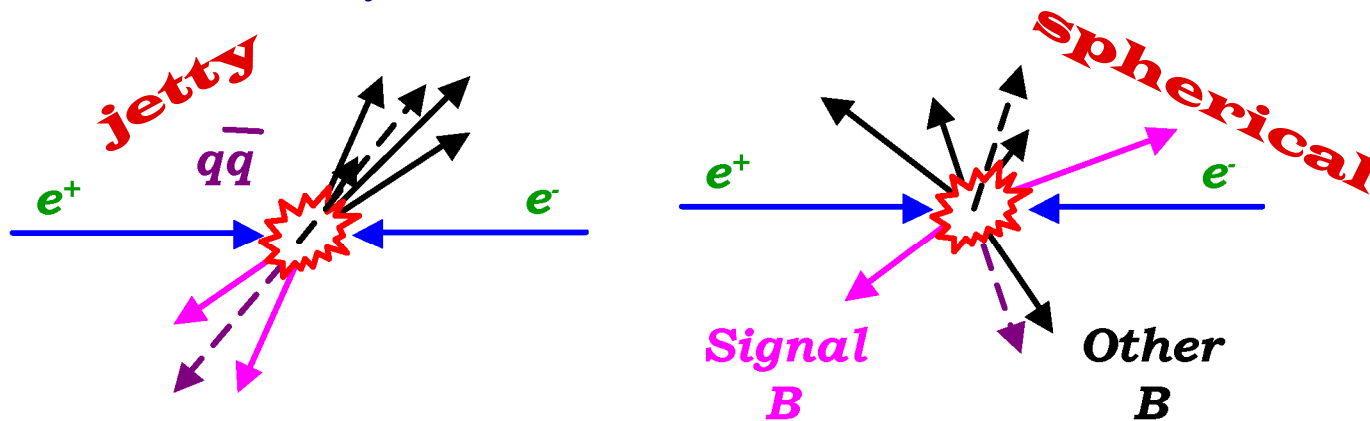
continuum background in the Grand Side Band





Common discriminating variables: event shape

- The principal source of background to rare B decays: random track/neutral combinations from quark-pair ($u\bar{d}sc$) production in the continuum:
 - total $u\bar{d}sc$ cross section ~ 3.4 nb, compared to ~ 1.1 nb for $Y(4S)$
 - $u\bar{d}sc$ events have jet-like topology, while B decays are nearly spherical in CM
 - several topological variables are employed to suppress this background
- Backgrounds from $\tau^+\tau^-$ production and two-photon physics are usually negligible
- Backgrounds from other B decays tend to be small



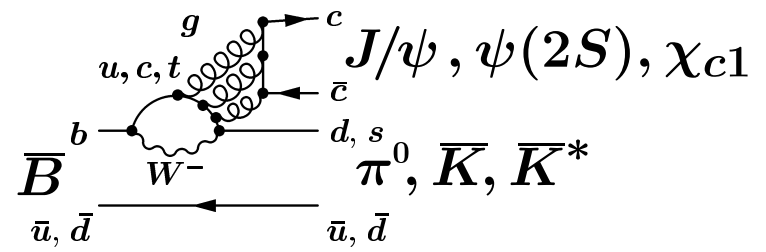
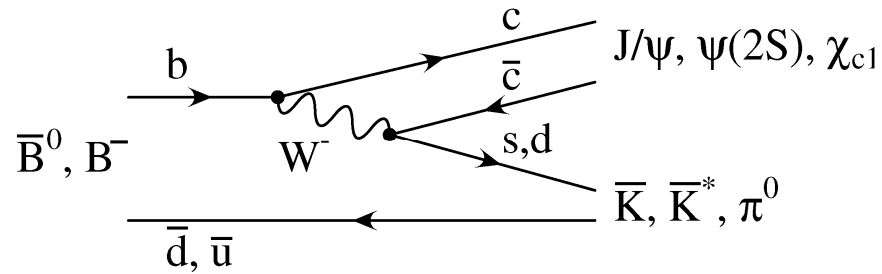
Analyses are blind until the methodology has been finalized and frozen



sin2β in “golden” modes:

The highest-precision test of the KM mechanism of CP violation in meson decays

branching fractions $O(10^{-3})$



“Golden” modes: color-suppressed tree dominates; the t -quark penguin has the same weak phase as the tree. In SM, therefore,

$$S_{\text{golden}} = \eta_{CP} \times \sin 2\beta, \quad C_{\text{golden}} = 0 \quad (\eta_{CP} = \pm 1)$$

Theoretical uncertainties:

- *an example* of a model-independent, data-driven calculation:
 - assuming $SU(3)_{\text{flavor}}$ invariance, use $B^0 \rightarrow J/\psi \pi^0$ data to constrain penguin pollution in $J/\psi K^0 \Rightarrow \Delta S_{J/\psi K^0} = S_{J/\psi K^0} - \sin 2\beta = 0.000 \pm 0.012$
- theoretical estimates of the biases due to u - and c -quark penguins, etc.:
 - $\Delta S_{J/\psi K^0} = S_{J/\psi K^0} - \sin 2\beta \sim O(10^{-3})$ H. Li, S. Mishima, hep-ph/0610120
 - $\Delta S_{J/\psi K^0} = S_{J/\psi K^0} - \sin 2\beta \sim O(10^{-4})$ H. Boos et al., Phys. Rev. D 73, 036006 (2006)

M. Ciuchini, M. Pierini, L. Silvestrini, Phys. Rev. Lett. **95**, 221804 (2005)

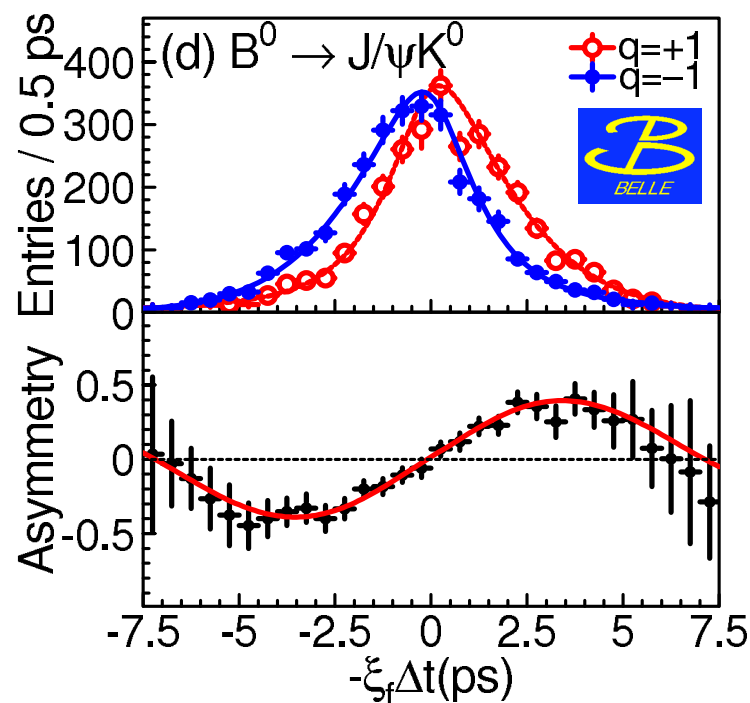
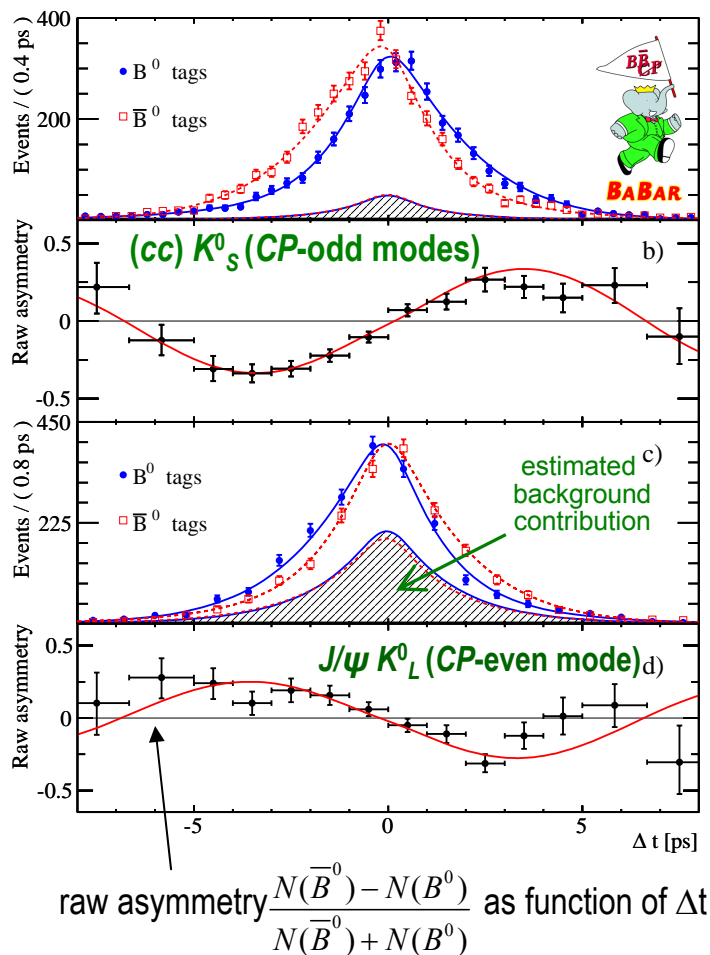


sin2β in “golden” modes: latest from *BaBar* and *Belle*

PRL 99, 171803 (2007)

First observed in 2001

PRL 98, 031802 (2007)



**CP violation in Standard Model is not small, it is $O(1)$.
Smallness of CPV in kaon decays is due to flavor suppression.
CP-violating phases in New Physics can also be $O(1)$.**

BaBar with 384×10^6 *BB* pairs:

$$\sin 2\beta = 0.714 \pm 0.032 \text{ (stat)} \pm 0.018 \text{ (syst)}$$


Belle with 535×10^6 *BB* pairs:

$$\sin 2\beta = 0.642 \pm 0.031 \text{ (stat)} \pm 0.017 \text{ (syst)}$$

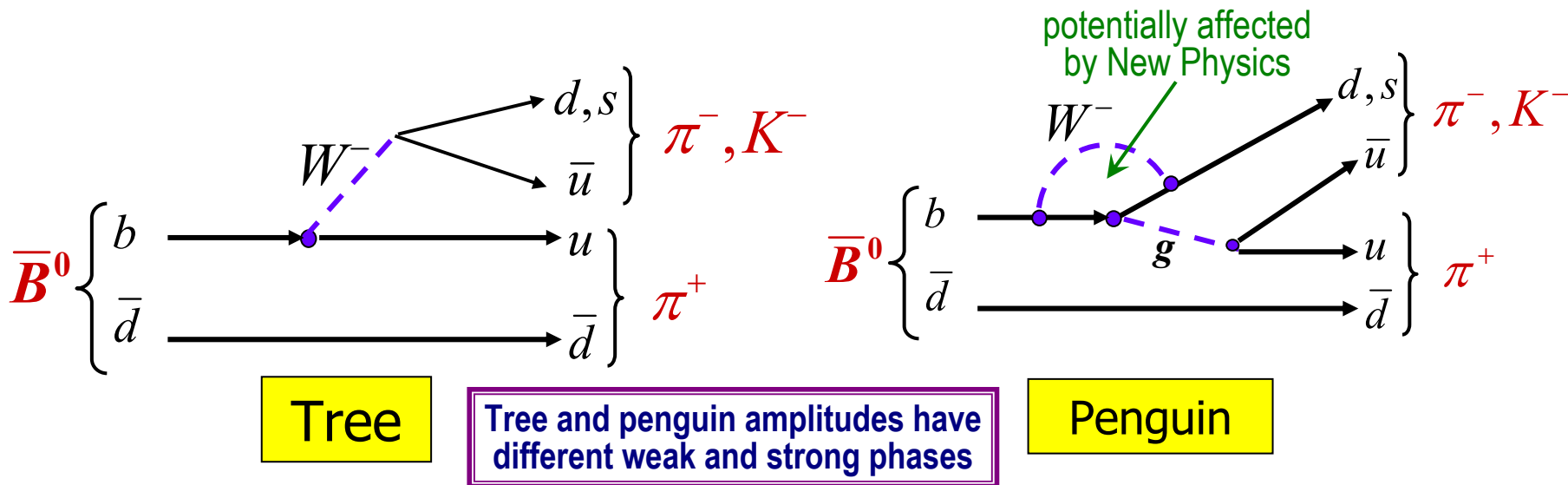


Measuring α with $B^0 \rightarrow \pi^+ \pi^-$

$\mathcal{A}_{CP}(t)$ in $b \rightarrow u\bar{u}d$ decay to a CP eigenstate at the tree level :


 Measure $180^\circ - \beta - \gamma = \alpha \equiv \arg \left[\frac{-V_{td} V_{tb}^*}{V_{ud} V_{ub}^*} \right]$ (in SM)

Penguins: $\mathcal{A}_{CP}(t) \Rightarrow \sin(2\alpha_{\text{eff}})$; $\alpha_{\text{eff}} = \alpha - \Delta\alpha$; **direct $\mathcal{A}_{CP} \neq 0$**



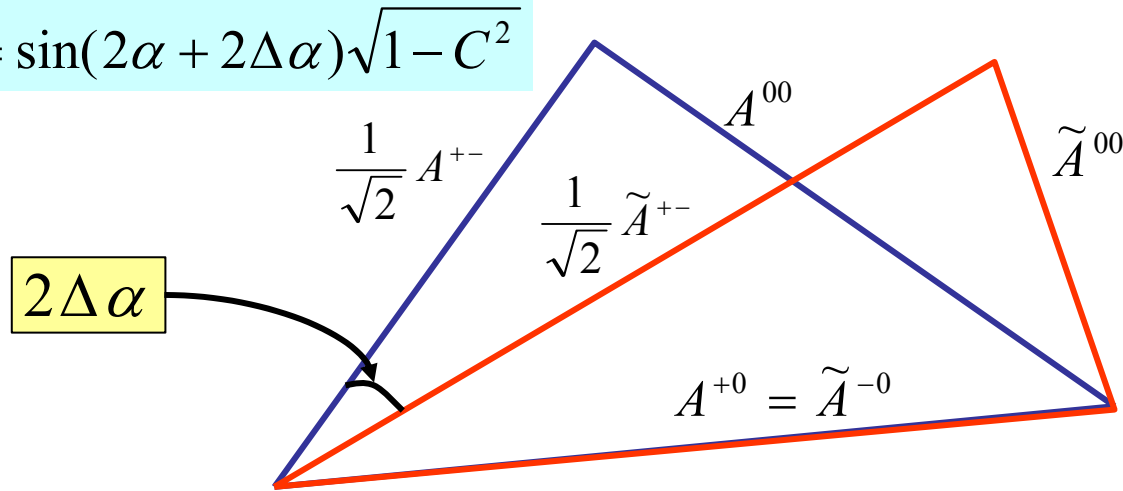


Isospin analysis in $B \rightarrow \pi\pi, \rho\rho$

Determines relative phase between B mixing and the tree, independent of the EW model

M. Gronau, D. London, *Phys. Rev. Lett.* **65**, 3381 (1990)

$$S = \sin(2\alpha + 2\Delta\alpha)\sqrt{1 - C^2}$$



$$A^{+-} = A(B^0 \rightarrow \pi^+ \pi^-)$$

$$\tilde{A}^{+-} = A(\bar{B}^0 \rightarrow \pi^+ \pi^-)$$

$$A^{00} = A(B^0 \rightarrow \pi^0 \pi^0)$$

$$\tilde{A}^{00} = A(\bar{B}^0 \rightarrow \pi^0 \pi^0)$$

$$A^{+0} = A(B^+ \rightarrow \pi^+ \pi^0)$$

$$\tilde{A}^{-0} = A(B^- \rightarrow \pi^- \pi^0)$$

In $B \rightarrow \rho\rho$, there are 3 such relations (one for each polarization)

6 unknowns, 6 observables in $\pi\pi$ (there is no vertex to measure $S_{\pi^0\pi^0}$)

5 observables in $\rho\rho$ (or 7, when both $C_{\rho^0\rho^0}$ and $S_{\rho^0\rho^0}$ are measured)

4-fold ambiguity in $2\Delta\alpha$: either triangle can flip up or down

$$A_{hh} = e^{+i\gamma}T + e^{-i\beta}P$$

$$\tilde{A}_{hh} = e^{-i\gamma}T + e^{+i\beta}P$$

Neglecting EW penguins, ± 0 is a pure tree mode, and so the two triangles share a common side:

$$A(B^+ \rightarrow h^+ h^0) = \tilde{A}(B^- \rightarrow h^- h^0)$$

$$A^{+0} = \frac{1}{\sqrt{2}} A^{+-} + A^{00}$$

$$\tilde{A}^{-0} = \frac{1}{\sqrt{2}} \tilde{A}^{+-} + \tilde{A}^{00}$$

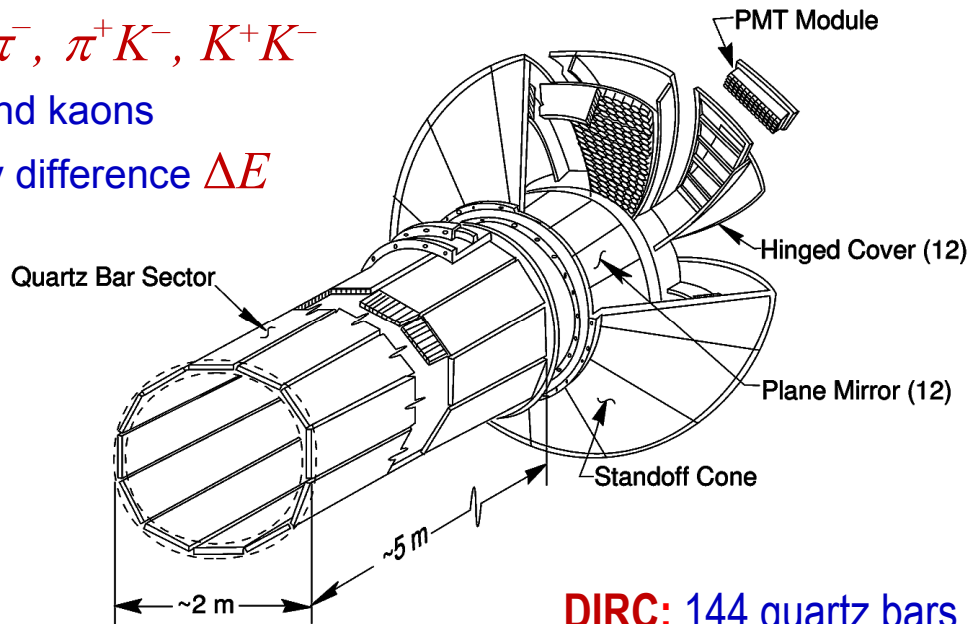
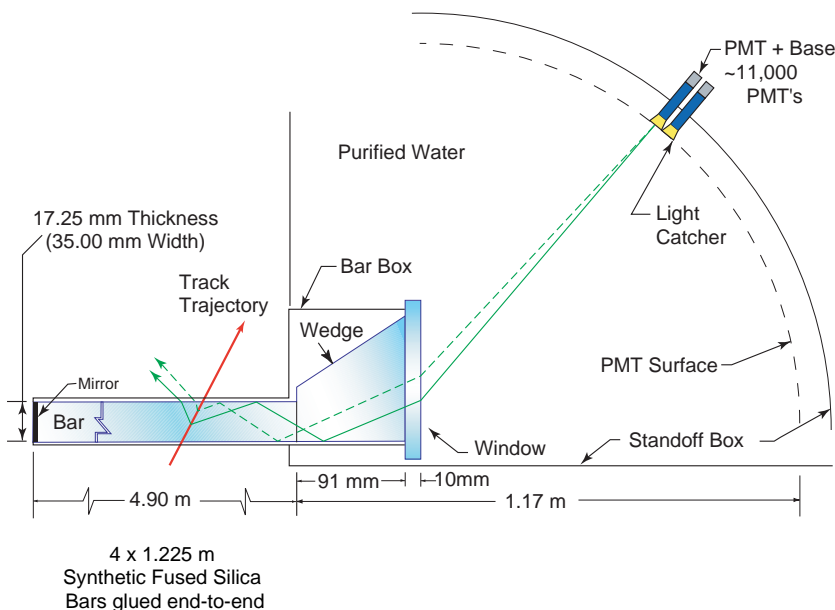


The "classic" $B^0 \rightarrow h^+h^-$ analysis

Simultaneous ML fit to $B^0 \rightarrow \pi^+\pi^-, K^+\pi^-, \pi^+K^-, K^+K^-$

Using DIRC Cherenkov angle to identify pions and kaons

Additional $\pi\pi/K\pi/KK$ separation from energy difference ΔE



DIRC: 144 quartz bars

0.84 x 4π coverage; 91% eff. with $n_\gamma > 5$

13σ π/K separation at 1.5 GeV/c, 2.5σ at 4.5 GeV/c

Calibration sample: $B^- \rightarrow \pi^- D^0, D^0 \rightarrow \pi^+ K^-$

Belle: π/K separation with aerogel (threshold Cherenkov) and dE/dx ; ΔE

Due to higher particle ID acceptance, Belle had x1.4 more h^+h^- per fb^{-1} than BaBar

CDF: dE/dx and B mass; $\mathcal{A}_{K^+\pi^-}$ only



π/K separation with DCH dE/dx :

Catching up with Belle's $B^0 \rightarrow h^+h^-$ reconstruction efficiency

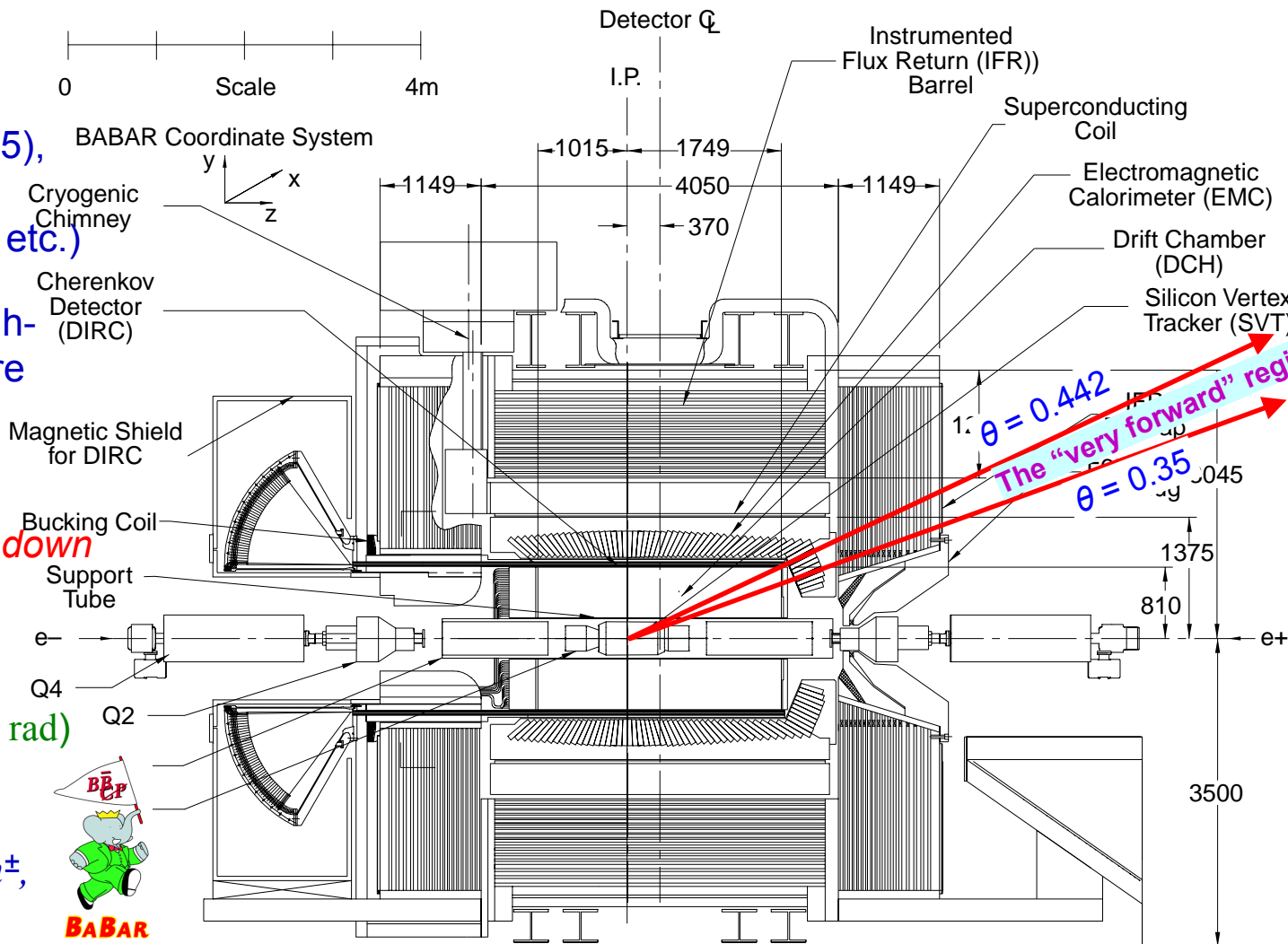
In the barrel ($\theta > 0.445$),
the DIRC is $\sim 9.3\%$
inefficient (ϕ cracks, etc.)

Another $\sim 12\%$ of high-
momentum tracks are
outside the DIRC
acceptance in θ .

We use DCH tracks down
to $\theta = 0.35$ rad

($J/\psi \rightarrow \mu^+\mu^-$ in $\sin 2\beta$
analysis: down to 0.30 rad)

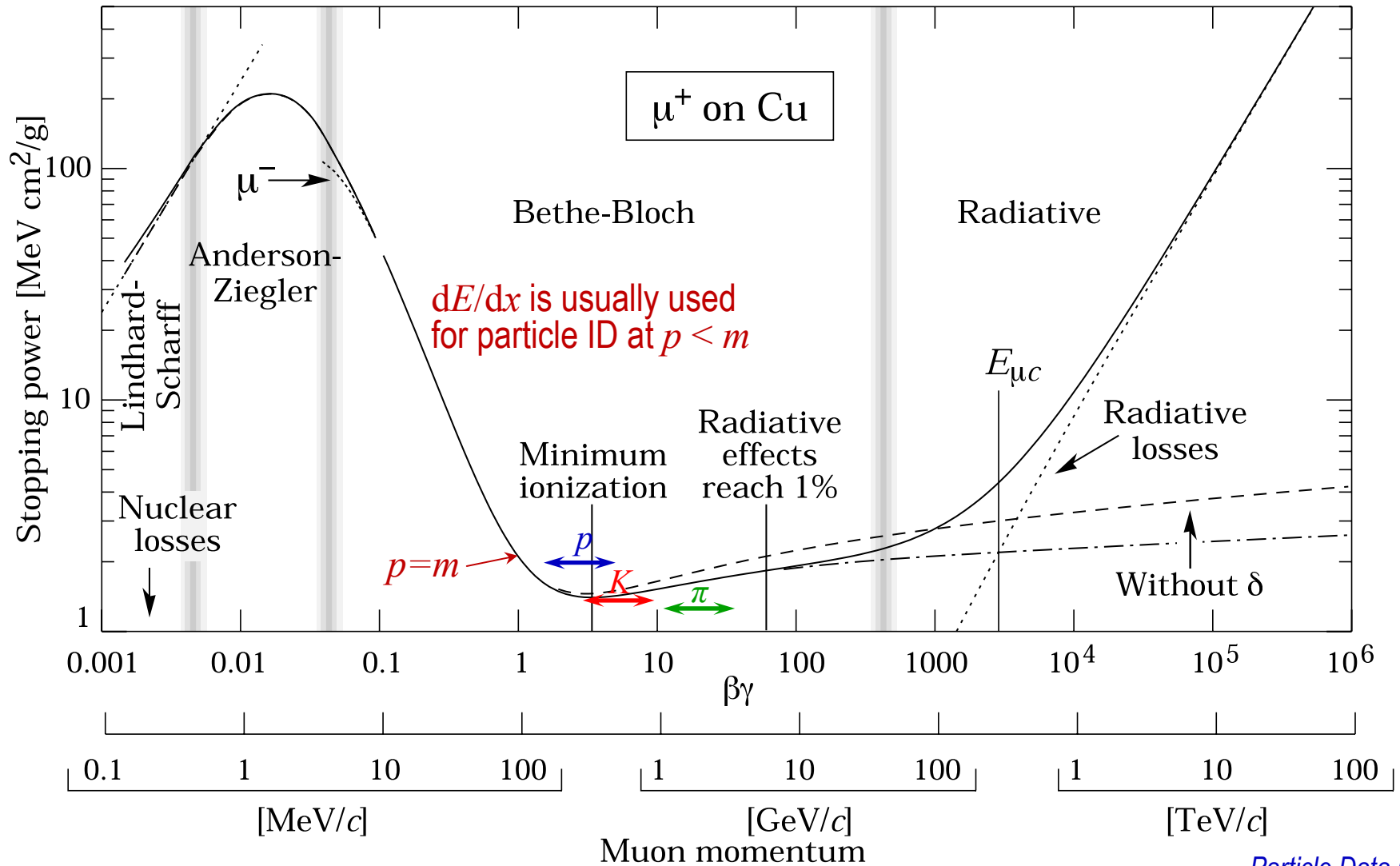
$\rightarrow 16\%$ event-yield
increase for $B \rightarrow Xh^\pm$,
35% for $B^0 \rightarrow h^+h^-$



3-2001
8583A50



Ionization energy loss for $B \rightarrow Xh^\pm$



Particle Data Group

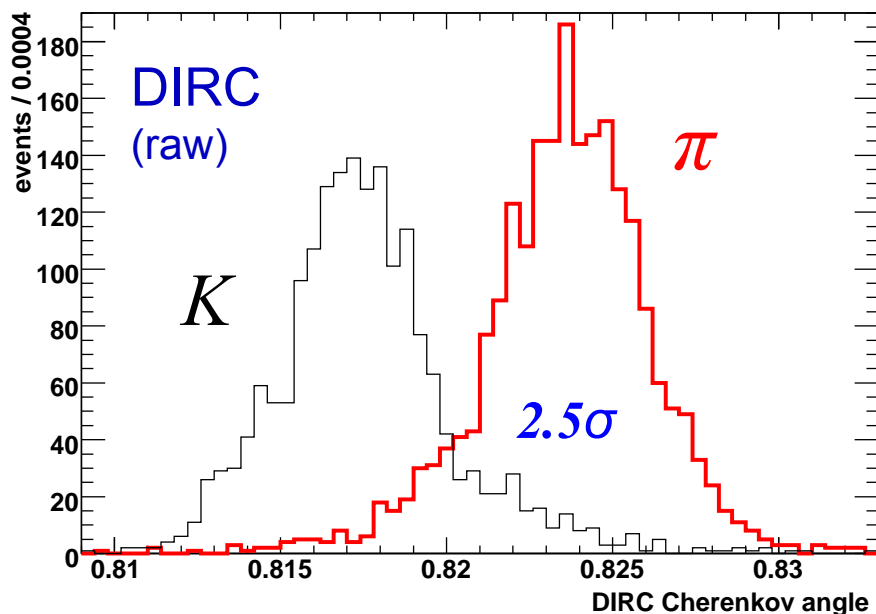


DCH dE/dx vs. the DIRC

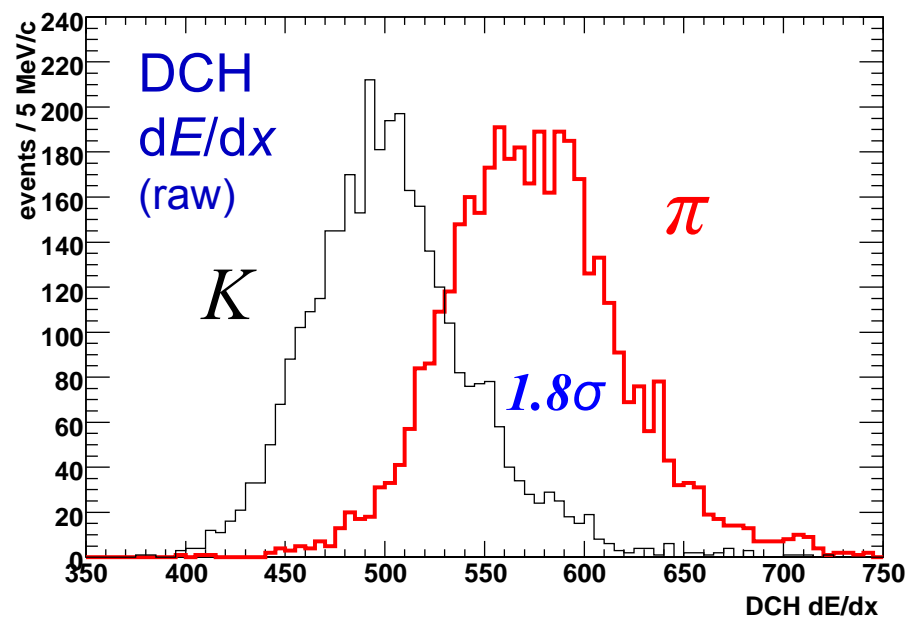
In the forward region, DCH dE/dx is not much worse than the DIRC—and is 100% efficient!

$0.4 < \theta < 0.7$, $3.8 < p < 4.2$ GeV/c
1999-2006 data

DIRC Cherenkov angle, pions and kaons, charge +1, $0.4 < \theta < 0.7$, $3.8 < p < 4.2$ GeV/c, data



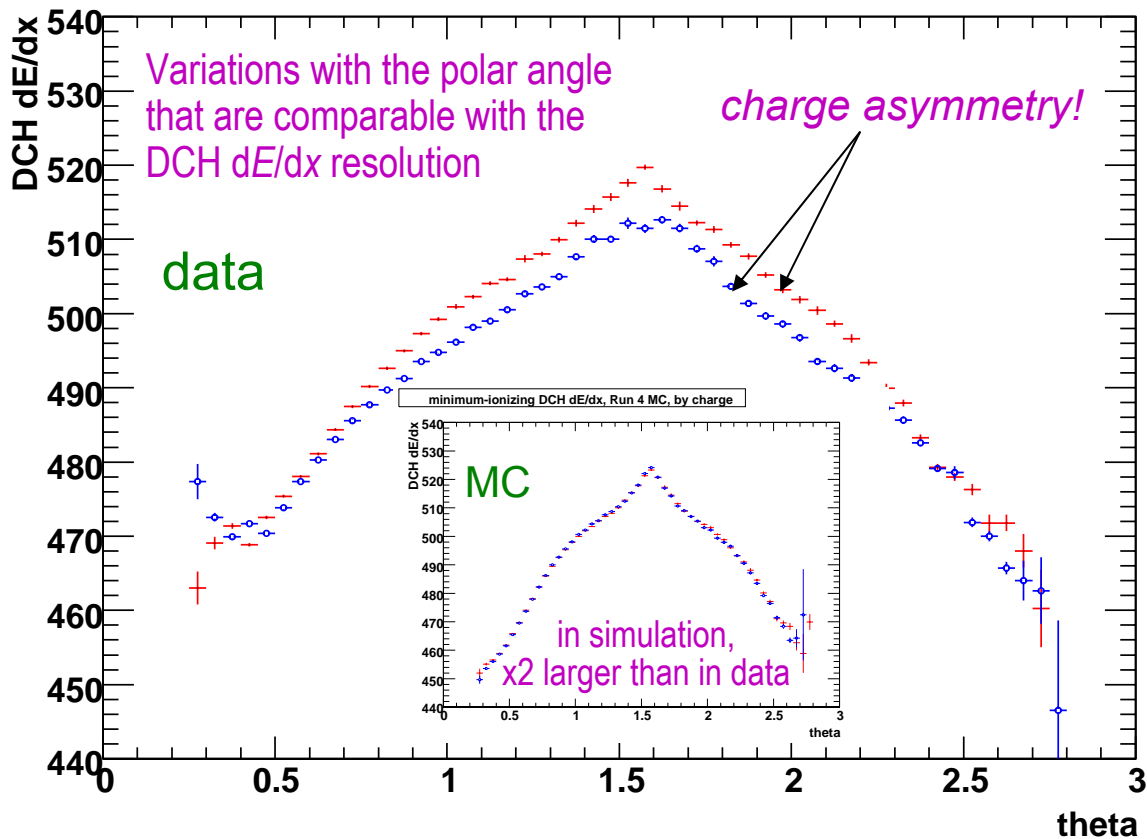
DCH dE/dx , pions and kaons, charge +1, $0.4 < \theta < 0.7$, $3.8 < p < 4.2$ GeV/c, data



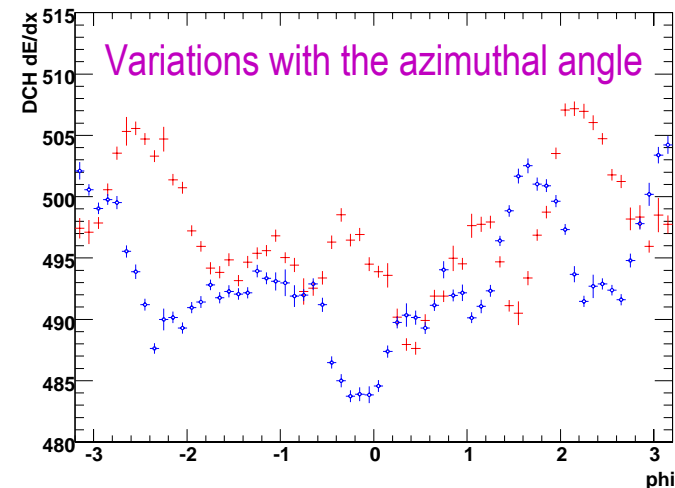


There were many reasons DCH dE/dx failed to work in likelihood-based $B \rightarrow Xh^\pm$ analyses in the past

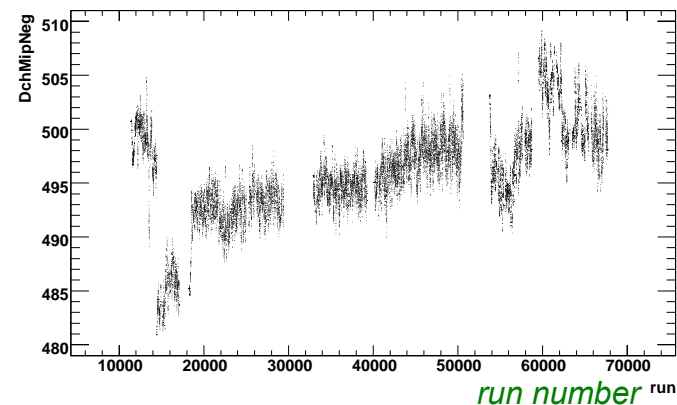
minimum-ionizing DCH dE/dx , Run 4 data, by charge



minimum-ionizing DCH dE/dx , Run 5a1 data, by charge



Minimum-ionizing DCH dE/dx vs. time

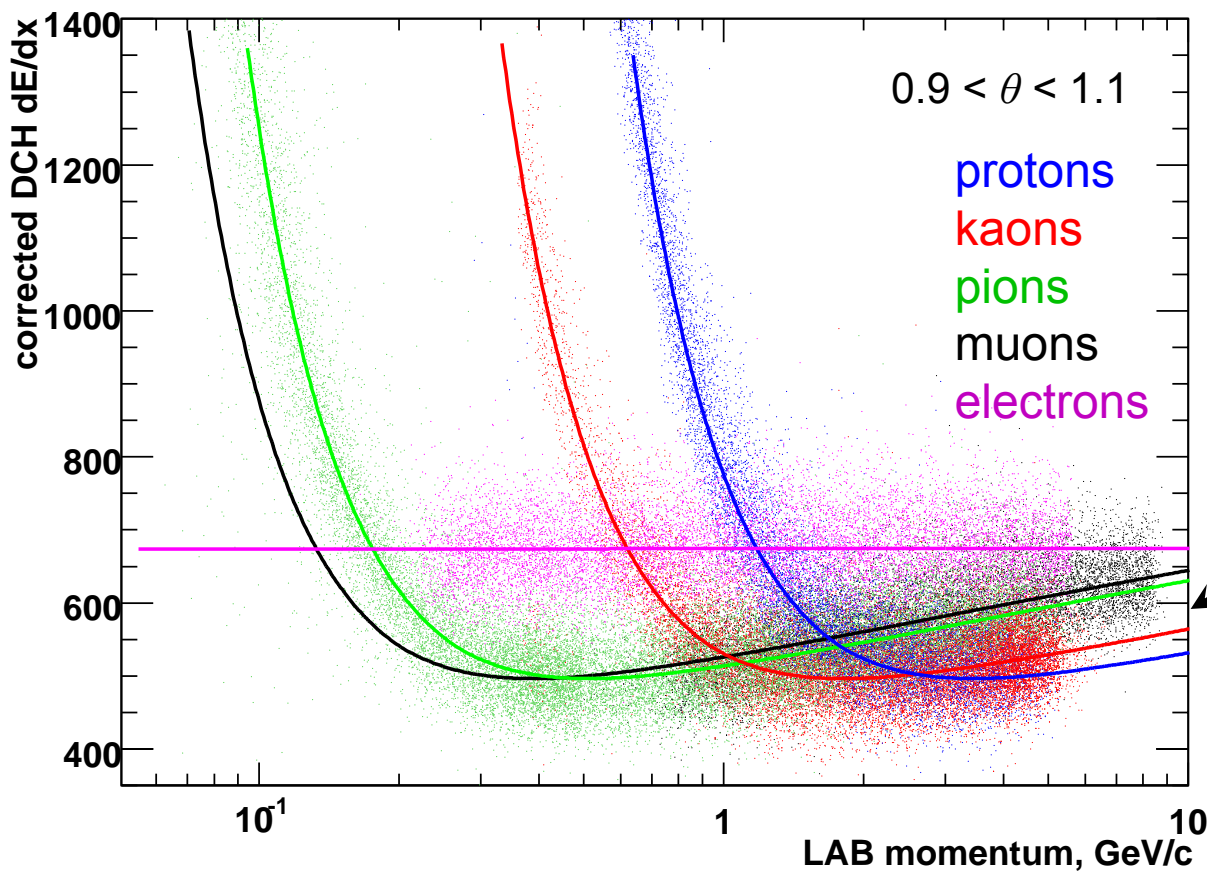




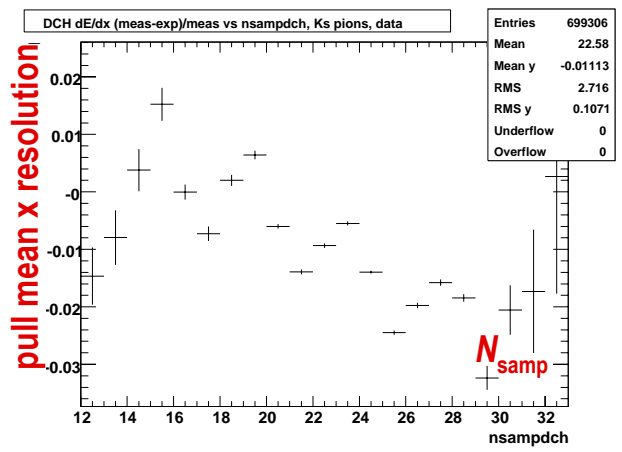
dE/dx in DCH

This is what DCH dE/dx looks like after the new calibration

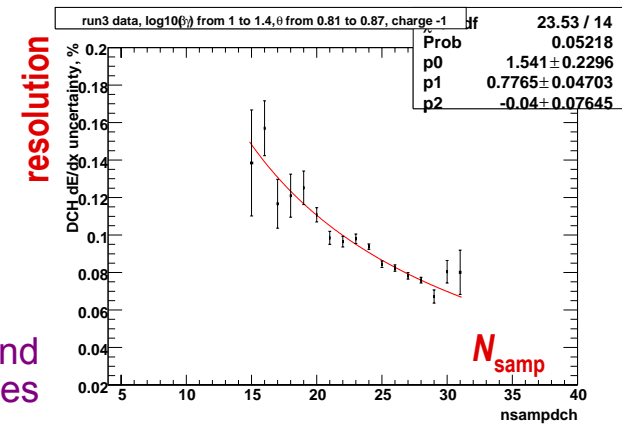
DCH dE/dx for various particle types, Run 3, data



DCH dE/dx mean and resolution also depend on the number of dE/dx samples



up to 1.9σ
 π/K separation



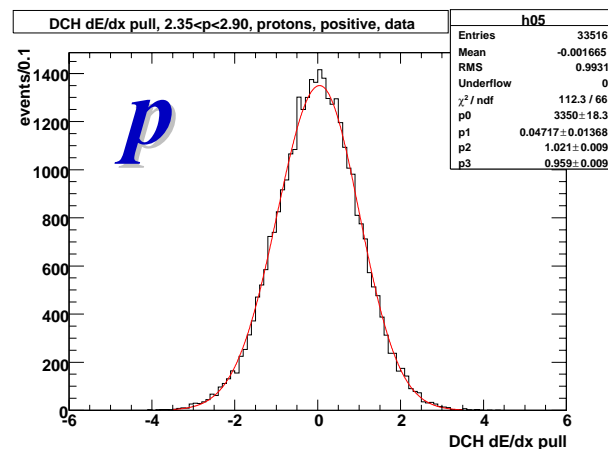
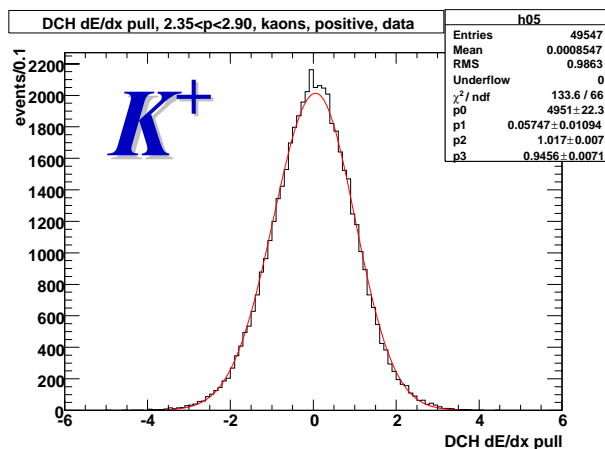
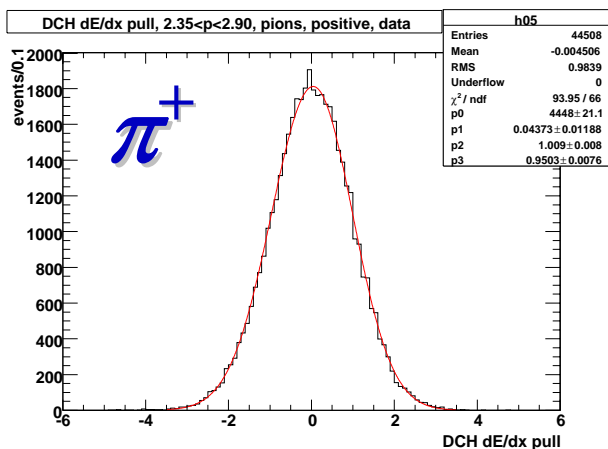


New, detailed DCH dE/dx parametrization:

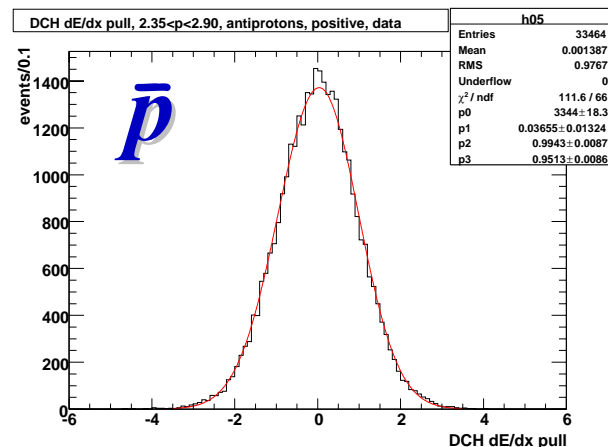
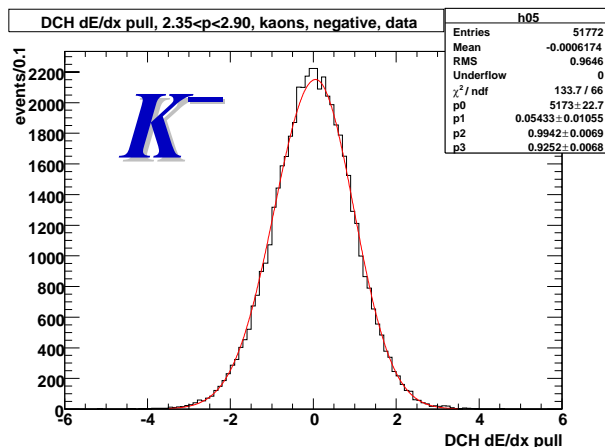
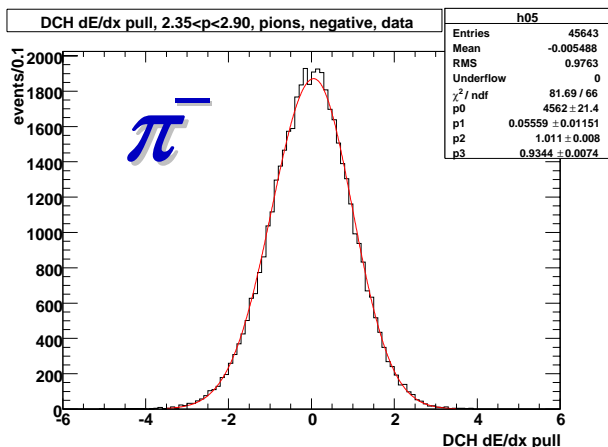
π, K, p pulls for $B \rightarrow Xh^\pm$ track momenta

$\text{pull} \equiv (\text{measured} - \text{expected}) / \text{error}$

“measured” values in these plots are taken from appropriate calibration samples

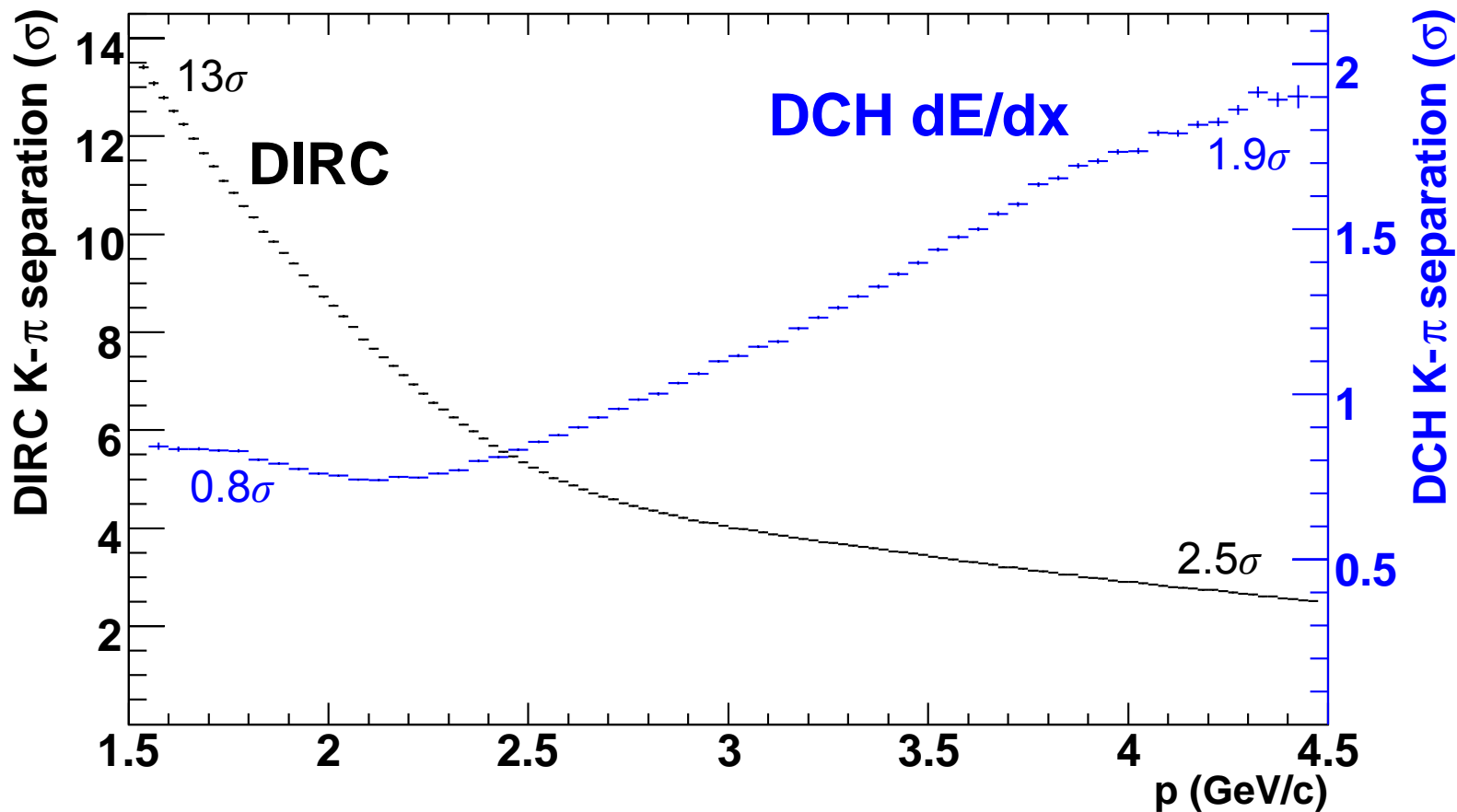


Pulls are controlled at a <1% level; non-Gaussian tails are absent by construction





DCH dE/dx K - π separation in $B \rightarrow Xh^\pm$ complementary to DIRC



(for tracks that have good DIRC information, we use both DIRC and dE/dx)

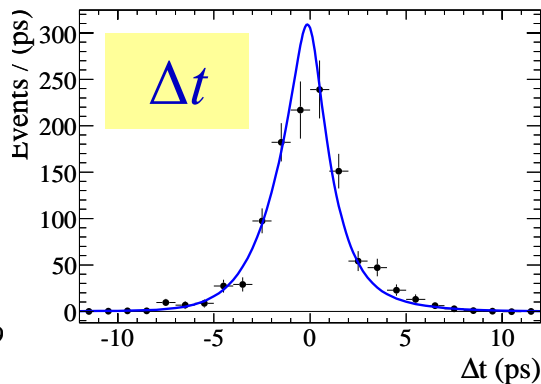
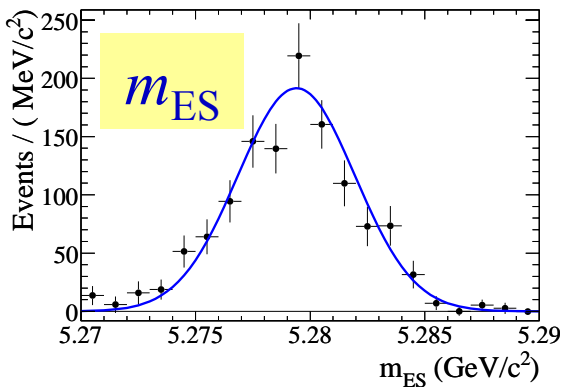


Our new result: $B^0 \rightarrow \pi^+ \pi^-$

$$N_{\pi^+ \pi^-} = 1139 \pm 49$$

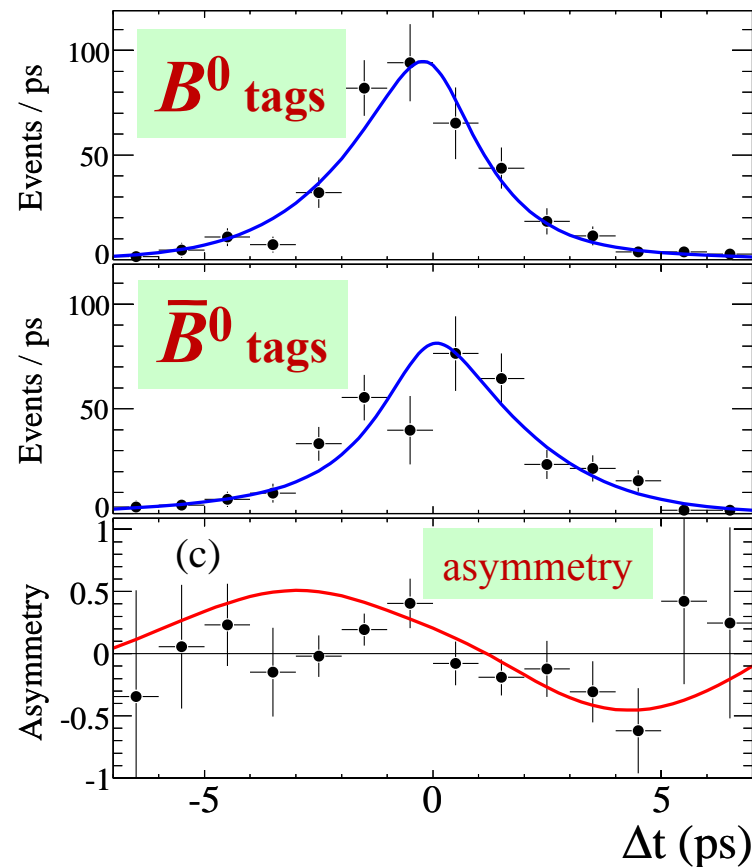
52% overall increase
in signal reconstruction efficiency:

35% from DCH dE/dx ,
13% mainly from reoptimizing the event-shape $|\cos\theta_{\text{spher}}|$ cut



sPlot:

Builds a histogram of x excluding it from the Maximum-Likelihood fit, assigning a weight to each event, keeping all signal events, getting rid of all background events, and keeping track of the statistical errors in each x bin



M. Pivk and F. R. Le Diberder,
“sPlot: a statistical tool to unfold data distributions,”
Nucl. Instrum. Meth. A **555**, 356 (2005)
[arXiv:physics/0402083]



Observation of CP Violation in $B^0 \rightarrow K^+ \pi^-$ and $B^0 \rightarrow \pi^+ \pi^-$

BaBar has made a **5.4 σ** observation of **CP** violation in $B^0 \rightarrow \pi^+ \pi^-$

$$S_{\pi\pi} = -0.60 \pm 0.11 \pm 0.03 \text{ (5.1}\sigma\text{)}$$

$$C_{\pi\pi} = -0.21 \pm 0.09 \pm 0.02 \text{ (2.3}\sigma\text{)}$$

BaBar: 383 million $B\bar{B}$ pairs, $1139 \pm 49 \pi^+ \pi^-$

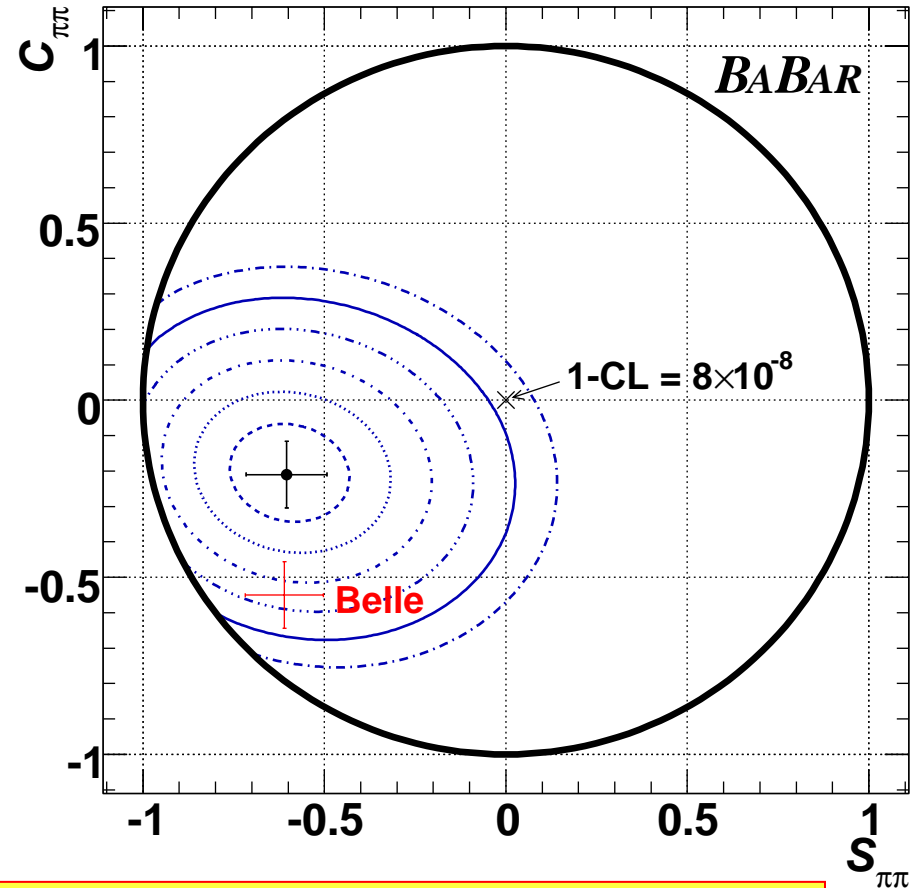
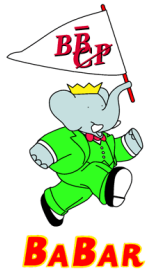
Belle: 535 million $B\bar{B}$ pairs, $1464 \pm 65 \pi^+ \pi^-$

BaBar now has 9% more $\pi^+ \pi^-$ per fb^{-1} than **Belle**

Belle PRL 98, 211801 (2007)

$$S_{\pi\pi} = -0.61 \pm 0.10 \pm 0.04 \text{ (5.3}\sigma\text{)}$$

$$C_{\pi\pi} = -0.55 \pm 0.08 \pm 0.05 \text{ (5.5}\sigma\text{)}$$



also:

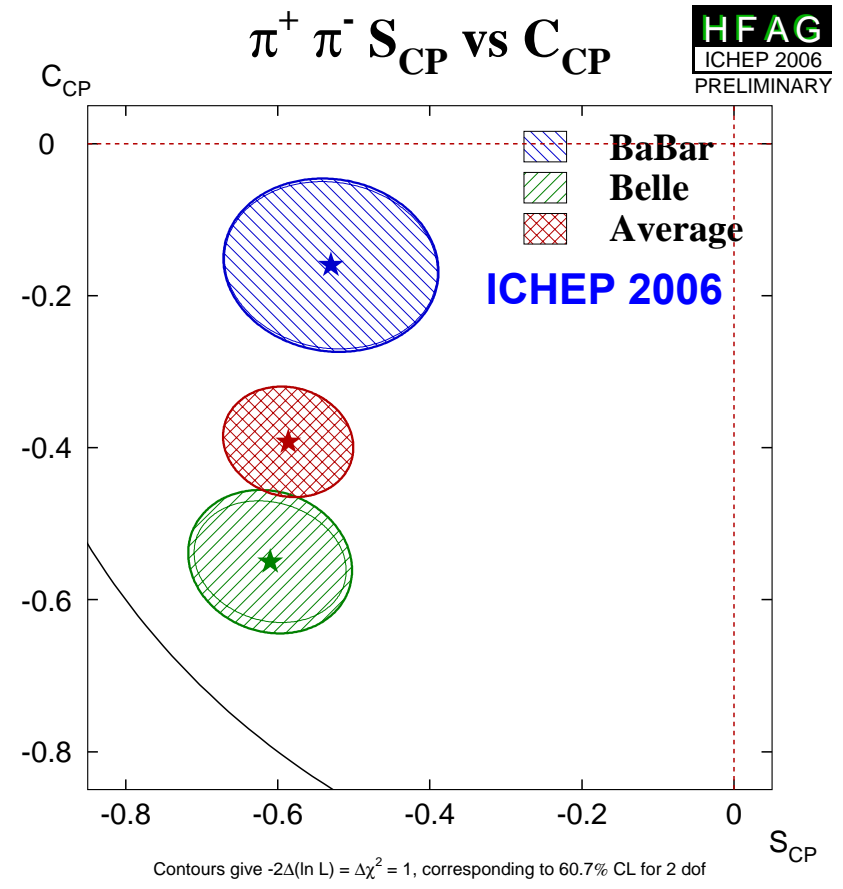
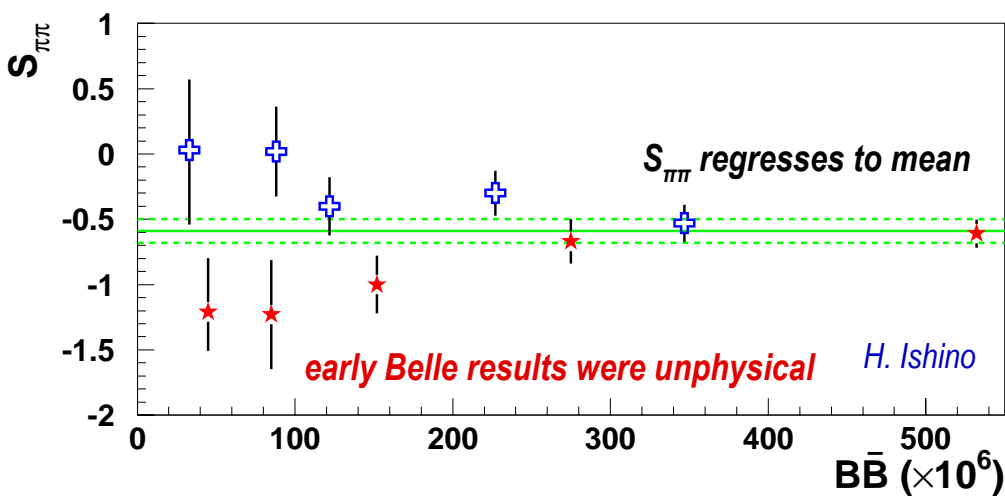
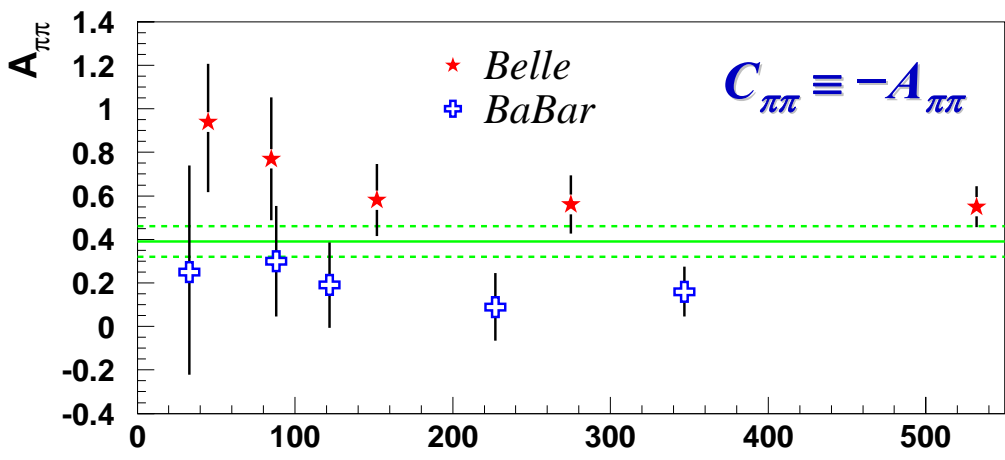
$$\mathcal{A}_{K^+ \pi^-} = -0.107 \pm 0.018^{+0.007}_{-0.004} \text{ (5.5}\sigma\text{)}$$



History of $B^0 \rightarrow \pi^+ \pi^-$ results

There exists a history of disagreement between Belle and *BaBar* on $B^0 \rightarrow \pi^+ \pi^-$

A 2.6σ difference is still present, in $C_{\pi\pi}$





The $B \rightarrow \pi^\pm \pi^0, \pi^0 \pi^0$ analysis

Simultaneous fit to $B^0 \rightarrow \pi^+ \pi^0, K^+ \pi^0$ (using DIRC Cherenkov angle to separate pions and kaons)

$B^0 \rightarrow \pi^0 \pi^0$: branching fraction and time-integrated direct CP asymmetry

new: in addition to $\pi^0 \rightarrow \gamma\gamma$, we use merged π^0 and $\gamma \rightarrow e^+e^-$ conversions

\Rightarrow 10% efficiency increase per π^0 (4% from merged π^0 , 6% from γ conversions)

At a Super B -meson factory, $\gamma \rightarrow e^+e^-$ conversions would make $S_{\pi^0 \pi^0}$ determination possible!

merged π^0 :

the two photons are too close to one another in the EMC to be reconstructed individually; can be recovered using

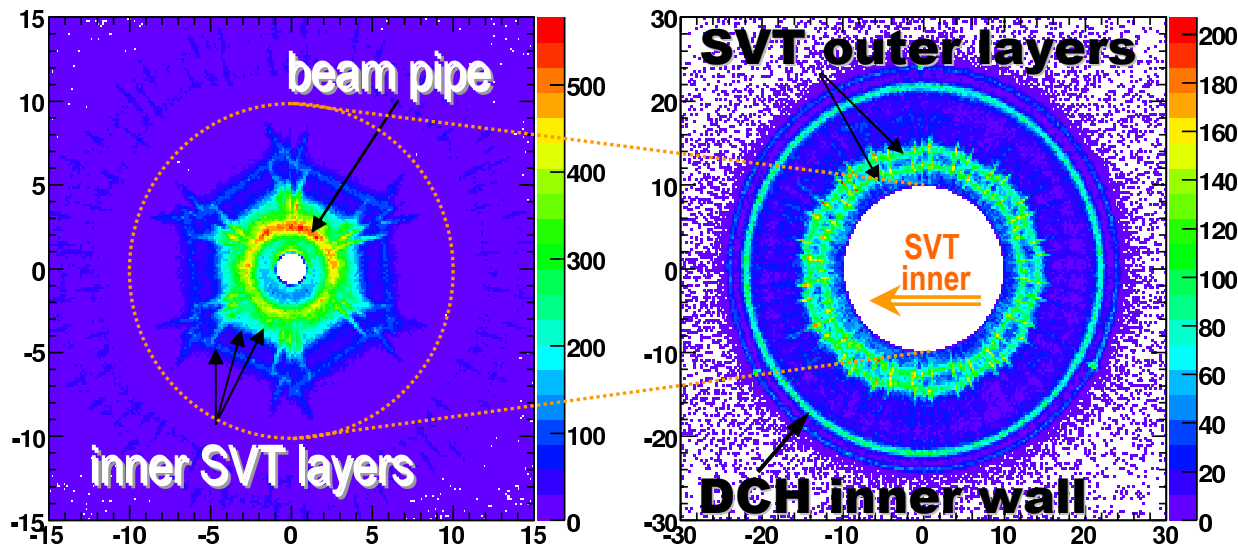
$$M_{\pi^0}^2 \approx E_{\pi^0}^2 (S_{\pi^0} - S_\gamma),$$

where S is the second EMC moment of the merged $\pi^0 \rightarrow \gamma\gamma$

The control sample: $\tau \rightarrow \rho\nu$

$\gamma \rightarrow e^+e^-$ conversions:

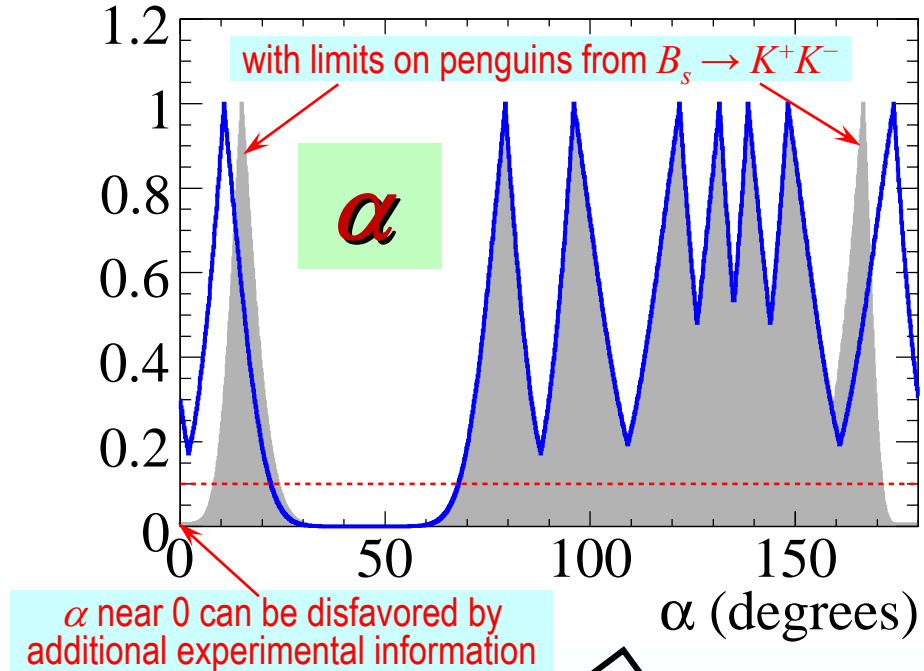
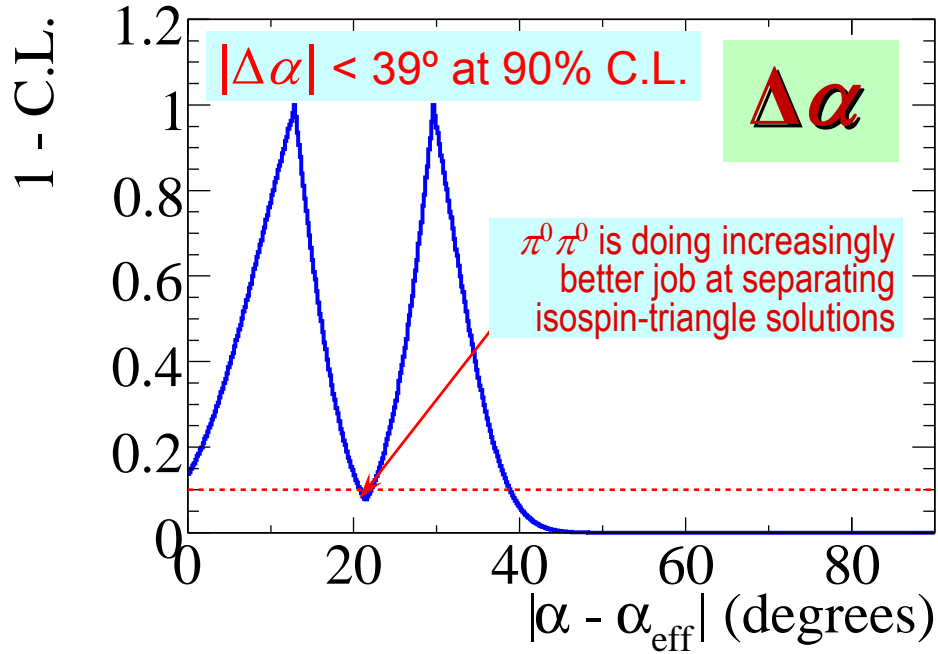
result from interactions with detector elements





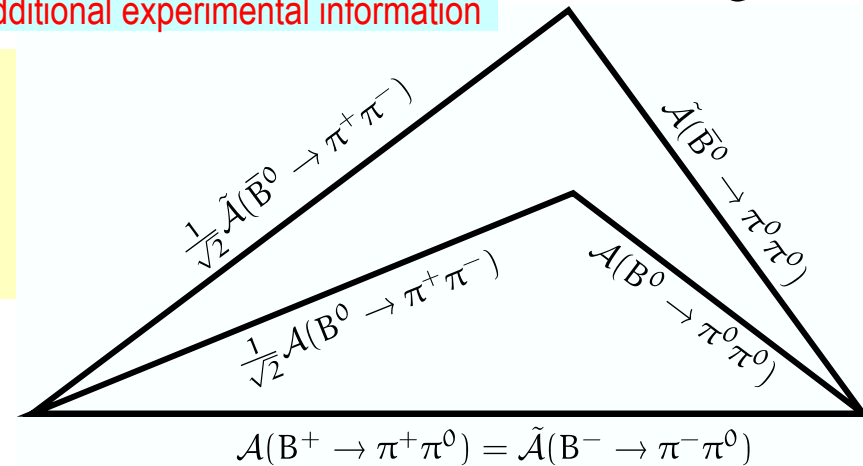
An interpretation of our new $B \rightarrow \pi\pi$ results

PRD 76, 091102 (2007)



This is a frequentist interpretation: we use only the $B \rightarrow \pi\pi$ isospin-triangle relations in arriving at these constraints on $\Delta\alpha = \alpha - \alpha_{\text{eff}}$ and on α itself

Here is one of the possible solutions to the Gronau-London isospin triangle in $B \rightarrow \pi\pi$ according to the central values of the Summer 2006 BaBar results:





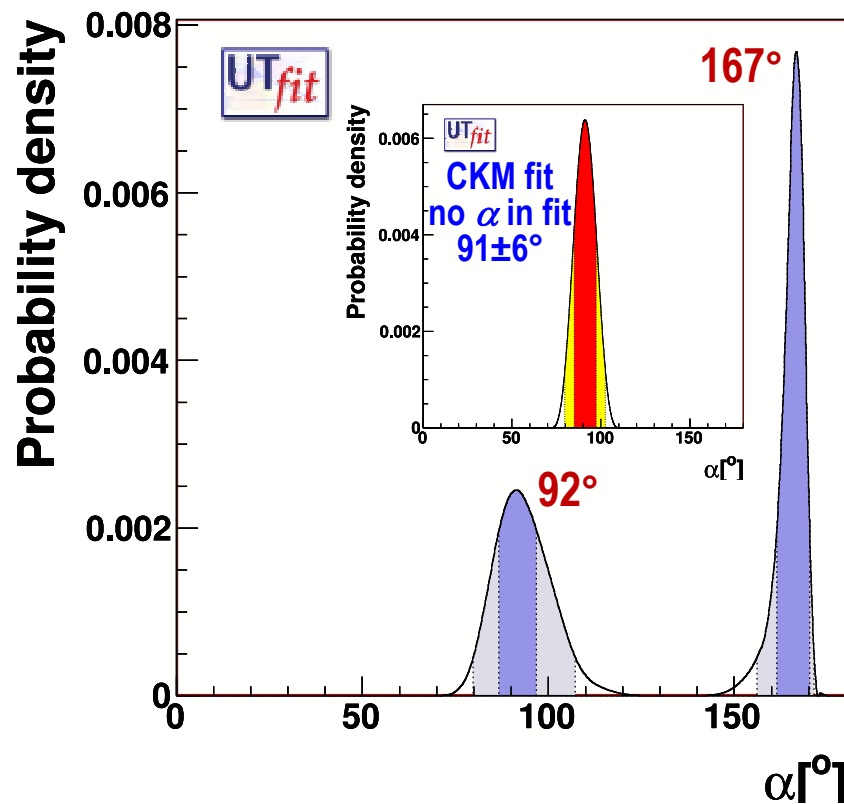
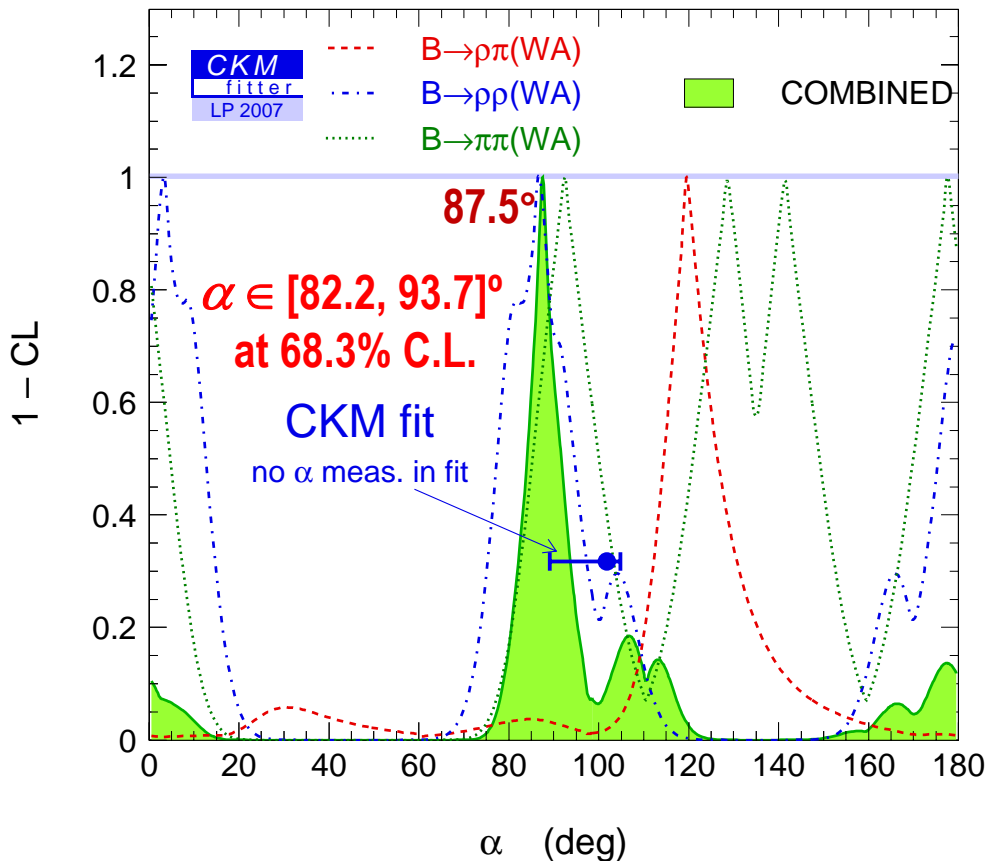
Global fits for the value of α

The second-highest-precision test of the KM mechanism of CP violation in meson decays

Two interpretations currently exist that convert the $B \rightarrow \pi\pi, \rho\pi, \rho\rho$ measurements to constraints on α :

A frequentist interpretation

A Bayesian interpretation, with model-dependent choices of priors



CKMfitter Group (J. Charles et al.), Eur. Phys. J. C41, 1-131 (2005), [hep-ph/0406184], updated results and plots available at <http://ckmfitter.in2p3.fr>

M. Ciuchini, G. D'Agostini, E. Franco, V. Lubicz, G. Martinelli, F. Parodi, P. Roudeau, A. Stocchi, JHEP 0107 (2001) 013 [hep-ph/0012308], updated results and plots available at <http://www.utfit.org>



Our new result: $B^0 \rightarrow K^+ \pi^-$

the first $>5\sigma$ observation of $\mathcal{A}_{K^+ \pi^-}$ (direct CP violation)

BaBar:

$$\mathcal{A}_{K^+ \pi^-} = -0.107 \pm 0.018^{+0.007}_{-0.004}$$

published in
PRL 99, 021603 (2007)

Belle:

$$\mathcal{A}_{K^+ \pi^-} = -0.093 \pm 0.018 \pm 0.008$$

ICHEP'06 preliminary

CDF:

$$\mathcal{A}_{K^+ \pi^-} = -0.086 \pm 0.023 \pm 0.009$$

DPF'06 preliminary

$$\mathcal{A}_{K^+ \pi^-}(\text{World Average}) = -0.097 \pm 0.012 \text{ (} 8\sigma \text{ from 0)}$$

$$\text{Predicted to be } \approx \mathcal{A}_{K^+ \pi^0}(\text{WA}) = +0.051 \pm 0.025$$

5.4 σ difference: an “ $\mathcal{A}_{K\pi}$ puzzle”

– Could be due to hadronic effects, not New Physics
(e.g., see R. Fleischer, hep-ph/0608010, hep-ph/0701217)

assuming unbroken $SU(3)_{\text{flavor}}$:

$$C_{\pi\pi} = 3 \times \mathcal{A}_{K^+ \pi^-}$$

with factorizable $SU(3)_{\text{flavor}}$ -breaking corrections:

$$C_{\pi\pi} = 2.5 \times \mathcal{A}_{K^+ \pi^-} = -0.24$$



$\mathcal{A}_{K^+\pi^-}$ bias due to material interactions determined and cross-checked using several independent approaches

- 1) Detailed GEANT4 v7.1-based simulation: bias of -0.005
- 2) From material accounting and material properties + cross sections tabulated in PDG-RPP

Material type	comment	density	λ_T	thickness	in λ_T	$\int \lambda_T$ from IP
Au	beam pipe	19.3	113.9	4 μm	0.07×10^{-3}	
Be	beam pipe	1.848	55.8	1.36 mm	4.50×10^{-3}	
H ₂ O	beam pipe	1.000	60.1	1.48 mm	2.46×10^{-3}	0.703%
Si	SVT modules	2.33	70.6	1.7 mm	5.61×10^{-3}	
Kapton + glue	SVT fanouts	1.4	60.3	0.50 mm	1.16×10^{-3}	
Cu + Cr	SVT fanouts	9.0	85.6	24 μm	0.25×10^{-3}	
Au	SVT fanouts	19.3	113.9	5 μm	0.09×10^{-3}	
Air	SVT	0.001205	62.0	20 cm	0.39×10^{-3}	
C	support tube	2.265	60.2	1.5 mm	5.57×10^{-3}	
Be	inner DCH wall	1.848	55.8	1.00 mm	3.31×10^{-3}	2.34%
80% He, 20% C ₄ H ₁₀	25 cm of DCH	0.000615	51.2	25 cm	0.30×10^{-3}	
Total (IP to DCH)	90° GTL track					2.37% λ_T
80% He, 20% C ₄ H ₁₀	the rest of DCH	0.000615	51.2	32 cm	0.38×10^{-3}	
C	DCH outer wall	2.265	60.2	3.8 mm	14.3×10^{-3}	
Al	DCH outer wall	2.70	70.6	125 μm	0.48×10^{-3}	
Al	DRC before SiO ₂	2.70	70.6	3.2 mm	12.3×10^{-3}	
Total IP to DIRC	90° track					5.1% λ_T

- 3) Asymmetry in the continuum background (uncorrected): -0.011 ± 0.004 (stat)

$\mathcal{A}_{K^+\pi^-}$ bias correction: $+0.005^{+0.006}_{-0.003}$

← conservative
← bias cannot be zero



Summary and outlook

By gaining a better, more detailed and precise understanding of BaBar detector performance, we have:

- Found ways to increase the $B^0 \rightarrow h^+h^-$ reconstruction efficiency by 52%
- Observed CP violation the time distribution of $B^0 \rightarrow \pi^+\pi^-$ decays
- Observed CP violation the charge asymmetry of $B^0 \rightarrow K^+\pi^-$ decays
The first single-experiment 5-sigma observation of direct CP violation in a process where direct CP violation is not tightly coupled to another type of CP violation (as in ε' or $C_{\pi\pi}$). CPV in decay rules out “superweak” models.

The Kobayashi-Maskawa phase has been demonstrated to be the dominant source of CP violation in meson decays.

Constraints on the CKM Unitarity Triangle from angle measurements are comparable with constraints from all other sources – and mutually consistent.

The absence of a statistically significant incompatibility between them contains a wealth of information about the New Physics that we can be expected at the TeV scale. In particular, it pushes up the scale of new CP -violating physics in most models.

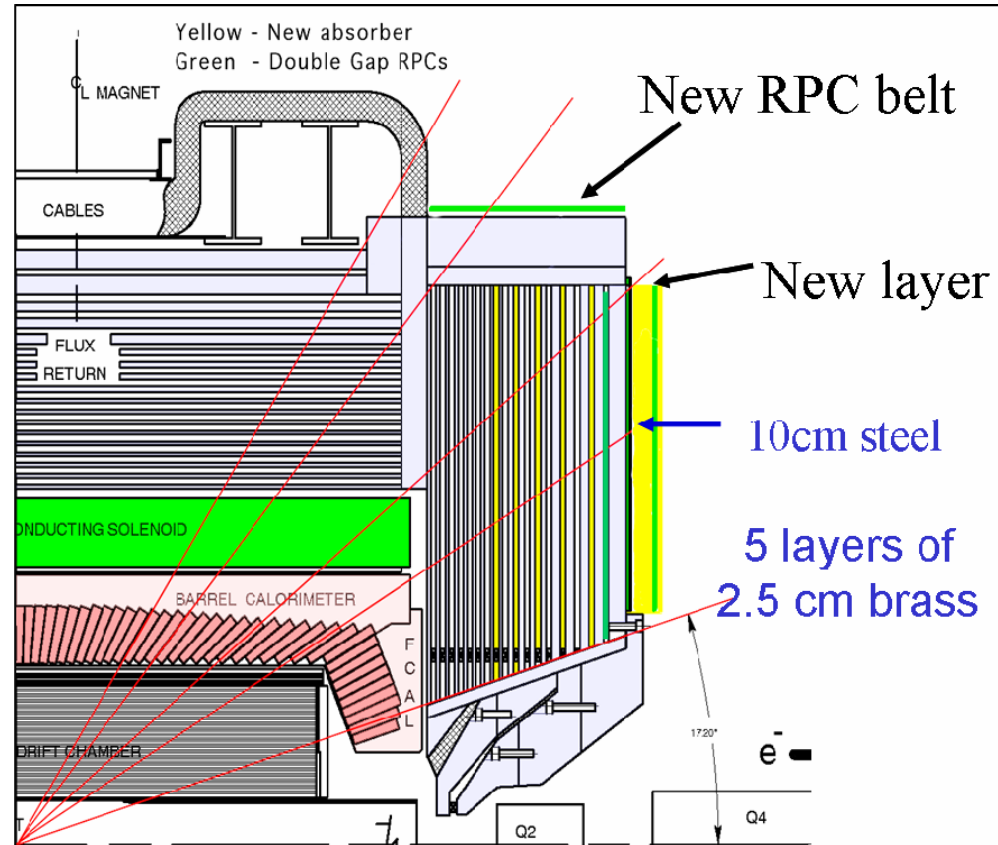


We continue to improve *BaBar* !



Problems with Resistive-Plate Muon Chambers

- The only major problem with the *BaBar* detector, known since 1999: RPC efficiency deteriorating at $\sim 10\text{-}20\%/ \text{year}$.
- Good muon and K_L identification efficiency essential in many searches for New Physics at *BaBar*.
- Forward End Cap upgraded in 2002:
- The Barrel has $5.1 \lambda_{\text{int}}$ with Layer 19 RPCs (dying, inaccessible), $4.5 \lambda_{\text{int}}$ without – marginal for a muon system.
- Adding six 2.2 cm layers of brass increases barrel thickness to $5.3 \lambda_{\text{int}}$.
- The technology chosen for Barrel RPC replacement is the **Limited Streamer Tube (LST)**. Installed in 2004-06.

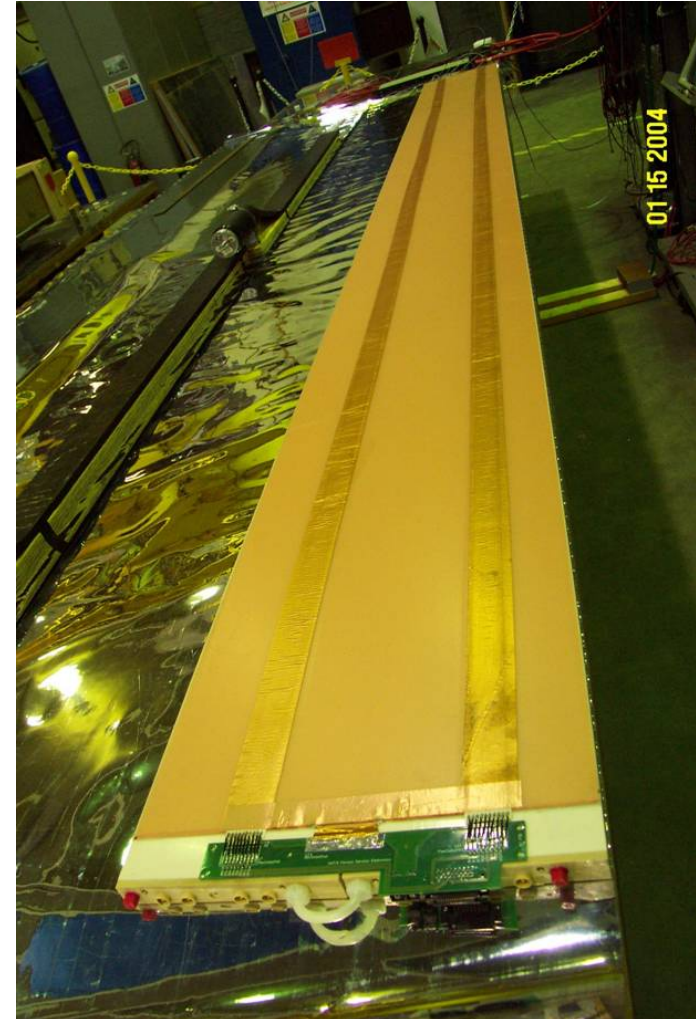




LST technology



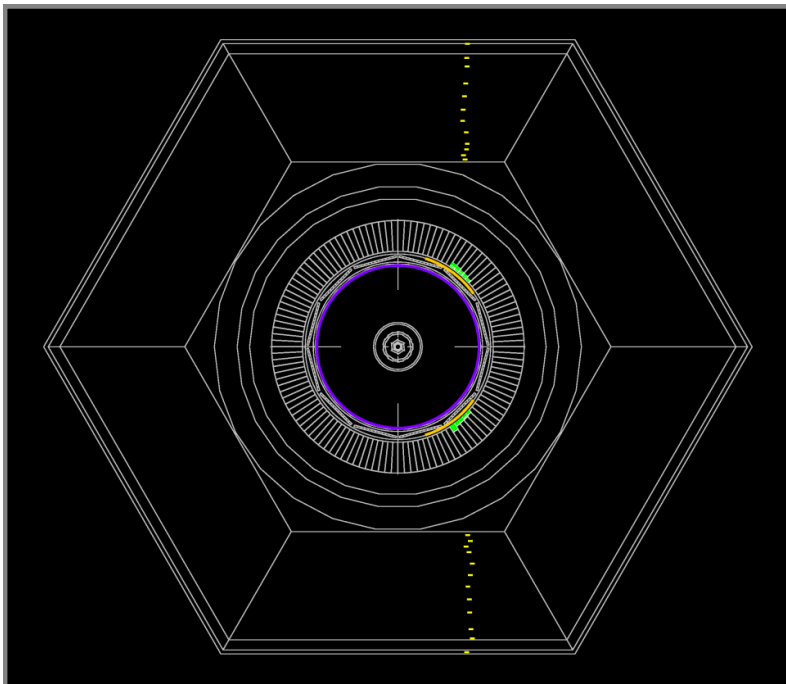
- LST is a wire chamber operating in the self-limiting streamer mode, so
- Signal does not depend on the amount of initial ionization.
- Non-flammable gas (CO_2 / $i\text{-C}_4\text{H}_{10}$ /Ar).
- 17x15 mm cells with three walls covered with conductive paint (graphite/PVAC).
- Tubes with 7 or 8 cell, 13 to 20 tubes per layer.
- phi position read off the wires (4 channels per tube), 94% eff., multiplicity mostly 1.
- Z strips span the entire width of a layer.



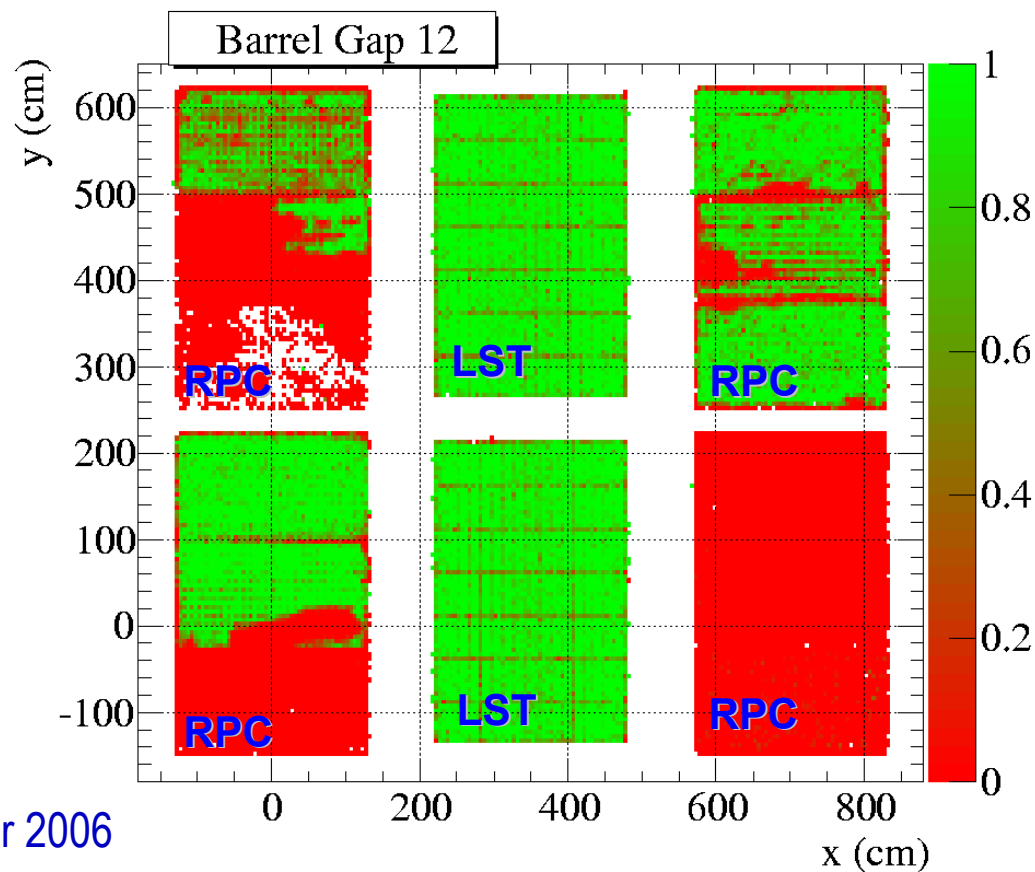


LST performance

First cosmics: run 50769, event 50170, Sep 28, 2004



Example of single-layer efficiency, May 2005



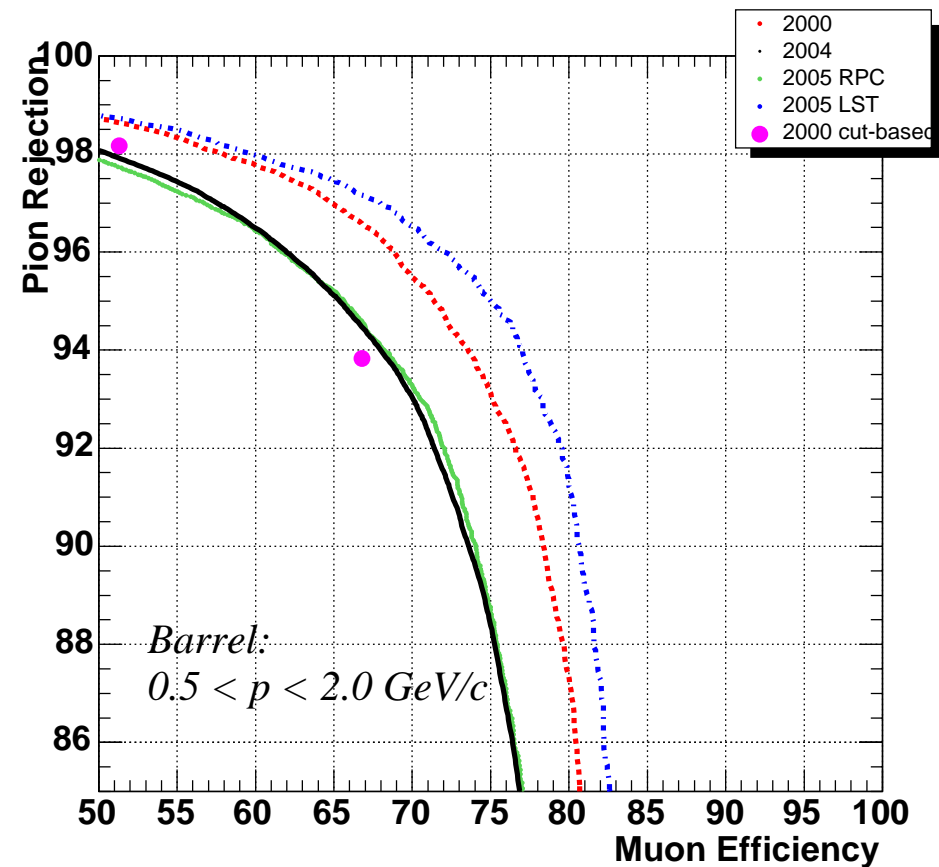
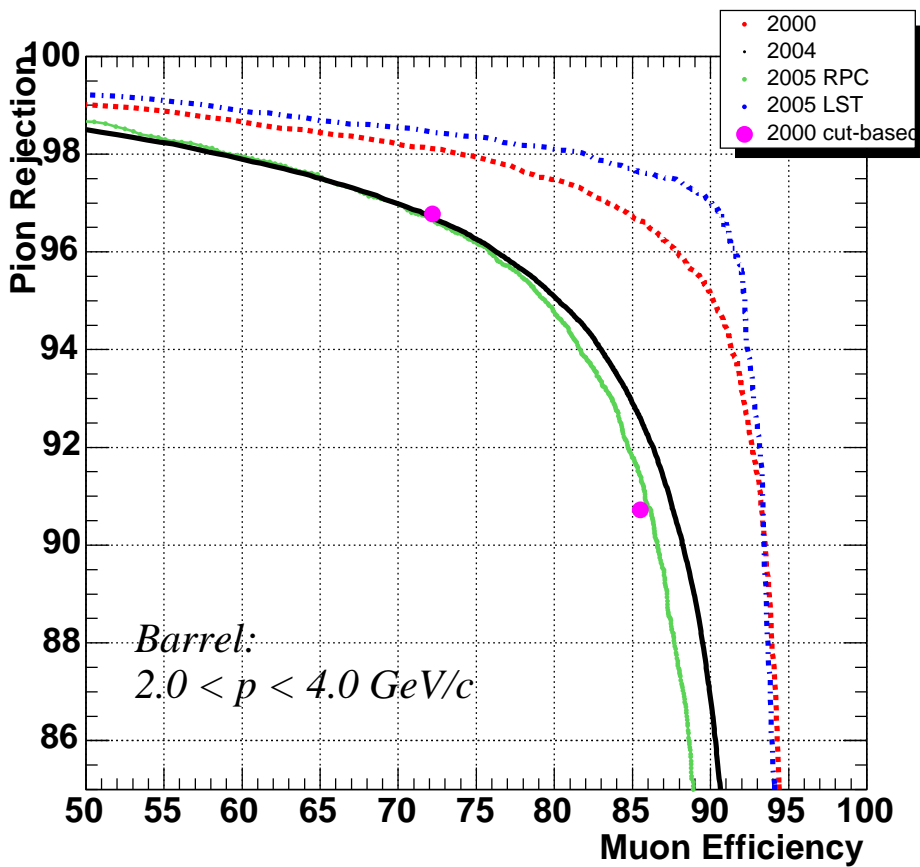
Upgrade of entire barrel completed in October 2006



Muon ID with LSTs

before many further improvements

LSTs were immediately doing better than RPCs ever did!



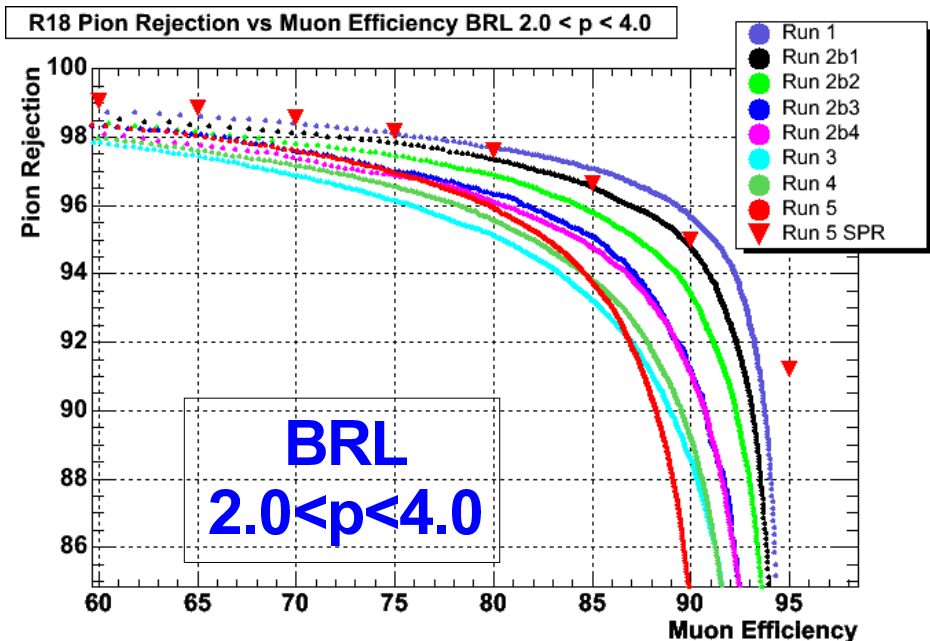
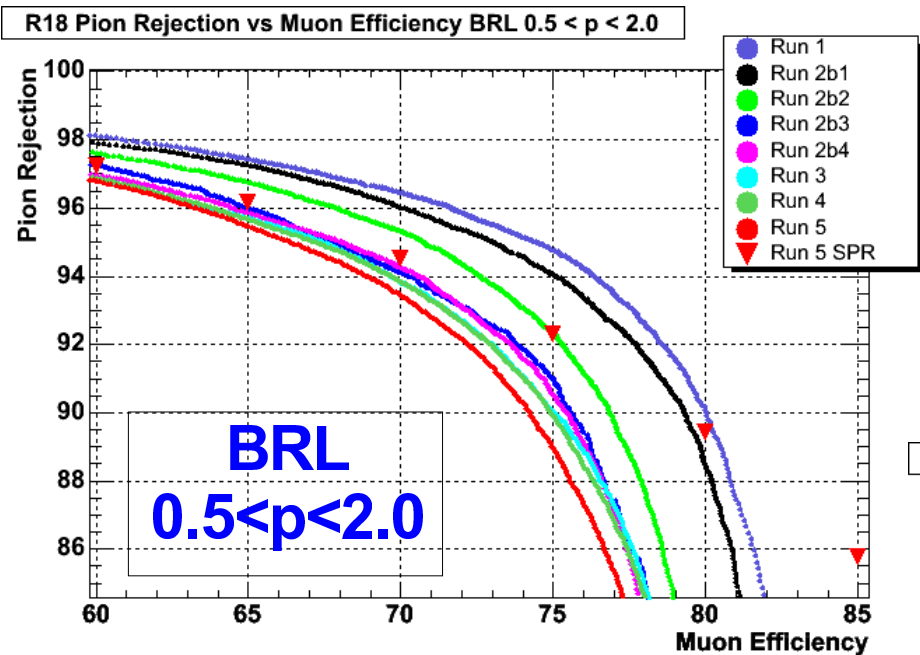


Muon ID idea:

Neural Net vs. Bagged Decision Trees

First try, some time in 2005... (Ilya Narsky, Caltech, CMS)

The more heterogeneous the inputs and/or the number of discriminating variables, the greater the advantage of Decision Trees over Neural Nets





Where does the room for improvement come from ?

1) Better parameterization of detector response:

- DCH dE/dx significantly improved
- SVT dE/dx greatly improved

Existing likelihood-based selectors will benefit “automatically” from the new dE/dx

First physics publications to use the new dE/dx parameterizations:

“Observation of CP Violation in $B^0 \rightarrow K^+ \pi^-$ and $B^0 \rightarrow \pi^+ \pi^-$ ”, PRL 99, 021603 (2007)

“Evidence for D^0 -anti- D^0 Mixing”, PRL 98, 211802 (2007)

“Search for D^0 -anti- D^0 mixing using doubly flavor tagged semileptonic decay modes”, PRD 76, 014018 (2007)

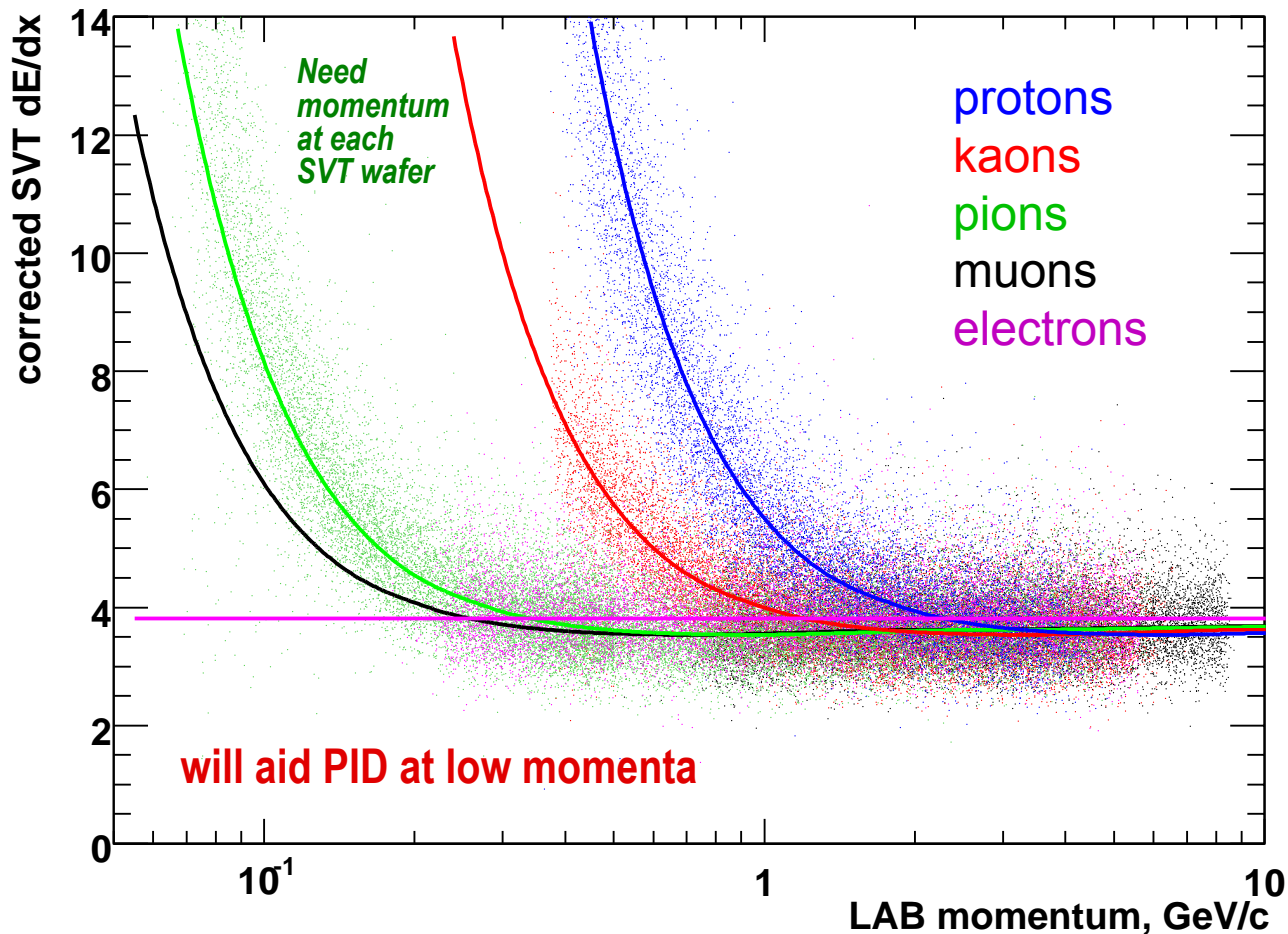
More analyses in the pipeline: all $\tau \rightarrow KX$, all charm-mixing with a D^{*+} , $B^+ \rightarrow K_s h^+$, ...



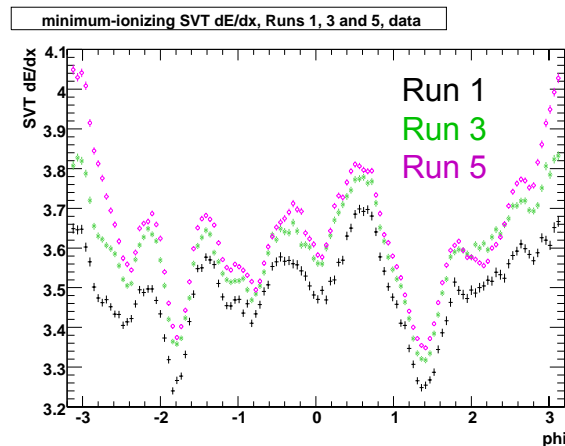
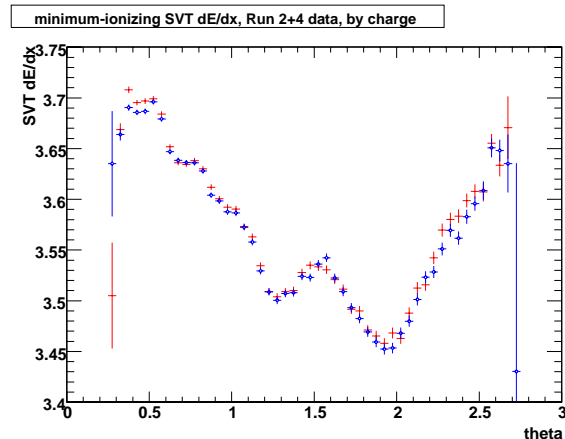
dE/dx in *BaBar* Silicon Vertex Tracker: π/e , K/e separation; stand-alone SVT tracking

This is what SVT dE/dx looks like after the new detailed calibration

SVT dE/dx for various particle types, Run 3, data



Suffers from the same systematics as DCH dE/dx , only worse





Where does the room for improvement come from ?

2) previously unused quantities (“weak classifiers”), e.g.:

- number of signal and background photons in the DIRC (in a non-“binary” fashion)
- last layer hit in DCH (for kaon and pion decays in flight)
- DIRC and DCH dE/dx for muon ID
- full set of EMC quantities for all particle types
- SVT dE/dx for electron ID
- Flattening of the training-sample spectra in theta, phi and charge allows the use of these variables as input parameters

3) new quantities:

- longitudinal EMC shower depth
- using geometry to predict “dead spots” in detector acceptance

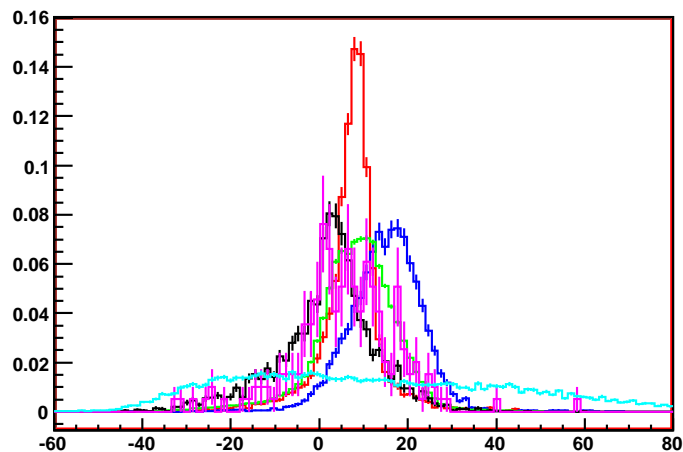
4) advanced statistical techniques from StatPatternRecognition

- SPR developed by Ilya Narsky (Caltech), SourceForge SPR V06-00-02
- we use SPR Boosted Decision Trees (muonBDT and kaonBDT selectors)
- and SPR multi-class learner (electron/pion/kaon/proton “KM” selectors)

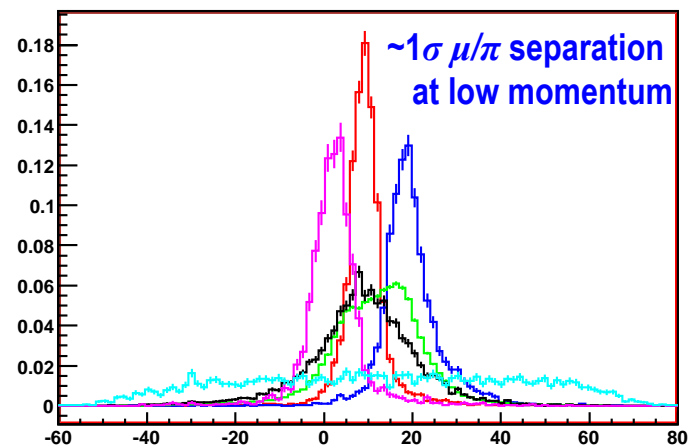


EMC longitudinal shower depth

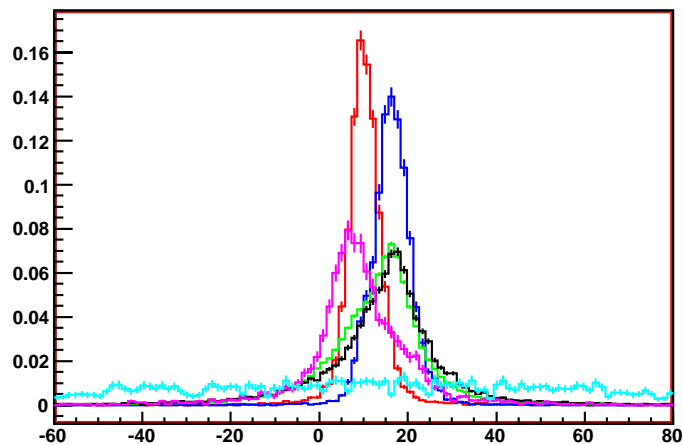
Shower Depth, $0.00 < \text{Track P} < 0.40 \text{ GeV:cm}$



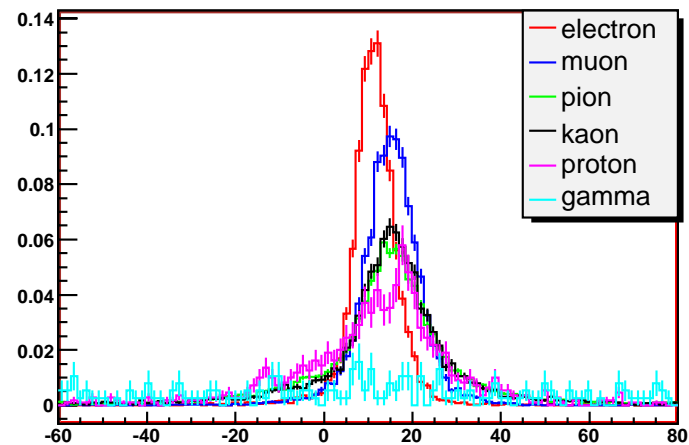
Shower Depth, $0.40 < \text{Track P} < 0.60 \text{ GeV:cm}$



Shower Depth, $0.60 < \text{Track P} < 1.25 \text{ GeV:cm}$



Shower Depth, $1.25 < \text{Track P} < 10.00 \text{ GeV:cm}$

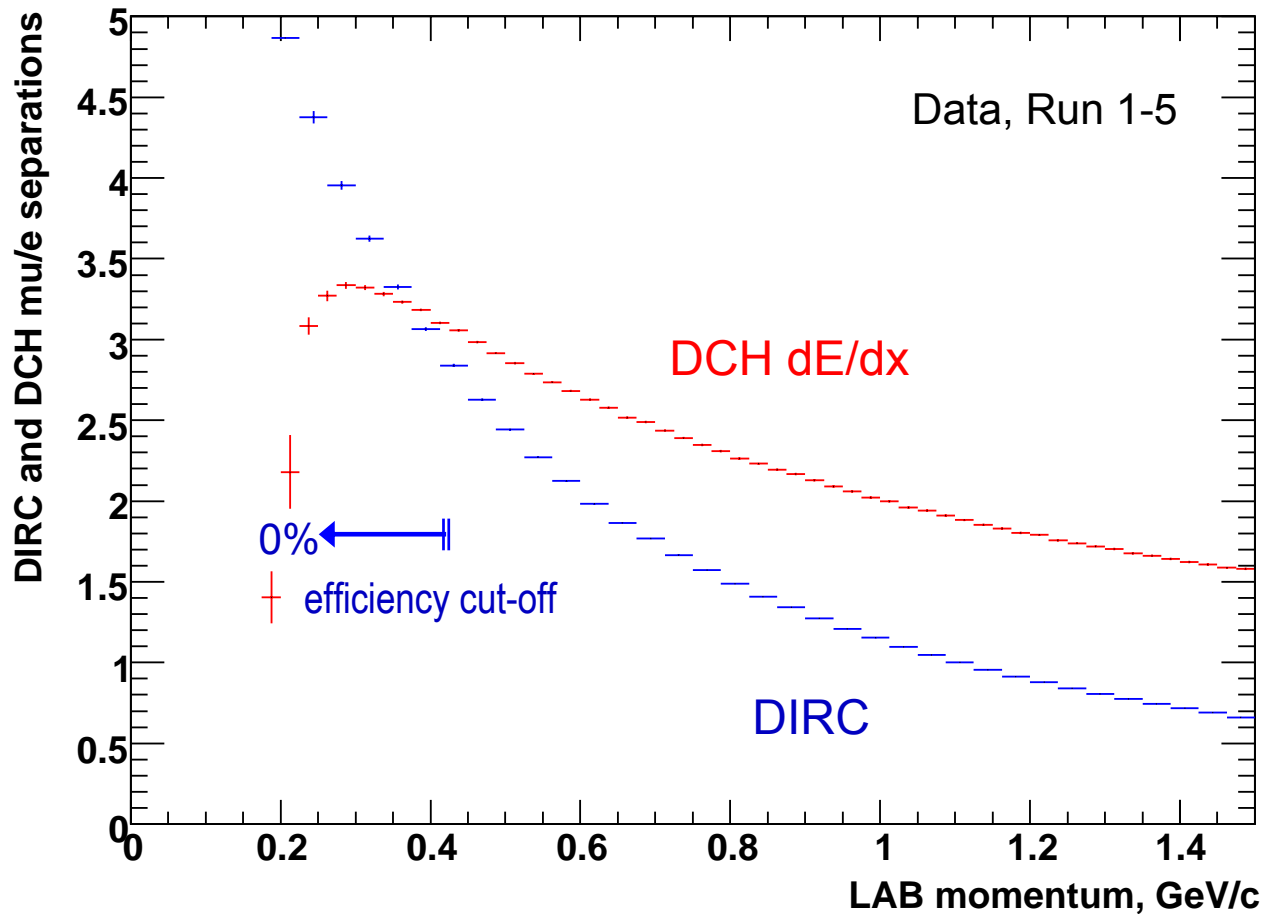


a NIM paper is planned



DCH dE/dx $\pi, \mu/e$ separation

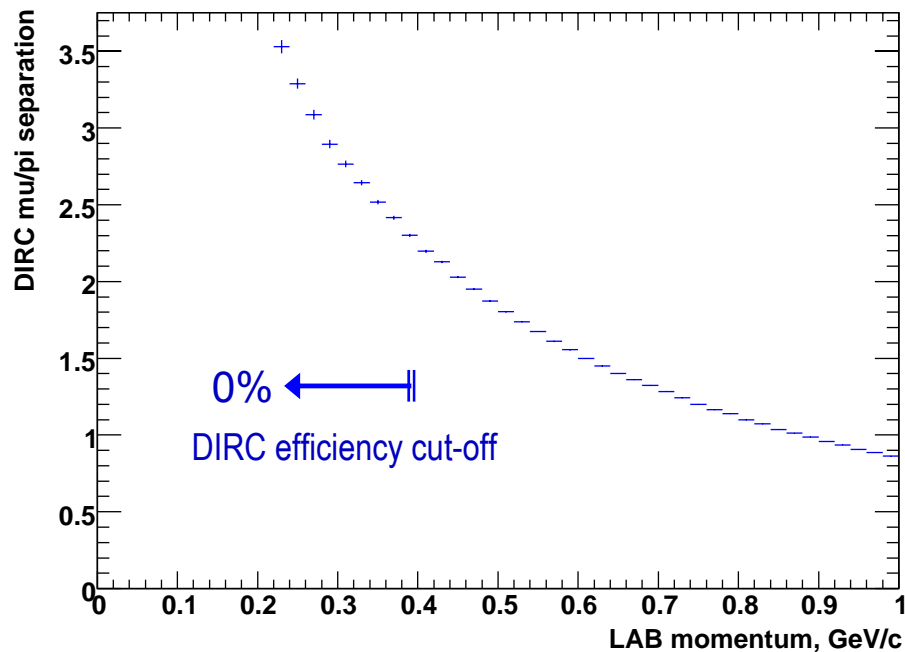
DIRC and DCH μ/e separation





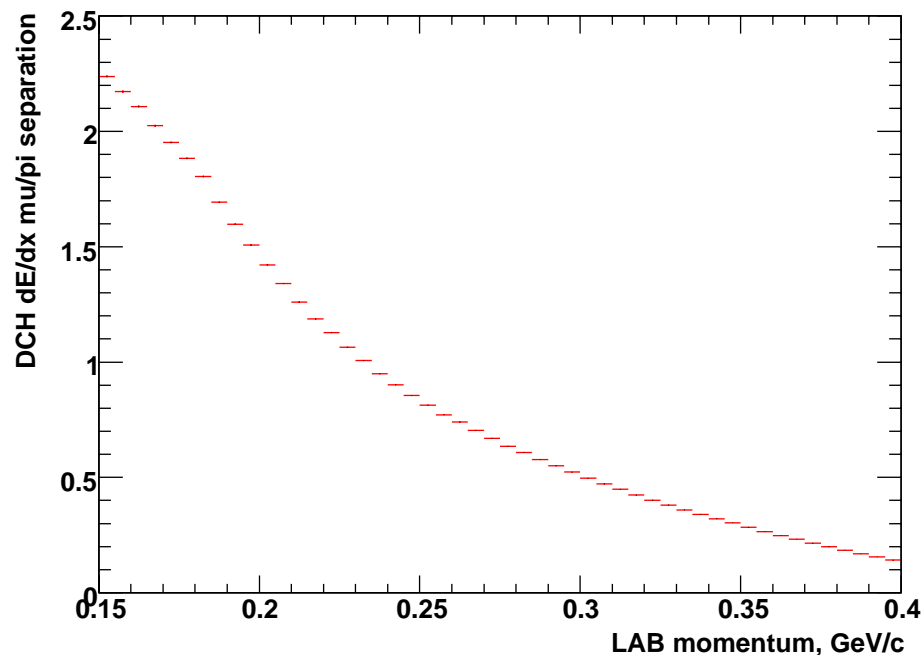
DIRC and DCH dE/dx π/μ separation

DIRC mu/pi separation, data



Muons below ~ 0.9 GeV/c do not have enough energy to punch through the IFR;
Until now, neither the DIRC nor the DCH were used in the muon selectors, nor a full set of EMC quantities

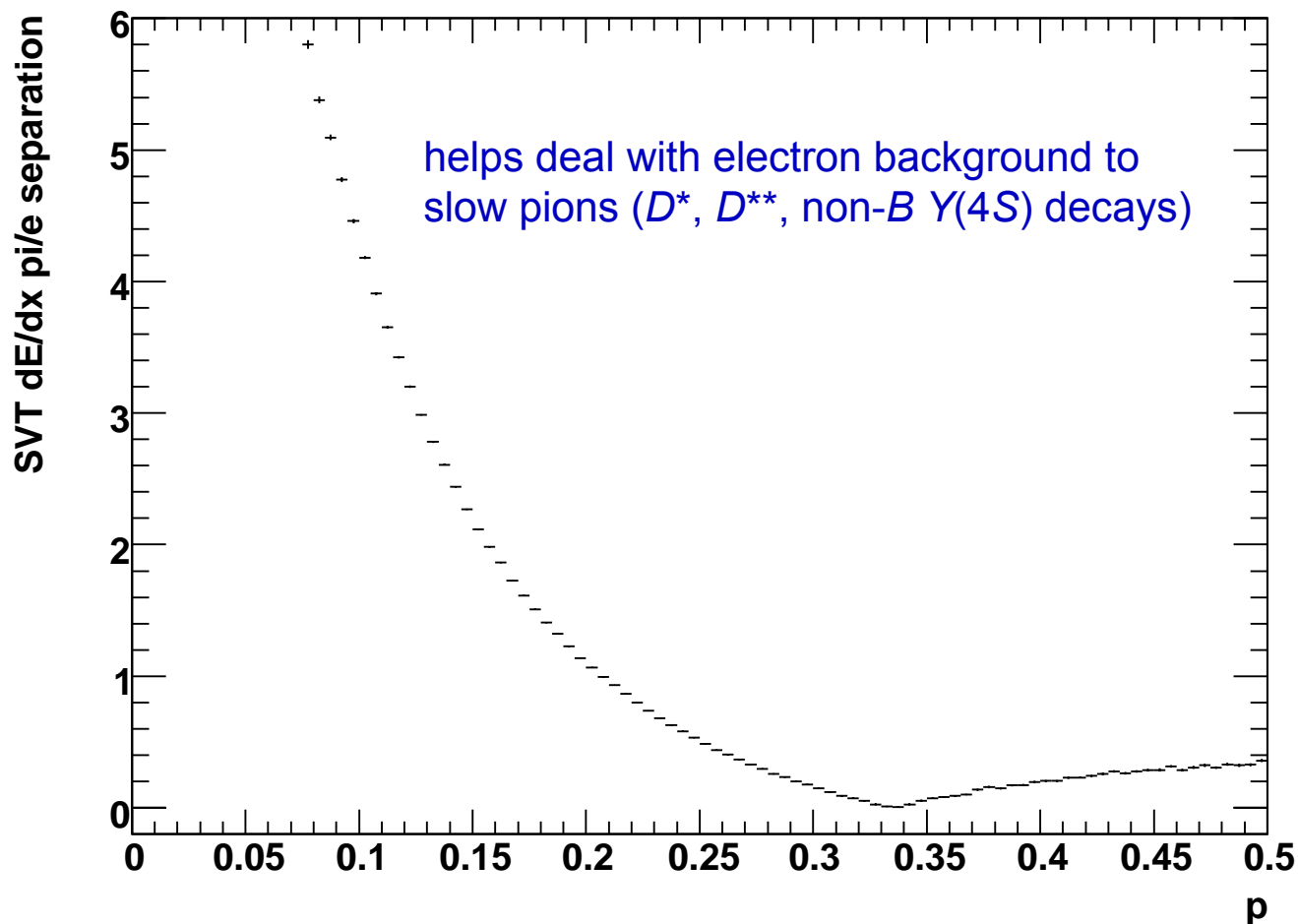
DCH dE/dx mu/pi separation, data





SVT dE/dx π/e separation

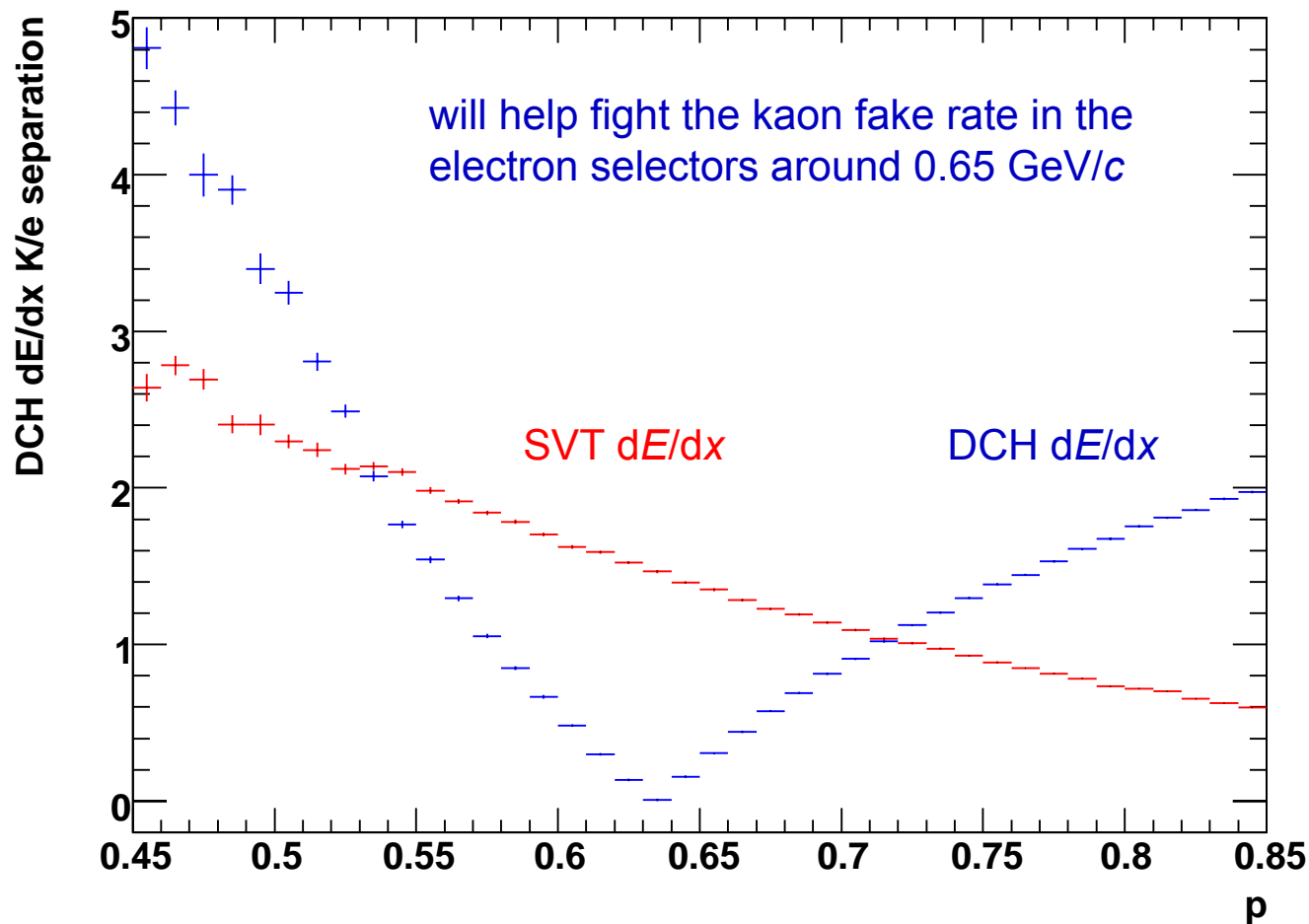
SVT dE/dx π/e separation, Run 3 data





SVT dE/dx K/e separation

DCH and SVT dE/dx K/e separation, data

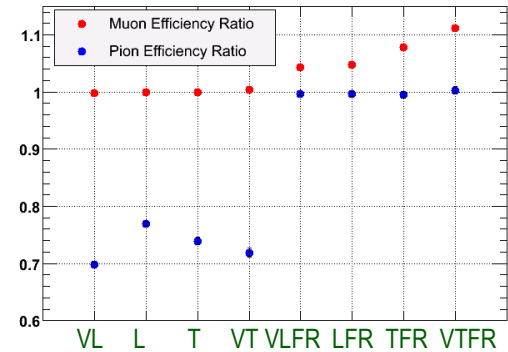




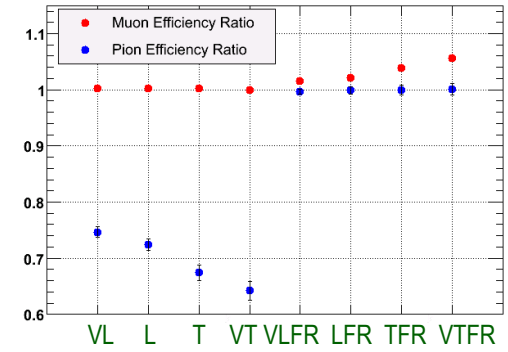
"Ultimate" muon ID now: 30 discriminating variables used

- p
- theta
- charge
- ifrns
- ifrmatchchi2
- ifrfitchi2
- ifrcont
- ifrsigmu
- ifrmeasintlen
- ecal
- lmom
- zmom20
- zmom42
- ncry
- s1s9
- s9s25
- secmom
- dEdxdchPullmu
- drcmuprob
- smsdrcmuprob
- emcdepth
- date
- drcpiprob
- drckprob
- nphot
- ndch
- lhit
- ifrcrackphi
- deltalambda
- ecaldivp

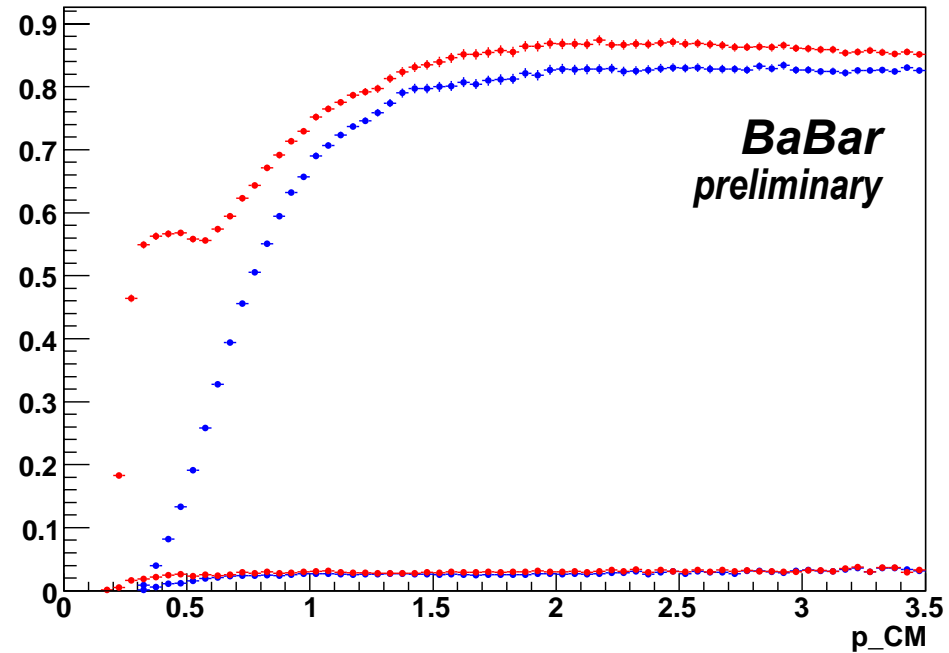
Barrel BDT/NN Efficiency Ratios (Runs 1-6, $1.5 < p < 4.0$ GeV/c)



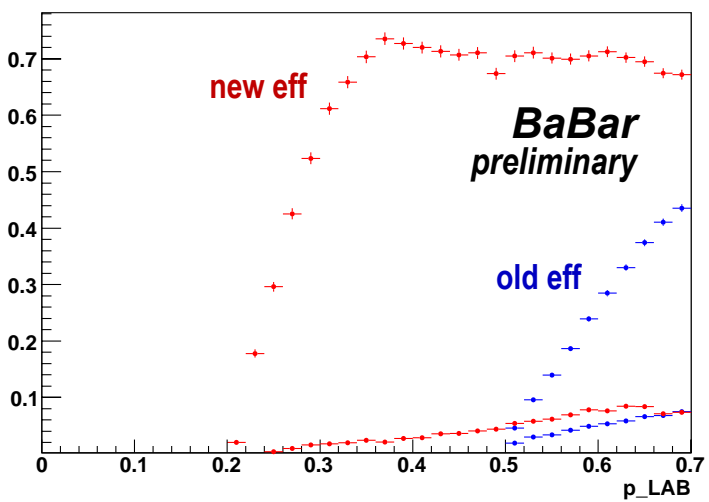
Forward BDT/NN Efficiency Ratios (Runs 1-6, $1.5 < p < 4.0$ GeV/c)



μ NN (blue) and BDT (red) eff and π misID rate (LooseFakeRate)

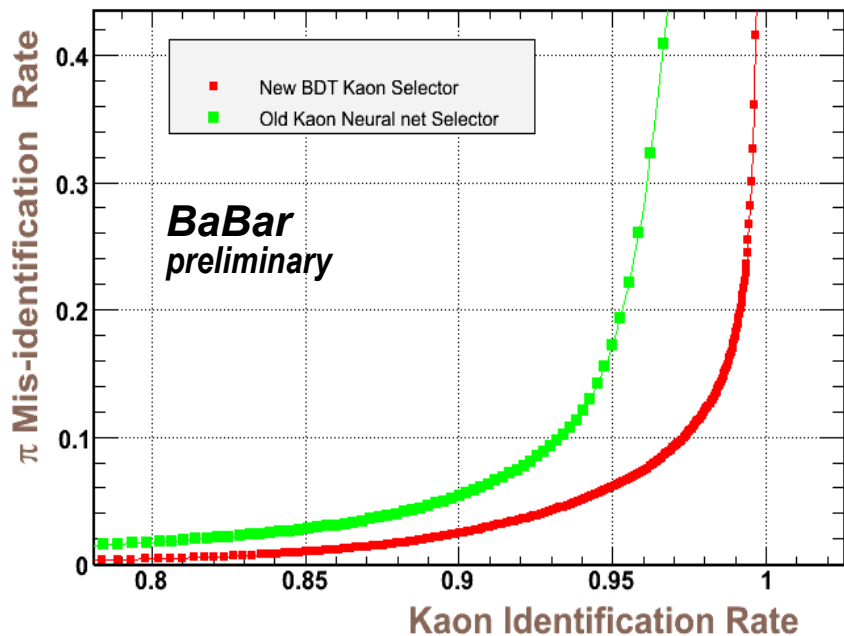


μ NN VeryLoose (blue) and BDTLoP Loose (red) eff and π misID rate



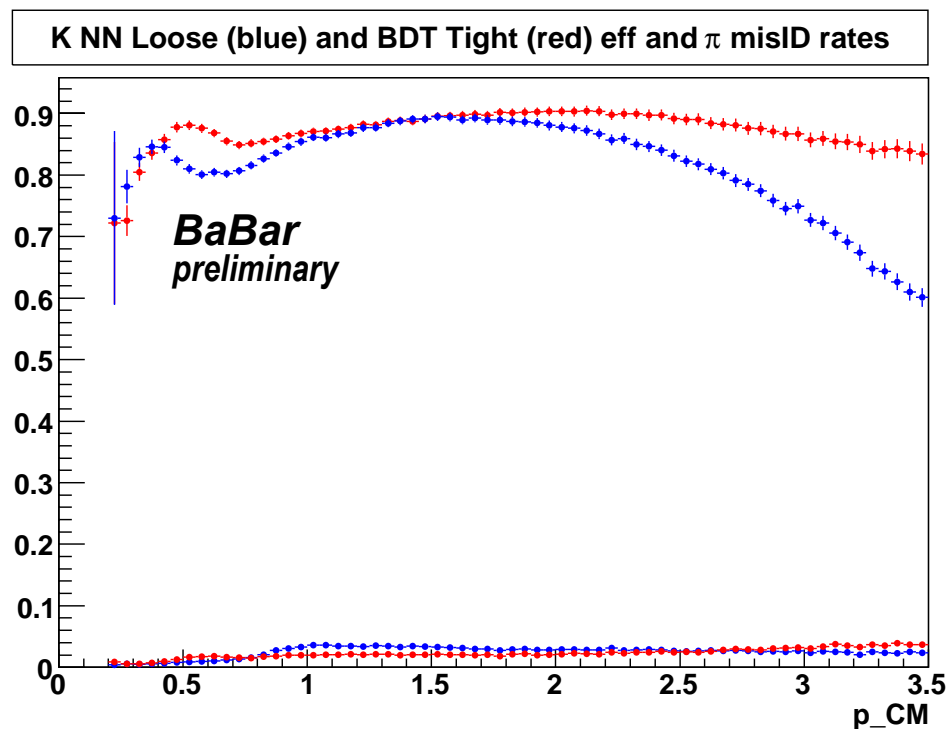


kaonBDT selector: replacement for kaonNN



Plot made for D^* pion and kaon spectra, relevant to B^0 flavor tagging

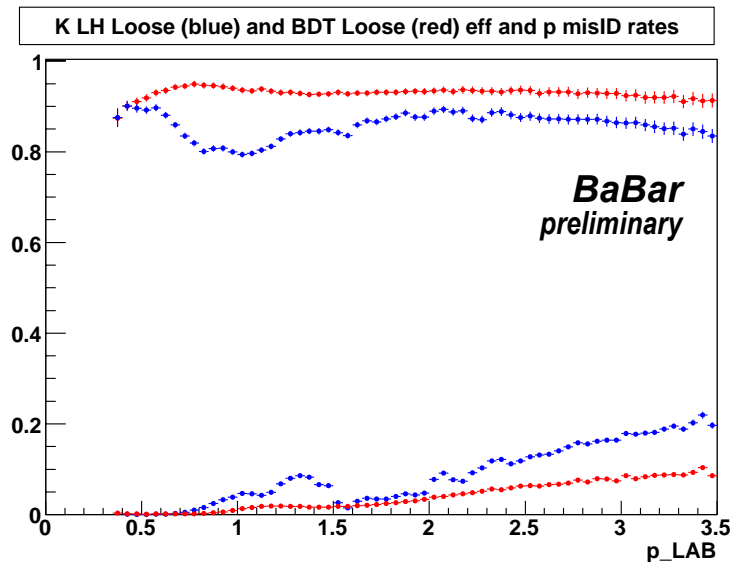
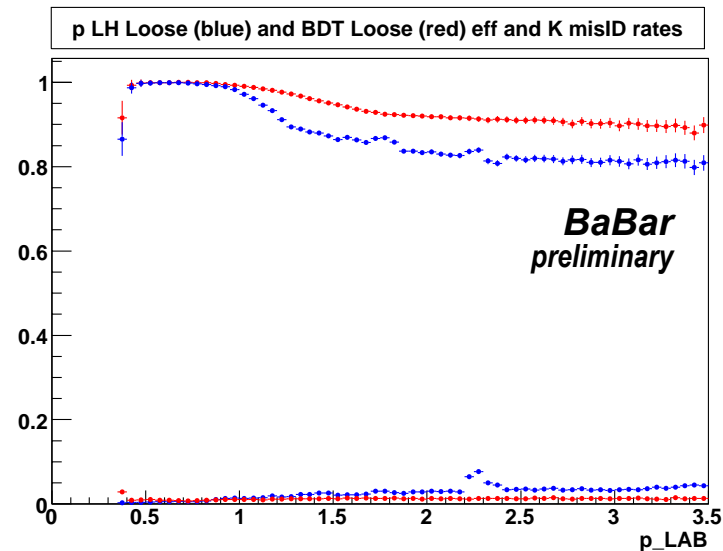
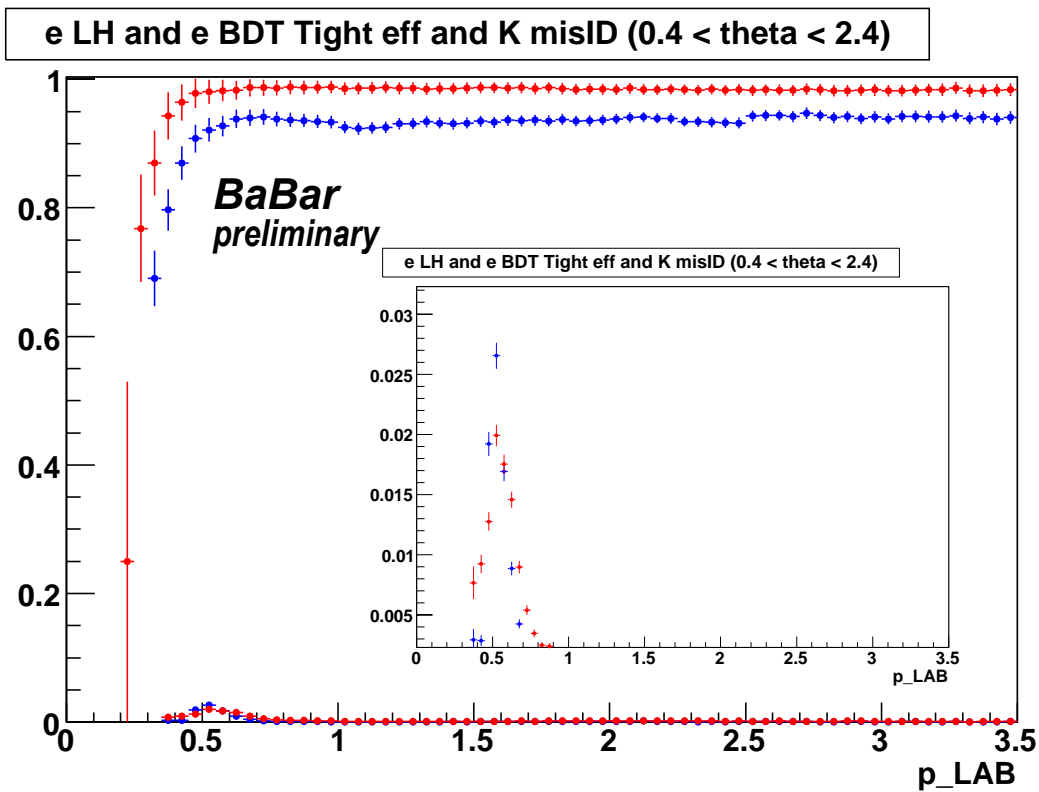
To be presented on February 11 at Caltech workshop on multivariate analysis





More new selectors

How much would you pay for 6% more electrons? ☺





The Motto

**Make impossible measurements possible
and make the possible measurements first
(or the best)**

**by gaining detailed understanding
of your detector;**

Make it in a way the entire Collaboration would benefit



The future: physics at the Large Hadron Collider

There is little doubt that previously unseen phenomena will be observed at LHC.

My primary goal for the next several years is to make important contributions to measurements that are key to establishing a pattern among these new phenomena.

So, we see something that could be a Higgs. What is it that we see?

- Is it responsible for the Electro-Weak Symmetry Breaking? Are the WW_h and ZZ_h couplings SM-like, $\sim O(1)$? Important: $h \rightarrow \gamma\gamma, bb, WW^*$.
- Are there other particles (Higgses?) responsible for the large top-quark mass? Does the tth coupling agree with SM? Important: $h \rightarrow \gamma\gamma, bb$.
- Are there other Higgses? $H/A \rightarrow \tau^+\tau^-$? $H^+ \rightarrow tb$?

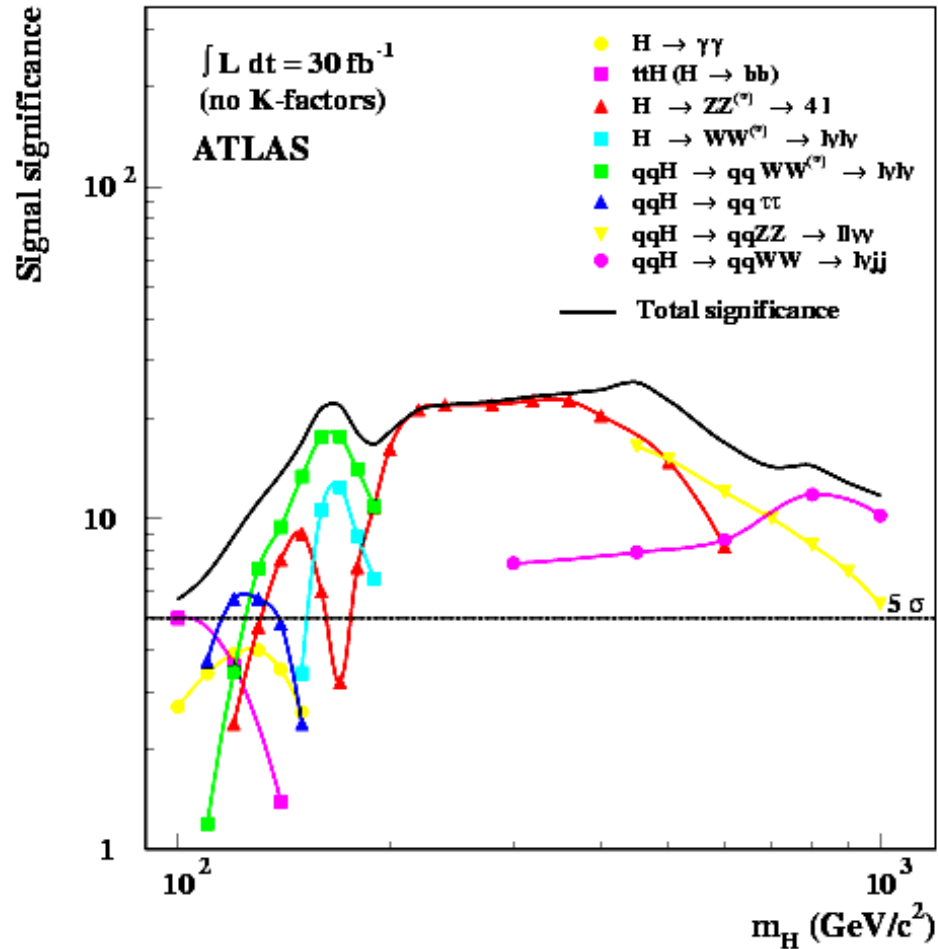
How can we use constraints on New Physics from B, B_s, K, D decays to be smart about the searches we focus on?



Answering these questions will be a challenge

Excellent understanding of the detector and sophisticated tools will be required

alignment
tracking
muons
electrons
 b jets
 τ jets
momentum scale
etc.

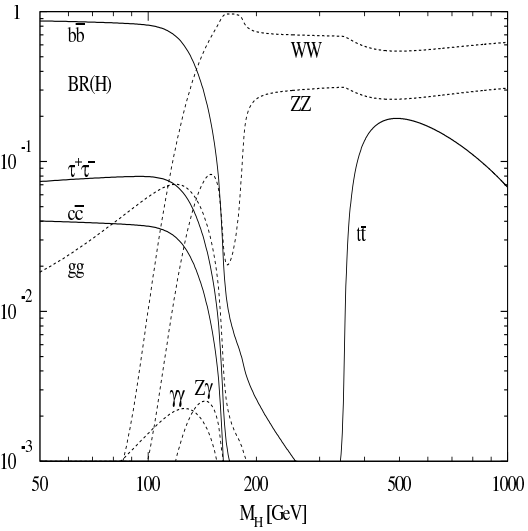
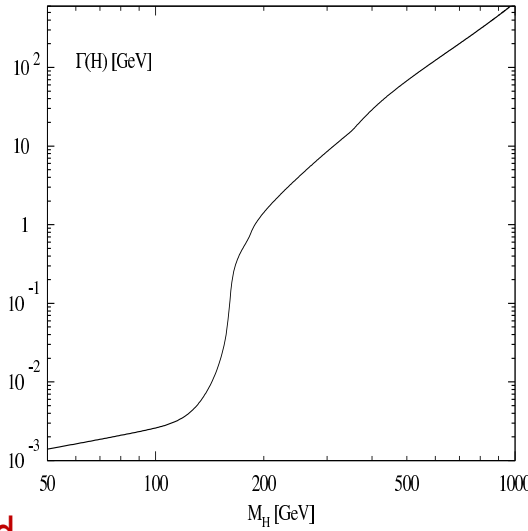
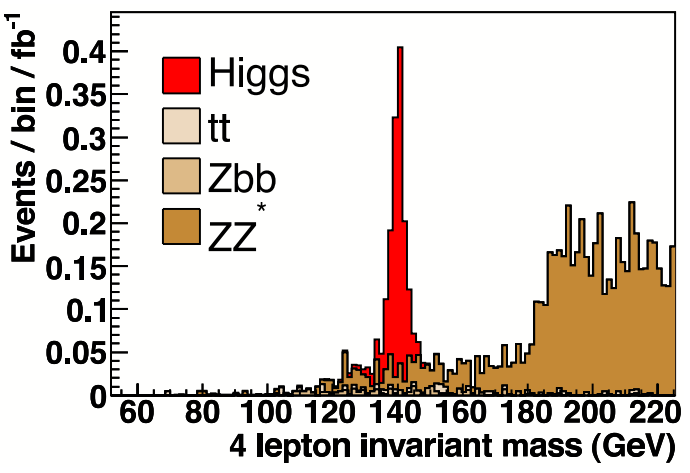




Muons at ATLAS: not merely a discovery tool!

$$H \rightarrow ZZ(*) \rightarrow e^+e^-\mu^+\mu^-$$

If Nature is kind, LHC may be able to make a precision measurement of Higgs mass, and perhaps even width



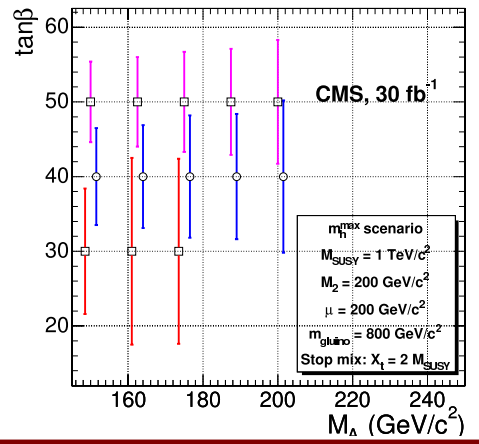
Excellent muon isolation/identification and momentum determination will be essential:

- alignment
- magnetic field map
- unbiased momentum determination
- multivariate "particle ID" (suppression of fakes)?

$$\Gamma(H/A \rightarrow \mu^+\mu^-):$$

Calibration sample:

$$Z \rightarrow \mu^+\mu^-$$



Not to mention scenarios with a Z' or other dimuon resonances...